Directed alternating lattices and the site-to-bond ratio for animals and trees

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(Received 16 December 1991)

We present here a brief summary of results on so-called "exotic" directed lattices having nonregular periodicity. Such lattices, e.g., the Archimedean nets, are characterized by different site types and a spread of coordination numbers. The evidence adduced here shows that the exponent structure for the growth of animals and trees on such lattices is of the form predicted, with θ_0 for trees equivalent to θ for unrestricted animals, independent of the periodicity property. This supports the general form $\theta_c = \theta_0 - c$, for cycles $c = 0$. We quote values for the growth parameter (or inverse critical fugacity) λ for bond trees and animal growth on selected directed lattices. The convergence of such series is well known to be subject to the influence of subdominant singularities in addition to θ , and we report on results obtained using the second-log-derivative scheme for the lattices of interest. Recent results from percolation studies for the alternating nets [H. J. Ruskin, Phys. Lett. A 162, 215 (1992)] have proved particularly encouraging for the Archimedean lattices, with $\overline{z}=2.498+-0.010$, but uncertainties for the honeycomb were found to be large. Investigation of the site-to-bond ratios for animals and trees on the honeycomb gives somewhat smoother series behavior, which, though subject to confluence effects, supports a much lower value of the effective coordination number. We quote $\overline{z}=1.450+-0.050$.

PACS number(s): $05.50.+q$

INTRODUCTION

The statistics of finite cluster growth on regular lattices has proved a useful study in terms of the understanding of various biological growth processes ([1] and references cited therein) for dilute polymer problems [2], and more recently in the formulation and understanding of various lattice gauge theories [3]. For isotropic lattices, trees and full animals have been shown to have the same critical exponents in all dimensions [4,5] and their hierarchical structure has been explored [6,7]. It was early noted [8] that on general types of lattice all site trees are neighboravoiding walks, and hence no proper zero-cycle configurations are possible. This phenomenon is observed for all lattices which are bond-to-site transformations of other nets, with additional numbers of loops dependent on the valence partition. The establishment of

FIG. 1. (a) Directed Kagomé lattice showing different site types. (b) Directed Archimedean lattice showing different site types (different origins for cluster growth).

a similar hierarchical structure for directed animals [9] has meant that the exploitation of recent developments with respect to site-to-bond conversion for directed graphs [10,11] has provided the means of direct investigation of the cyclomatic number partition of directed bond configurations. Utilizing the information from site-tobond conversion, we explore here the singularity structure for directed animals and trees for the Archimedean (3,3,3,4,4) and honeycomb nets (bond}, which fall into the class of "exotic" lattices. These lattices are characterized by alternating periodicity with coordination number, z, having more than one value which is dependent on the site type. The connectivity features of the different site types thus provide different origins for cluster growth [illustrated in Fig. 1(b} for the Archimedean]. Results for the alternating nets are contrasted with those of the Kagomé lattice, which, despite regular periodicity (singlevalued z), shares several similar properties [Fig. 1(a), with site types distinguished by the number of incoming and outgoing bonds on the hexagon].

EXPONENT STRUCTURE AND SITE-TO-BOND RATIOS

The critical exponent structure which we investigate for these exotic lattices assumes a generating function of the form

$$
A_c(x) = 1 + \sum_{n=1} A_{n,c} x^n , \qquad (1)
$$

with

$$
A_{n,c} \sim n^{-\theta_c} \lambda_c^n = \sum_c A_n \lambda_0^n n^{-\theta_0 + c} \tag{2}
$$

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and

$$
\theta_c = \theta_0 - c \tag{3}
$$

where *n* is the size of the site (bond) cluster, λ is the growth parameter or inverse critical fugacity, and $A_{n,c}$ is the total number of directed animals with size n having c cycles (or loops). λ can also be regarded as tree multiplicity, where $c = 0$ for trees. Further, the presence of subdominant singularities is known to influence the convergence of animal series, so that we may extend Eq. (1) above $[3,12]$ to write for animals of loop size c

$$
A_{n,c} \sim \lambda_0^n n^{-\theta_c} (1 + B n^{-w} + \cdots) \tag{4}
$$

Thus by applying the second log-derivative scheme, we may obtain information on the effect of the confluences.

