

## Quantum search without entanglement

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Entanglement of quantum variables is usually thought to be a prerequisite for obtaining quantum speedups of information processing tasks such as searching databases. This paper presents methods for quantum search that give a speedup over classical methods, but that do not require entanglement. These methods rely instead on interference to provide a speedup. Search without entanglement comes at a cost: although they outperform analogous classical devices, the quantum devices that perform the search are not universal quantum computers and require exponentially greater overhead than a quantum computer that operates using entanglement. Quantum search without entanglement is compared to classical search using waves.

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Quantum computers exploit quantum coherence to perform computations in ways that classical computers cannot [1–5]. Despite the considerable difficulties involved constructing quantum computers [6,7] simple quantum logic devices have been built and prototype quantum computations have been performed [8–16]. Quantum computation is known to be able to solve some problems more rapidly than is possible classically [17–24]. Some problems, such as factoring and quantum simulation, can apparently be solved exponentially faster on a quantum computer than on a conventional digital computer [19–21]. Other problems, such as database search [22–24], can be solved polynomially faster on a quantum computer. The goal of this paper is to clarify what aspects of quantum mechanics are responsible for these speedups. In particular, it is often claimed that quantum speedups arise out of the quantum phenomenon known as entanglement [25]. This paper shows that this claim, while accurate by and large, is incomplete: although digital quantum computers that operate on quantum bits or qubits typically exhibit entanglement in the course of computation, when operating on only a few qubits, they can obtain a speedup over the best classical device without becoming entangled. In addition, it is possible to obtain quantum speedups using special-purpose devices that do not exhibit entanglement. Grover's algorithm for database search, for example, searches a database with  $n$  slots using only  $O(\sqrt{n})$  queries, while the best classical algorithm requires  $O(n)$  queries. Although Grover's algorithm as originally formulated induces entanglement in the qubits of a quantum computer performing the algorithm, Farhi and Gutmann's work [24] on continuous analogs of Grover's algorithm do not require entanglement; this paper builds on the work of Farhi and Gutmann to show that it is possible to construct quantum search devices that also give  $\sqrt{n}$  speedup over classical devices, but that do not require entanglement. These devices rely not on entanglement to obtain their quantum speedup, but on interference. Such devices are not general purpose quantum computers: to perform quantum searches without entanglement, they incur an exponentially greater overhead in their incidental operations than a universal quantum com-

puter that operates using entanglement. They nonetheless provide a speedup over classical devices. Finally, the paper shows how it is possible to construct classical search devices using waves that provide a  $\sqrt{n}$  speedup over the best classical search device that uses particles.

Entanglement is a peculiarly quantum phenomenon that is responsible for a variety of counterintuitive effects such as apparent quantum nonlocality, quantum teleportation, etc. [26,27]. A pure state  $|\psi\rangle$  for a quantum system composed of two or more subsystems is said to be entangled if it cannot be written in tensor product form:  $|\psi\rangle \neq |\psi_1\rangle \otimes |\psi_2\rangle \otimes \dots$ . Note that entanglement is not a property of the state  $|\psi\rangle$  on its own, but rather of the state and the way in which the system is divided up into subsystems. The claim that a quantum computation requires entanglement relies on a division of the quantum computer into quantum bits or qubits.

In Grover's algorithm for database search, a single item located in one of  $n$  slots in a database is located with only  $O(\sqrt{n})$  queries of the database [22]. This clearly gives a speedup over classical database search, in which  $n-1$  queries are required in the worst case and  $n/2$  queries are required on average. Grover's algorithm is normally taken to involve entanglement. In Grover's original version of this algorithm, he took  $n=2^r$  and performed the specified operations using quantum logic on  $r$  qubits. For  $r>2$  these operations entail putting the qubits in an entangled state at some point in the operation. However, as will now be shown, this entanglement is not an essential for obtaining a speedup over a classical device, but rather a by-product of the mapping of the steps above onto qubits. The following implementation allows one to perform quantum search in a way that does not require entanglement, but that nonetheless does better than the best classical device.

Consider a box with  $n$  slots through which a coin can be dropped. In all but one of the slots, when the coin goes in heads it comes out heads: the slot does not flip the coin. In the remaining slot, when a coin goes in heads, it comes out tails: the slot flips the coin. The problem is to find which slot flips the coin. One way to find out is to take the box apart and look to see which slot has a twist: but let us suppose that one is only allowed to put coins in and see if they come out flipped or not. In this case, one has to put  $n/2$  coins through on average and  $n-1$  in the worst case to locate the slot that

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flips the coin. If there is a meter on the box that charges a dollar for each coin that goes through, searching a box with 100 slots costs \$50, on average. This problem is clearly a version of the database search problem. It differs from Grover's original formulation only in that Grover assumed the slots in his database to be labeled by binary numbers, whereas in the "box" version no particular labeling need be specified.

