# **Engineering functional quantum algorithms**

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Suppose that a quantum circuit with *K* elementary gates is known for a unitary matrix *U*, and assume that  $U^m$  is a scalar matrix for some positive integer  $m$ . We show that a function of  $U$  can be realized on a quantum computer with at most  $O(mK+m^2\ln m)$  elementary gates. The functions of *U* are realized by a generic quantum circuit, which has a particularly simple structure. Among other results, we obtain efficient circuits for the fractional Fourier transform.

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Let *U* be a unitary matrix,  $U \in \mathcal{U}(2^n)$ . Suppose that a fast quantum algorithm is known for *U*, which is given by a factorization of the form

$$
U = U_1 U_2 \cdots U_K, \tag{1}
$$

where the unitary matrices  $U_i$  are realized by controlled-NOT gates or by single-qubit gates  $[1]$ . We are interested in the following question: *Are there efficient quantum algorithms for unitary matrices, which are functions of U*?

The question is puzzling, because the knowledge of the factorization  $(1)$  of *U* does not seem to be of much help in finding similar factorizations for, say,  $V = U^{1/3}$ . The purpose of this paper is to give an answer to the above question for a wide range of unitary matrices *U*.

Our solution to this problem is based on a generic circuit which implements arbitrary functions of *U*, assuming that *U<sup>m</sup>* is a scalar matrix for some positive integer *m*. If *m* is small (that is, polylogarithmic in  $n$ ), then our method provides an efficient quantum circuit for *V*.

*Notations.* We denote by  $U(m)$  the group of unitary  $m \times m$  matrices, by 1 the identity matrix, and by C the field of complex numbers.

## **I. PRELIMINARIES**

We recall some standard material on matrix functions, see Refs.  $[2-4]$  for more details. Let *U* be a unitary matrix. The spectral theorem states that *U* is unitarily equivalent to a diagonal matrix *D*, that is,  $U = TDT^{\dagger}$  for some unitary matrix *T*. The elements  $\lambda_i$  on the diagonal of *D*  $=$ diag( $\lambda_1$ , ..., $\lambda_{2^n}$ ) are the eigenvalues of *U*.

Let  $f$  be any function of complex scalars such that its domain contains the eigenvalues  $\lambda_i$ ,  $1 \le i \le 2^n$ . The matrix function  $f(U)$  is then defined by

$$
f(U) = T \operatorname{diag}(f(\lambda_1), \ldots, f(\lambda_{2^n}))T^{\dagger},
$$

where *T* denotes the diagonalizing matrix of *U*, as above.

Notice that any two scalar functions *f* and *g*, which take the same values on the spectrum of *U*, yield the same matrix value  $f(U) = g(U)$ . In particular, one can find an interpolation polynomial *g*, which takes the same values as *f* on the eigenvalues  $\lambda_i$ . It is possible to assume that the degree of *g* is smaller than the degree of the minimal polynomial of *U*. In other words,  $V = f(U)$  can be expressed by a linear combination of integral powers of the matrix *U*,

$$
V = f(U) = \sum_{i=0}^{m-1} \alpha_i U^i,
$$
 (2)

where *m* is the degree of the minimal polynomial of the matrix *U*, and  $\alpha_i \in \mathbb{C}$  for  $i=0, \ldots, m-1$ . In order for *V* to be unitary, it is necessary and sufficient that the function *f* maps the eigenvalues  $\lambda_i$  of *U* to elements on the unit circle.

*Remark.* There exist several different definitions for matrix functions. The relationship between these definitions is discussed in detail in Ref.  $[5]$ . We have chosen the most general definition that allows to express the function values by polynomials.

# **II. THE GENERIC CIRCUIT**

Let *U* be a unitary  $2^n \times 2^n$  matrix with minimal polynomial of degree *m*. We assume that an efficient quantum circuit is known for *U*. How can we go about implementing the linear combination (2)? We will use an ancillary system of  $\mu$ quantum bits, where  $\mu$  is chosen such that  $2^{\mu-1} < m \le 2^{\mu}$ holds. This will allow us to create the linear combination by manipulating somewhat larger matrices, which on input  $|0\rangle$  $\otimes |\psi\rangle \in \mathbb{C}^{2^{\mu}} \otimes \mathbb{C}^{2^{n}}$  produce the state  $|0\rangle \otimes V|\psi\rangle$ .

