¹D. R. Johnson, G. J. Oudemans, and R. H. Cole, J. Chem. Phys. <u>33</u>, 1310 (1960).

 2 R. H. Orcutt and R. H. Cole, J. Chem. Phys. <u>46</u>, 697 (1967).

³A. Dalgarno and A. E. Kingston, Proc. Roy. Soc. (London) <u>A259</u>, 424 (1960).

 4 M. N. Grasso, K. T. Chung, and R. P. Hurst, Phys. Rev. <u>167</u>, 1 (1968).

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ERRATA

Laser-Induced Line-Narrowing Effects in Coupled Doppler-Broadened Transitions. II. Standing-Wave Features, B. J. Feldman and M. S. Feld [Phys. Rev. A 5, 899 (1972)]. The sentence on p. 908 which begins 11 lines from the bottom of the second column has been garbled. It should read: "Accordingly, when E_2 is detuned, the $\langle R_{02} \rangle$ curve is simply the superposition of the broad and narrow traveling-wave resonances, well resolved from each other and symmetrically located about ω_1 [Fig. 6(a)]."

S-Matrix Formulation of Statistical Mechanics,

R. Dashen, S. Ma, and H. J. Bernstein [Phys. Rev. <u>187</u>, 345 (1969)]. (a) Eq. (3.28), $V(G_0 - G_0^{\dagger})V$ should be $V(G - G^{\dagger})V$. (b) The first sign in Eq. (3.41) should be +, and the proof is therefore wrong. The right proof is the following. Integrating by parts (the first term over E'' and the second over E') we obtain

$$R = \delta' (E - E'') [T_{EE} \delta(E - E') T_{EE'}^{\dagger},] - \delta' (E - E') [T_{EE}^{\dagger} \delta(E - E'') T_{EE'}].$$

Summing over E', E'' and other variables not explicitly shown and applying the unitarity condition, we obtain

$$\sum_{E'E''} R = -\sum_{E''} \frac{\delta'(E - E'')(T_{EE''} - T^{\dagger}_{EE''})}{2\pi i} + \sum_{E'} \frac{\delta'(E - E')(T_{EE'} - T^{\dagger}_{EE'})}{2\pi i} = 0. \quad Q. E. D.$$

(c) In the seventh line above Eq. (4.3), "...which leave none of the elements...," should read "... which leave none of the subsets of the elements ...". (d) Equation (5.10) should be $\int_0^{\infty} d\epsilon \ (4\pi i)^{-1} \times e^{-\beta\epsilon} \dots$

For more general discussions on the subject of Sec. VII D, see R. Dashen and S. Ma [J. Math. Phys. 11, 1136 (1970); 12, 689 (1971)]. Extension of the theory to correlation functions can be found in R. Dashen and S. Ma [J. Math. Phys. <u>12</u>, 1449 (1971); Phys. Rev. A <u>4</u>, 700 (1971)].

⁵P. Langhoff, J. D. Lyons, and R. P. Hurst, Phys.

 6 J. Heinriches, J. Chem. Phys. <u>52</u>, 6316 (1970). ⁷B. P. Tripathi, R. K. Laloraya, and S. L. Srivastava,

⁹J. C. Slater, *Quantum Theory of Atomic Structure* (McGraw-Hill, New York, 1960), Vol. I, Appendix 11.

⁸J. P. Vinti, Phys. Rev. <u>42</u>, 632 (1932).

Atomic *M*-Shell Coster-Kronig, Auger, and Radiative Rates and Fluorescence Yields for Ca-Th, Eugene J. McGuire [Phys. Rev. A 5, 1043 (1972)]. Because of an error the fluorescence yields ω_{M4} for $32 \le Z \le 57$ in Table V are incorrect. The corrected table reads:

Ζ	$\omega_{_{M4}}$
32	1.80 E-5
36	1.49 E-4
40	4.35 E-4
44	2.62 E-4
47	2.56 E-4
50	3.19 E-4
54	4.39 E-4
57	6.88 E-4
60	2.6 E-3

I wish to thank Professor Chandra Bhalla for pointing out this error.

Collision Times and Kinetic Theory for Superfluid Helium, Wayne M. Saslow [Phys. Rev. A 5, 1491 (1972)]. On p. 1491, the second and third sentences on the second paragraph should read: "In particular, measurement of first- and secondsound attenuation can give the longitudinal relaxation time τ_i associated with the second viscosity ξ_2 . Further, if measurements can be made on shear waves in the hydrodynamic regime, they will provide a test for the calculation by Landau and Khalatnikov of τ_t , the transverse relaxation time.¹⁵"

On p. 1497, the sentence beginning eight lines from the bottom of the first column should read: "This lends support to the use of collision-time models at low frequencies, but it is not clear that errors may not arise at higher frequencies, where the works of Disatnik and KC differ."