The investigation of a critical site-to-bond ratio for directed animals examines the form [10] for sizedependent growth by means of the bond perimeter distribution, i.e.,

$$
\lambda(a) = (a+1)^{a+1}/a^a - k(a)
$$
, $k(a) > 0$ for $a > a_c$, (5)

where

 \equiv

 \equiv

$$
a = t/b = \bar{z} - 1 + (\bar{z}/b) - (\bar{z}c/b) , \qquad (6)
$$

with \tilde{z} the "effective" coordination number, t the perimeter, b the number of bonds, and a_c the critical ratio. In a previous paper, Ruskin [13] has considered percolation

TABLE I. Honeycomb bond animals.

thresholds for selected alternating bond lattices, and has obtained estimates of \tilde{z} for these cases. For the Archimedean lattice, results were good, with $\tilde{z} \sim 2.5$, in agreement with the conjectured exact value based on averaging the number of sites of different types. Estimates were less satisfactory for the honeycomb lattice, given the characteristic oscillation of sign for the percolation series. However, there was some evidence to show that \tilde{z} could be as low as 1.5, which would support a conjecture based on similar configurational considerations. Furthermore, Eqs. (5) and (6) above clearly indicate that an investigation of trees (cycle or loop size $c = 0$) and animals with $c > 0$ should provide an alternative estimate of \tilde{z} . The hierarchical structure of exponents for these fixed-cycle clusters is expected to be of the form given in Eq. (2), with some support for this assumption provided by earlier studies [2,11,14]. Here we investigate series for both trees and animals for selected directed nets with regular and alternating periodicity.

RESULTS

Series data on the total number of directed animals are given by Ruskin [14] for the Archimedean (3,3,3,4,4) to order 14 (15) for the bond (site), and for the Kagome bond and site to order 22. The honeycomb bond animals are given in Table I, and the bond trees for Archimedean, Kagome, and honeycomb lattices are given in Table II. Relevant site-to-bond ratios, denoted $\langle t \rangle/b$ and $\langle s \rangle/b$,

TABLE II. Bond trees.

	TABLE 1. Honeycomb bond animals.				
	$\boldsymbol{2}$		Honeycomb	Kagomé	Archimedean
\overline{c}	$\overline{\mathbf{3}}$		$\overline{\mathbf{c}}$	4	5
3	6	$\overline{2}$	3	10	16
4	11	3	6	27	56
5	22	4	11	76	206
6	44	5	22	221	781
	89	6	43	657	3019
8	182	$\overline{7}$	86	1982	11831
9	374	8	173	6039	46837
10	775	9	350	18543	186876
11	1611	10	712	57305	750267
12	3363	11	1452	178049	3027441
13	7042	12	2975	555 650	12 267 584
14	14790	13	6108	1740 394	49 885 928
15	31139	14	12577	5468 171	203 473 454
16	65 697	15	25 9 44	17 226 909	
17	138869	16	53 626	54 399 777	
18	294 029	17	111020	172 142 129	
19	623492	18	230189	545 724 610	
20	1323898	19	477902	1732 885 298	
21	2814533	20	993 373	5510656445	
22	5990156	21	2067084	17 547 367 745	
23	12761699	22	4305 588	55 942 528 625	
24	27 213 259	23	8976406		
25	58079355	24	18729822		
26	124 051 993	25	39 111 032		
27	265 155 766	26	81728463		
28	567 141 107	27	170 896 956		
29	1213817958	28	357 570 769		
30	2599 382 044	29	748 580 760		