Now look at a quantum-mechanical version of this problem. Use a quantum particle such as a neutron as a quantum "coin:" the box is constructed so that when a neutron with polarization  $\uparrow$  goes through all but one of the slots it emerges with polarization  $\uparrow$ . But when it goes through the remaining slot, it is flipped and emerges with polarization  $\downarrow$ . Equivalently, when it goes through that slot with a polarization  $\rightarrow = (1/\sqrt{2})(\uparrow - \downarrow)$ , the neutron acquires a phase of  $-1$ . For example, the slot that does the flipping could contain a magnetic field along the  $\rightarrow$  axis that flips the spin about that axis. Let  $|l\rangle$  be the state in which the neutron is in the mode that goes through the  $l$ th slot. Let  $u$  be the label of the slot that flips the neutron ( $u$  for "unknown"). Using neutrons with polarization  $\rightarrow$ , the effect of the box is to take the incoming state  $|l\rangle \rightarrow (1 - 2\delta_{lu})|l\rangle \equiv \mathcal{O}|l\rangle$ .  $\mathcal{O} = e^{-i\pi|u\rangle\langle u|}$  is the unitary operator that gives the effect of the box on the neutron.

Now pose the question: "How many times must one put a neutron or quantum coin through the quantum box to figure out which slot flips the neutron?" The answer is  $O(\sqrt{n})$  times as the following procedure shows. We have  $n$  translational modes of the neutron, one going through each of the  $n$  slots in the box. Let  $\mathcal{B}$  be the unitary operator representing the action of a beam splitter that takes a neutron from one mode and divides it equally among all the modes: such a beam splitter can be constructed from  $O(n)$  two-mode beam splitters. Let  $\mathcal{B}^\dagger$  be the unitary operator corresponding to the "inverse" beam splitter that undoes the action of the first. Finally, let  $\mathcal{I} = -e^{-i\pi|1\rangle\langle 1|}$  be the unitary operator that corresponds to an inverter that gives every mode except for the first a phase of  $-1$ . The inverter could also be constructed from a magnetic field.

The procedure for finding the unknown slot is as follows. Take a neutron in mode 1 and put it in sequence through the beam splitter, then the box, then the inverse beam splitter, then the inverter. The net effect is to apply the operator  $\mathcal{I}\mathcal{B}^\dagger\mathcal{O}\mathcal{B}$  to the initial state  $|1\rangle$ . By comparison with Grover's original algorithm, it is easily seen that  $\mathcal{B}$  gives an action analogous to a Hadamard transformation on the original state,  $\mathcal{O}$  gives the same action as Grover's "quantum phase oracle" that gives the effect of querying the database,  $\mathcal{B}^\dagger$  gives an action analogous to an inverse Hadamard transformation, and  $\mathcal{I}$  gives the same action as the operation Grover calls "inversion about average." Now take the neutron and put it through the beam splitter, box, inverse beam splitter, and inverter again, and again,  $O(\sqrt{n})$  times. By the same calculation as in Grover's original algorithm, the neutron is now with high probability emerging from the  $u$ th slot, and detection of its position will reveal  $u$ . The location of the slot that flips the quantum coin has been determined by putting a neutron through the box only  $\sqrt{n}$  times. If the box has a

meter that charges a dollar each time a neutron goes through, searching a box with 100 slots costs only \$10. The best classical strategy using coins, by comparison, costs on average \$50: the quantum version does better. There is no entanglement as there is only one neutron, and nothing for it to be entangled with.

The key to seeing that entanglement is not required for quantum search is to note that there is nothing in Grover's description of his search algorithm that requires the  $n$ -state system to be composed of qubits. The  $n$  states could just as well be discrete states of a single quantum variable. In such a "unary" representation there is no entanglement as there is only one quantum variable. There is nothing to be entangled with. There is nothing in either the classical or the quantum search problem that requires that the problem be formulated in a "binary" representation in which the slots in the database are labeled by binary numbers. Indeed, the unary representation of the search problem is more "natural" in that the slots are labeled by natural numbers without requiring that a particular base (2 or 10, e.g.) be specified. Even if one demands that a base be specified, then as long as the base is greater than  $n$ , no entanglement is required. In fact, the point that entanglement is not required in few-qubit quantum algorithms has been noted before [28], but with the misleading conclusion that the algorithms are not quantum-mechanical because entanglement is not present. Clearly, the unary representation still gives a  $\sqrt{n}$  speedup over the classical search problem in the sense that in the quantum version the quantum "coin" need only be passed through the box  $\sqrt{n}$  times.