We first bring the ancillary system into a superposition of the first *m* computational base states, such that an input state  $|0\rangle \otimes |\psi\rangle \in \mathbb{C}^{2^{\mu}} \otimes \mathbb{C}^{2^n}$  is mapped to the state

$$
\frac{1}{\sqrt{m}} \sum_{i=0}^{m-1} |i\rangle \otimes |\psi\rangle.
$$
 (3)

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FIG. 1. A quantum circuit realizing the block diagonal matrix  $A = \text{diag}(1, U, U^2, \ldots, U^{2^{\mu-1}}).$ 

This can be done by acting with a  $2^{\mu} \times 2^{\mu}$  unitary matrix *B* on the ancillary system, where the first column of *B* is of the form  $1/\sqrt{m}(1, \ldots, 1, 0, \ldots, 0)^t$ . Efficient implementations of *B* exist.

Notice that there exists an efficient implementation of the block diagonal matrix  $A = diag(1, U, U^2, \ldots, U^{2^{\mu}-1})$ . Indeed, *A* can be composed of the matrices  $U^{2^{\eta}}$ ,  $0 \le \eta \le \mu$ , conditioned on the  $\mu$  ancillae bits. The resulting implementation is shown in Fig. 1. The state  $(3)$  is transformed by this circuit into the state

$$
\frac{1}{\sqrt{m}} \sum_{i=0}^{m-1} |i\rangle \otimes U^i |\psi\rangle.
$$
 (4)

In the next step, we let a  $2^{\mu} \times 2^{\mu}$  matrix *M* act on the ancillae bits. We choose  $M$  such that the state  $(4)$  is mapped to

$$
\frac{1}{\sqrt{m}} \sum_{k=0}^{m-1} |k\rangle \otimes U^k V |\psi\rangle.
$$
 (5)

It turns out that *M* can be realized by a unitary matrix, assuming that the minimal polynomial of *U* is of the form *x<sup>m</sup>*  $-\tau$ ,  $\tau \in \mathbb{C}$ . This will be explained in some detail in the following section.

We apply the inverse  $A^{\dagger}$  of the block diagonal matrix A. This transforms the state  $(5)$  to

$$
\frac{1}{\sqrt{m}} \sum_{k=0}^{m-1} |k\rangle \otimes V|\psi\rangle.
$$
 (6)

We can clean up the ancillae bits by applying the  $2^{\mu} \times 2^{\mu}$ matrix  $B^{\dagger}$ . This yields then the output state

$$
|0\rangle \otimes V|\psi\rangle = |0\rangle \otimes f(U)|\psi\rangle. \tag{7}
$$

The steps from the input state  $|0\rangle \otimes |\psi\rangle$  to the final output state  $|0\rangle \otimes V|\psi\rangle$  are illustrated in Fig. 2 for the case  $\mu=2$ .



FIG. 2. Generic circuit realizing a linear combination *V*. The case  $\mu$ =2 is shown.

The following theorem gives an upper bound on the complexity of the method. We use the number of elementary gates (that is, the number of single-qubit gates and controlled-NOT gates) as a measure of complexity.

*Theorem 1*. Let *U* be a  $2^n \times 2^n$  unitary matrix with minimal polynomial  $x^m - \tau$ ,  $\tau \in \mathbb{C}$ . Suppose that there exists a quantum algorithm for *U* using *K* elementary gates. Then a unitary matrix  $V = f(U)$  can be realized with at most  $O(mK+m^2\ln m)$  elementary operations.

*Proof.* A matrix acting on  $\mu \in O(\ln m)$  qubits can be realized with at most  $O(m^2 \ln m)$  elementary operations, cf. Ref. [1]. Therefore, the matrices  $B, B^{\dagger}$ , and *M* can be realized with a total of at most  $O(3m^2 \ln m)$  operations.

If *K* operations are needed to implement *U*, then at most 14*K* operations are needed to implement  $\Lambda_1(U)$ , the operation *U* controlled by a single qubit. The reason is that a doubly controlled-NOT gate can be implemented with 14 elementary gates  $[6]$ , and a controlled single-qubit gate can be implemented with six or fewer elementary gates  $[1]$ .

