ERRATA

P Autoionization States of He and H^-, A. K. Bhatia and A. Temkin [Phys. Rev. 182, 15 (1969)]. The value of the reduced rydberg in the above reference for H^- (not for He) was miscopied and the incorrect value used for H^- results in Tables II and V. The changes are small but they are correctly: The effective rydberg for $H^- = 13.597940$ eV and $\mathcal{S} = 9.720$ eV for $H^-(^3P)$ in Tables II and V. Finally, the values of Δ and Γ remain unchanged to the accuracy given, thus E for $H^-(^3P)$ in Table V should be 9.733 eV.

Inelastic Neutron Scattering from bcc ^{4}He, E. B. Osgood, V. J. Minkiewicz, T. A. Kitchens, and G. Shirane [Phys. Rev. A 5, 1537 (1972)]. Many α^*'s were put in place of a^*'s. Also, the first paragraph, beginning in the right-hand column of p. 1541 in our discussion of the anomalous scattering results, contains typesetting errors which have bollixed the point of discussion. It is presented below correctly.

For a direct comparison of the shapes and the expected intensities, the data taken at (1.5, 0, 0) are plotted with the background subtracted in Fig. 6. The Gaussian fit to the (0.5, 0, 0) data is also plotted on the same figure. The amplitude is corrected for the difference in counting time and for the factor $\vec{Q}^2 e^{-2W}$. For $|\vec{Q}| = 0.5a^*$, the profile is symmetric and for $|\vec{Q}| = 1.5a^*$ there is more intensity than expected, as is seen in the figure. The profile with extra intensity is asymmetric with a sharp rise at low-energy transfer, and peaks around an energy of 1.4 meV, and thereafter slowly falls off at higher energies. This is also characteristic of most profiles with extra intensity. The same features are presented for phonon profiles at $|\vec{Q}| \sim 1.6a^*$ along other symmetry directions, such as the $L[111]$ and $T[100]$. Outside the shaded area, groups were observed with distorted profiles but without additional intensity.

Adsorption of Helium on Rare-Gas-Plated Graphite, Anthony D. Novaco and Frederick J. Milford [Phys. Rev. A 5, 783 (1972)]. The contribution to the total potential due to the graphite substrate is in error by a factor of 2. In Eq. (1), the $4\pi\epsilon_0'$ should be replaced by $8\pi\epsilon_0'$. This correction results, for $2.0 < z < 5$ Å, in an almost rigid downward shift of about 10% in the potential profiles. The only effect is to lower the ground-state energy by the same amount, while keeping intact the densities of states relative to the ground-state energy. The distribution functions $P_1(x, y)$ and $P_2(z)$, along with the specific heat $C(T)$, remain unaltered. The correct ground-state energies are

He^4 on argon-plated graphite: $-71\,^\circ$K,

He^3 on argon-plated graphite: $-66\,^\circ$K,

He^4 on xenon-plated graphite: $-77\,^\circ$K,

He^3 on xenon-plated graphite: $-72\,^\circ$K.

Collision Spectroscopy. IV. Semiclassical Theory of Inelastic Scattering with Application to $He^+ + Ne$, R. E. Olson and F. T. Smith [Phys. Rev. A 3, 1607 (1971)]. The recent calculations by E. E. Nikitin and A. I. Reznikov [this issue, Phys. Rev. A 6, 522 (1972)] and by A. M. Woolley (private communication) have prompted us to redo the Landau-Zener-Stueckelberg transition probabilities shown in our Fig. 4(a). We have found that our Gaussian quadrature used by us was not fully converged and required a smaller grid step. When this was done the phases shifted by 0.28π radians at $l = 0$ and correspondingly less for larger l values. We now are in complete agreement with Nikitin and Reznikov and with Woolley. Our Fig. 4(a) will now be the same as the Nikitin and Reznikov's figure when $\gamma_l = \frac{1}{4}\pi$. A typographical error also exists in Eq. (42) and $(4U_{12})^2$ should read $4U_{12}{}^2$. All other results will be unchanged.

We would like to thank Dr. A. M. Woolley of the University of Copenhagen and Dr. Nikitin for communicating their results to us.