The use of a unary representation does not come without cost, however. The conventional binary version of Grover's algorithm uses  $O(\log_2 n)$  qubits and requires  $O(\log_2 n)$  operations to perform each inversion about average and to determine the final result. The unary version of Grover's algorithm, in contrast, although it requires only  $O(1)$  operation to perform the inversion about average (all that is required is a single phase delay on the first mode) requires  $O(n)$  two-mode beam splitters to manipulate the neutron and  $O(n)$  detectors to read out the result. Although both devices give a  $\sqrt{n}$  speedup over the analogous classical device in the sense that they have to consult the "database" fewer times, the unary version requires exponentially more resources than the qubit version to perform the incidental operations. Like unary optical simulators of quantum logic [29] such devices are emphatically not universal quantum computers. The number of resources required to simulate an  $N$  qubit quantum computation using such a unary representation goes as  $2^N$ . Accordingly, unary devices cannot provide an exponential speedup over classical devices: the best they can do is reduce the number of times that they consult the database or "oracle." To map the operation of Shor's algorithm to a unary device, for example, would require exponential resources.

Before turning to classical search using waves, a further discussion of entanglement is in order. As noted above, entanglement is not a property of the state of a system on its own, but rather of the state and the way in which one divides the system up into subsystems. By changing the way one divides up the system, it is always possible to represent an

unentangled state as an entangled state and *vice versa*. For example, if one describes the single neutron interferometric database search method in a “second quantized” picture, in which the state  $|0\rangle_1 \cdots |1\rangle_l \cdots |0\rangle_n$  represents a state in which the neutron is in the  $l$ th mode, then the initial state  $(1/\sqrt{n})(|1\rangle_1 \cdots |0\rangle_n + \cdots + |0\rangle_1 \cdots |1\rangle_n)$  exhibits entanglement between the  $n$  modes. We now present two further unary versions of quantum in which such a second quantized picture is less applicable.

First, the  $n$  states could be different energy levels of a single atom. In this case, the action of  $\mathcal{B}$  above could be accomplished by a shaped, broadband pulse that takes the atom from the state  $|1\rangle$  to an equal superposition of the first  $n$  energy levels; the box could effect the phase inversion  $\mathcal{O}$  of the unknown state  $|u\rangle$  by driving a  $2\pi$  pulse between  $|u\rangle$  and the ground state  $|0\rangle$  (recall that we are not allowed to look inside the box and determine  $u$  by detecting this pulse); the action of  $B^\dagger$  could be accomplished by a shaped, broadband pulse that inverts  $B$ ; and the inversion about average  $\mathcal{I}$  could be accomplished by driving a  $2\pi$  pulse between  $|1\rangle$  and  $|0\rangle$  as for  $\mathcal{O}$  followed by a broadband  $2\pi$  pulse between all the states and the ground state. After  $O(\sqrt{n})$  iterations of the operations  $\mathcal{I}\mathcal{B}^\dagger\mathcal{O}\mathcal{B}$ , the system is in the state  $|u\rangle$  and a measurement of its energy will reveal the position in the database. This measurement could be performed, for example, by interchanging each state  $|x\rangle$  with the ground state in turn, and by driving a cycling transition that induces fluorescence if and only if the system is in the ground state. Although such a measurement requires up to  $n$  steps, it does not require any further passages through the box.

A second example of database search without entanglement is given by the Farhi-Gutmann continuous version of Grover’s algorithm [24]. Here, the system could be a spin with  $n$  states, and the database is given by a box that applies a Hamiltonian  $|u\rangle\langle u|$ . You are allowed to prepare the spin of any desired state, and to add your own preferred Hamiltonian to the unknown Hamiltonian applied by the box. Farhi and Gutmann show that if you start the spin in an *arbitrary* state  $|\psi\rangle$  and add to the database Hamiltonian the Hamiltonian  $|\psi\rangle\langle\psi|$ , then after a period of time proportional to  $\sqrt{n}$  the spin has rotated to the state  $|u\rangle$  with high probability.

One might ask whether by discarding the qubit representation one might be able to improve on the  $\sqrt{n}$  speedup that Grover’s algorithm gives over the classical bound. Grover’s algorithm is known to be optimal for qubits [30] and the proof of its optimality does not intrinsically rely on the qubit representation. Accordingly, the unary representations given here cannot provide any further speedup.