We observe that  $2^{\mu}-1$  copies of  $\Lambda_1(U)$  suffice to implement *A*. Indeed, we certainly can implement  $\Lambda_1(U^{2^k})$  by a sequence of  $2^k$  circuits  $\Lambda_1(U)$ . This bold implementation yields the estimate for *A*. Typically, we will be able to find much more efficient implementations. Anyway, we can conclude that *A* and  $A^{\dagger}$  can both be implemented by at most  $14(2^{\mu}-1)K \in O(14mK)$  operations. Combining our counts yields the result.

## **III. UNITARITY OF THE MATRIX** *M*

It remains to show that the state  $(4)$  can be transformed into the state  $(5)$  by acting with a unitary matrix *M* on the system of  $\mu$  ancillae qubits. This is the crucial step in the previously described method.

Let *U* be a unitary matrix with a minimal polynomial of degree *m*. A unitary matrix  $V = f(U)$  can then be represented by a linear combination

$$
V = \sum_{i=0}^{m-1} \alpha_i U^i.
$$
 (8)

We will motivate the construction of the matrix *M* by examining in some detail the resulting linear combinations of the matrices  $U^k V$ . From Eq. (8), we obtain

$$
U^{k}V = \sum_{i=0}^{m-1} \alpha_{i}U^{i+k}.
$$
 (9)

Suppose that the minimal polynomial of *U* is of the form  $m(x) = x^m - g(x)$ , with  $g(x) = \sum_{i=0}^{m-1} g_i x^i$ . The right-hand side of Eq.  $(9)$  can be reduced to a polynomial in *U* of degree less than *m* using the relation  $U^m = g(U)$ :

$$
U^k V = \sum_{i=0}^{m-1} \beta_{ki} U^i.
$$

The coefficients  $\beta_{ki}$  are explicitly given by

$$
(\beta_{k0},\beta_{k1},\ldots,\beta_{k(m-1)})=(\alpha_0,\alpha_1,\ldots,\alpha_{m-1})P^k,
$$

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where *P* denotes the companion matrix of  $m(x)$ , that is,

$$
P = \begin{pmatrix} 0 & 1 & 0 & \cdots & 0 \\ 0 & 0 & 1 & \cdots & 0 \\ \vdots & \vdots & \vdots & \ddots & \vdots \\ 0 & 0 & 0 & \cdots & 1 \\ g_0 & g_1 & g_2 & \cdots & g_{m-1} \end{pmatrix}.
$$

The  $2^{\mu} \times 2^{\mu}$  matrix *M* is defined by

$$
M = \begin{pmatrix} C & 0 \\ 0 & 1 \end{pmatrix},
$$

where  $C = (\beta_{ki})_{k,i=0,\dots,m-1}$ , and **1** is a  $(2^{\mu}-m) \times (2^{\mu})$  $-m$ ) identity matrix. Under the assumptions of Theorem 1, it turns out that the matrix *M* is unitary. Before proving this claim, let us formally check that the matrix *M* transforms the state  $(4)$  into the state  $(5)$ . If we apply the matrix *M* to the ancillary system, then we obtain from expression  $(4)$  the state

$$
\frac{1}{\sqrt{m}} \sum_{i=0}^{m-1} M|i\rangle \otimes U^{i}|\psi\rangle = \frac{1}{\sqrt{m}} \sum_{k,i=0}^{m-1} \beta_{ki}|k\rangle \otimes U^{i}|\psi\rangle
$$

$$
= \frac{1}{\sqrt{m}} \sum_{k=0}^{m-1} |k\rangle \otimes \sum_{i=0}^{m-1} \beta_{ki} U^{i}|\psi\rangle
$$

$$
= \frac{1}{\sqrt{m}} \sum_{k=0}^{m-1} |k\rangle \otimes U^{k}V|\psi\rangle,
$$

which coincides with the state  $(5)$ , as claimed.

*Lemma 2*. Let *U* be a unitary matrix with minimal polynomial  $m(x) = x^m - \tau$ . Let *V* be a matrix satisfying Eq. (2). If *V* is unitary, then *M* is unitary.

*Proof.* It suffices to show that the matrix *C* is unitary. Notice that the assumption on the minimal polynomial  $m(x)$ implies that *C* is of the form

$$
C=\left(\begin{array}{cccccc}\n\alpha_0 & \alpha_1 & \cdots & \alpha_{m-2} & \alpha_{m-1} \\
\tau\alpha_{m-1} & \alpha_0 & \cdots & \alpha_{m-3} & \alpha_{m-2} \\
\vdots & \vdots & \ddots & \vdots & \vdots \\
\tau\alpha_1 & \tau\alpha_2 & \cdots & \tau\alpha_{m-1} & \alpha_0\n\end{array}\right),
$$

that is, *C* is obtained from a circulant matrix by multiplying every entry below the diagonal by  $\tau$ . In other words, we have

$$
C = (\big[\tau\big]_{i > j} \alpha_{j-i \bmod m}\big)_{i,j=0,\ldots,m-1},
$$

where  $[\tau]_{i>i} = \tau$  if  $i > j$ , and  $[\tau]_{i>i} = 1$  otherwise.

Note that the inner product of row *a* with row *b* of matrix *C* is the same as the inner product of row  $a+1$  with row *b*  $+1$ . Thus, to prove the unitarity of *C*, it suffices to show that

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$$
\delta_{a,0} = \langle \text{row a} | \text{row 0} \rangle = \sum_{j=0}^{a-1} \overline{\tau} \overline{\alpha_{j-a}} \alpha_j + \sum_{j=a}^{m-1} \overline{\alpha_{j-a}} \alpha_j
$$
\n(10)

holds, where  $\delta_{a,0}$  denotes the Kronecker delta and the indices of <sup>a</sup> are understood modulo *m*.

Consider the equation

$$
\mathbf{1} = V^{\dagger} V = \left(\sum_{i=0}^{m-1} \overline{\alpha_i} U^{-i}\right) \left(\sum_{i=0}^{m-1} \alpha_i U^i\right). \tag{11}
$$

The right-hand side can be simplified to a polynomial in *U* of degree less than *m* using the identity  $\overline{\tau}U^m = 1$ . The coefficient of  $U^a$  in Eq. (11) is exactly the right-hand side of Eq.  $(10)$ . Since the minimal polynomial of *U* is of degree *m*, it follows that the matrices  $U^0, U^1, \ldots, U^{m-1}$  are linearly independent. Thus, comparing coefficients on both sides of Eq.  $(11)$  shows Eq.  $(10)$ . Hence the rows of *C* are pairwise orthogonal and of unit norm.

*A simple example.* Let  $F_n$  be the discrete Fourier transform matrix,

$$
F_n = 2^{-n/2} (\exp(-2\pi i k\ell/2^n))_{k,\ell=0,\ldots,2^{n}-1},
$$

with  $i^2 = -1$ . Recall that the Cooley-Tukey decomposition yields a fast quantum algorithm, which implements  $F_n$  with  $O(n^2)$  elementary operations. The minimal polynomial of  $F_n$ is  $x^4-1$  if  $n \ge 3$ . Thus, any unitary matrix *V*, which is a function of  $F_n$ , can be realized with  $O(n^2)$  operations.

For instance, if  $n \ge 3$ , then the fractional power  $F_n^x$ , x  $\in$  R, can be expressed as

$$
F_n^x = \alpha_0(x)I + \alpha_1(x)F_n + \alpha_2(x)F_n^2 + \alpha_3(x)F_n^3,
$$

where the coefficients  $\alpha_i(x)$  are given by (cf. Ref. [7])

$$
\alpha_0(x) = \frac{1}{2} (1 + e^{ix}) \cos x, \quad \alpha_1(x) = \frac{1}{2} (1 - ie^{ix}) \sin x,
$$
  

$$
\alpha_2(x) = \frac{1}{2} (-1 + e^{ix}) \cos x, \quad \alpha_3(x) = \frac{1}{2} (-1 - ie^{ix}) \sin x.
$$

In this case,  $F_n^x$  is realized by the circuit in Fig. 2 with *U*  $F_{n}$  and  $M = (\alpha_{j-i}(x))_{i,j=0,\ldots,3}$ . The circuit can be implemented with  $O(n^2)$  operations.

#### **IV. LIMITATIONS**

The previous sections showed that a unitary matrix  $f(U)$ can be realized by a linear combination of the powers  $U^i$ ,  $0 \le i \le m$ , if the minimal polynomial  $m(x)$  of *U* is of the form  $x^m - \tau$ ,  $\tau \in \mathbb{C}$ . One might wonder whether the restriction to minimal polynomials of this form is really necessary. The next lemma explains why we had this limitation.