Once it is clear that a representation in terms of qubits is not required for quantum search, many implementations are possible. In fact, as will now be seen, a *classical* implementation phrased in terms of waves can still give a  $\sqrt{n}$  speedup over the classical search problem phrased in terms of coins or particles. The set of search methods has now come full circle: it was by considering interference via classical waves emitted by an array of antennae that Grover arrived at his quantum algorithm in the first place [31].

Return now to the neutron interferometer picture of quantum search described above. What in this picture is quantum-mechanical? There are in fact only two points in which the quantum nature of the neutron appears. The first is in the billing procedure: the box charges on a per quantum basis. The second is the final click of the detector at the slot from which the neutron emerges. These are the only points at which the neutron is required to behave like a particle. At all other points in the search process, it is the wave aspect of the neutron that comes into play: it is the interference between the waves in the interferometer that lies behind the  $\sqrt{n}$  speedup.

This dominance of the wave aspect of quantum mechanics suggests the following purely classical wave method for search. Instead of quantum matter waves, use classical waves such as light or sound. At bottom, of course, such waves are composed of photons and phonons. But it is possible to reformulate the search problem in such a way that the particle aspect of the waves is unimportant. Let the unknown slot in the box flip the polarization of the waves, and suppose now that the box charges on the basis of the integrated intensity of the waves that pass through it rather than on a per particle basis. We are provided with detectors with a finite signal-to-noise ratio. How now does the cost of determining the unknown slot scale with the number of slots  $n$ ?

One way to search the box is simply to shine waves through all the slots at once and to determine which slot flips the polarization of the transmitted wave. Because of the finite signal-to-noise ratio of the detectors, the cost of this method is proportional to  $n$ . A second method is to recycle the waves through an interferometer constructed in exact analog to the neutron interferometer described above to give positive interference at the output of the unknown slot. Just as in the quantum case, the cost of this method is proportional to  $\sqrt{n}$ . So a purely classical wave search device can also find the unknown slot with an integrated intensity proportional to  $\sqrt{n}$ . Of course, if one tries to minimize costs by decreasing the intensity of the recycled waves and increasing the sensitivity of the detectors, one’s data will eventually arrive in the form of individual “clicks”: the quantum nature of the wave will reassert itself.

The interferometric versions of database search described above can be thought of as complementary to the well-known phenomenon of an interaction-free measurement [32–34]. In the most dramatic version of such an experiment, one wishes to use optical methods to detect the presence of a bomb that explodes if it absorbs a photon. By comparison, one can phrase quantum database search as a problem of finding and exploding a similar light-sensitive bomb in one of  $n$  slots while only firing photons at it  $1/\sqrt{n}$  times. Both bomb detection and bomb demolition are easier with quantum resources. It is important to note, however, that just as in the neutron interferometer search method above, what makes bomb detection and bomb demolition quantum mechanical is not the factors of  $\sqrt{n}$ , but rather the fact that the bomb explodes when it absorbs a single quantum of light. A “classical” bomb that explodes when it has been subjected to a



given integrated intensity of classical light can be detected in a nondemolition fashion using only a classical interferometer.

To summarize:

(i) A classical digital computer that searches a database with  $n$  slots requires  $O(\log_2 n)$  resources and has to look at the database  $O(n)$  times.

(ii) A quantum digital computer that searches a database with  $n$  slots requires  $O(\log_2 n)$  resources and has to look at the database  $O(\sqrt{n})$  times.

(iii) A classical device that determines which of  $n$  slots in a box flips a discrete object, such as a coin, requires  $O(n)$  resources and has to pass the coin through  $O(n)$  times.

(iv) A quantum device that determines which of  $n$  slots in a box flips a discrete object, such as a particle, requires  $O(n)$  resources and has to pass the particle through  $O(\sqrt{n})$  times.

(v) A classical wave device that determines which of  $n$  slots in a box flips the polarization of a wave requires  $O(n)$

resources and has to send the wave through  $O(\sqrt{n})$  times.

Special purpose quantum search devices can give a speedup over classical search devices without using entanglement. A quantum device that probes a system by sending discrete objects such as particles through it can acquire information about unknown features of the system more rapidly than analogous classical devices that probe a system by sending discrete objects through it. The  $\sqrt{n}$  speedup obtained by the quantum devices arises out of the wave nature of the particles sent through. Classical devices that rely on waves and interference can also give a  $\sqrt{n}$  speedup over classical devices that probe a system using particles alone.

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