*Lemma 3*. Let *U* be a unitary matrix with minimal polynomial  $m(x) = x^m - g(x)$ , deg  $g(x) \le m$ . If  $g(x)$  is not a constant, then the matrix *M* is in general not unitary.

*Proof.* Suppose that  $g(x) = \sum_{i=0}^{m-1} g_i x^i$ . We may choose for instance  $V = U^m = g(U)$ . Then the norm of first row in *M* is greater than 1. Indeed, we can calculate this norm to be  $|g_0|^2 + |g_1|^2 + \cdots + |g_{m-1}|^2$ . However,  $|g_0|^2 = 1$ , because  $g_0$ is a product of eigenvalues of *U*. By assumption, there is another nonzero coefficient  $g_i$ , which proves the result.  $\blacksquare$ 

## **V. EXTENSIONS**

We describe in this section one possibility to extend our approach to a larger class of unitary matrices *U*. We assumed so far that  $f(U)$  is realized by a linear combination  $(2)$  of *linearly independent* matrices *U<sup>i</sup>* . The exponents were restricted to the range  $0 \le i \le m$ , where *m* is degree of the minimal polynomial of *U*. We can circumvent the problem indicated in the preceding section by allowing *m* to be larger than the degree of the minimal polynomial.

*Theorem 4.* Let  $U \in \mathcal{U}(2^n)$  be a unitary matrix such that *U<sup>m</sup>* is a scalar matrix for some positive integer *m*, i.e., the quotient of any two eigenvalues of *U* is a root of unity. Suppose that there exists a quantum circuit which implements *U* with *K* elementary gates. Then a unitary matrix *V*  $f(U)$  can be realized with  $O(mK+m^2\ln m)$  elementary operations.

*Proof.* By assumption,  $U^m = \tau \mathbf{1}$  for some  $\tau \in \mathbb{C}$ . This means that the minimal polynomial  $m(x)$  of *U* divides the polynomial  $x^m - \tau$ , that is,  $x^m - \tau = m(x)m_2(x)$  for some  $m_2(x) \in \mathbb{C}[x]$ .

We may assume without loss of generality that the function *f* is defined at all roots of  $x^m - \tau$ . Indeed, we can replace *f* by an interpolation polynomial *g* satisfying  $f(U) = g(U)$  if this is necessary.

Choose any unitary matrix  $A \in U(2^n)$  with minimal polynomial  $m_2(x)$ . The minimal polynomial of the block diagonal matrix  $U_A = \text{diag}(U, A)$  is  $x^m - \tau$ , the least common multiple of the polynomials  $m(x)$  and  $m_2(x)$ . Express  $f(U_A)$  by powers of the block diagonal matrix  $U_A$ :

$$
f(U_A) = \text{diag}(f(U), f(A)) = \sum_{i=0}^{m-1} \alpha_i \text{diag}(U^i, A^i). \quad (12)
$$

The approach detailed in Sec. III yields a unitary matrix *M* to realize this linear combination. On the other hand, we obtain from Eq.  $(12)$  the relation

$$
f(U) = \sum_{i=0}^{m-1} \alpha_i U^i
$$

by ignoring the auxiliary matrices  $A^i$ ,  $0 \le i \le m$ . It is clear that a circuit of the type shown in Fig. 2 with  $\mu$  chosen such that  $2^{\mu-1}$  *m*  $\leq$  2<sup> $\mu$ </sup> implements this linear combination of the matrices  $U^i$ ,  $0 \le i \le m$ , provided we use the matrix *M* constructed above.

## **VI. CONCLUSIONS**

Few methods are currently known that facilitate the engineering of quantum algorithms. Linear algebra allowed us to derive efficient quantum circuits for  $f(U)$ , given an efficient quantum circuit for *U*, as long as *U<sup>m</sup>* is a scalar matrix for some small integer *m*. This method can be used in conjuction with the Fourier sampling techniques by Shor  $[8]$ , the eigenvalue estimation technique by Kitaev  $[9]$ , and the probability amplitude amplification method by Grover  $[10]$ , to design more elaborate quantum algorithms.

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