	Kagomé Animals $\langle s \rangle/b$	Honeycomb			
			Animals		Trees
		$\langle t \rangle / b$	$\langle s \rangle/b$	$\langle t \rangle/b$	$\langle s \rangle/b$
4	1.4634	0.8636		0.8636	
5	1.3579	0.7818		0.7818	
6	1.2870	0.7159	1.1629	0.7209	1.1667
7	1.2357	0.6774	1.1380	0.6844	1.1429
8	1.1966	0.6435	1.1188	0.6517	1.1250
9	1.1658	0.6197	1.1040	0.6298	1.1111
10	1.1410	0.5994	1.0919	0.6104	1.1000
11	1.1207	0.5829	1.0818	0.5955	1.0909
12	1.1037	0.5692	1.0735	0.5827	1.0833
13	1.0892	0.5576	1.0664	0.5720	1.0769
14	1.0768	0.5476	1.0603	0.5629	1.0714
15	1.0660	0.5389	1.0550	0.5549	1.0667
16	1.0565	0.5313	1.0503	0.5480	1.0625
17	1.0481	0.5245	1.0462	0.5418	1.0588
18	1.0406	0.5185	1.0425	0.5364	1.0556
19	1.0339	0.5132	1.0392	0.5316	1.0526
20	1.0278	0.5083	1.0362	0.5272	1.0500
21	1.0223	0.5039	1.0335	0.5232	1.0476
22	1.0173	0.4999	1.0310	0.5197	1.0455
23		0.4963	1.0288	0.5164	1.0435
24		0.4929	1.0267	0.5134	1.0417
25		0.4898	1.0248	0.5106	1.0400
26		0.4870	1.0230	0.5081	1.0385
27		0.4843	1.0214	0.5057	1.0370
28		0.4819	1.0199	0.5035	1.0357
29		0.4796	1.0184	0.5015	1.0345
Central					
limiting value	0.908	0.414	0.978	0.447	0.995

TABLE III. Bond averages.

are given in [13] for the Archimedean bond lattice, and Table III gives the appropriate ratios for the Kagomé and honeycomb lattices, together with the central value obtained from asymptotic extrapolation. For the Kagomé lattice (with $\tilde{z}=2$), $\langle s \rangle/b$ ratios are included for completeness but $\langle t \rangle/b$ averages are not given, since these follow directly [Eq. (6)] for \tilde{z} known and constant. For the honeycomb trees, $\langle s \rangle/b$ are counted according to the usual definition $s = b + 1$, and, as can be seen from Table III, closely parallel the animal sequences, giving slightly higher central values. Estimates of the growth parameter (or tree multiplicity) for the bond trees for all lattices and for the honeycomb bond animals are summarized in Table IV, together with central values for the dominant critical exponent (θ_0 and θ , respectively). These are seen to be in good agreement with the form predicted ([9] and references cited therein), with exponent estimates for trees well within the error bounds quoted by Ruskin [14] for the unrestricted animals. (The animal estimates for Archimedean and Kagomé lattices are repro-

TABLE IV. Summary of growth parameter and dominant exponent estimates for bond trees. Animal estimates are included in comparison.

		$\theta_0(\theta)$
Kagomé site animals	2.7000 ± 0.0025	0.505 ± 0.005
Kagomé bond animals	3.426 ± 0.006	0.515 ± 0.015
Kagomé bond trees	3.2660 ± 0.0005	0.500 ± 0.005
Archimedean site animals	3.416 ± 0.004	0.505 ± 0.010
Archimedean bond animals	4.477 ± 0.007	0.510 ± 0.020
Archimedean bond trees	4.2350 ± 0.0025	0.500 ± 0.005
Honeycomb bond animals	2.21875 ± 0.0005	0.500 ± 0.005
Honeycomb bond trees	2.1778 ± 0.0002	0.500 ± 0.001

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Lattice	Discrimination	\boldsymbol{A}	B	φ
Honeycomb	$\langle t \rangle/b$	0.41382	1.753 60	-0.97503
		0.413 64	1.74583	-0.97290
		0.41386	1.755 66	-0.97554
		0.41384	1.75481	-0.97531
		0.41388	1.75674	-0.97581
		0.41390	1.75806	-0.97614
		0.41394	1.76022	-0.97668
	$\langle s \rangle/b$	0.97651	1.032 19	-0.95129
		0.97652	1.03273	-0.95155
		0.97664	1.03744	$-0.953\,71$
		0.97668	1.03924	-0.95451
		0.97670	1.04023	-0.95495
		0.97675 0.97679	1.04252 1.04485	-0.95593 -0.95692
Kagomé	$\langle s \rangle/b$	0.90636	2.095 18	-0.95074
		0.90729	2.11548	-0.95644
		0.90751	2.121 10	-0.95796
		0.90740	2.11803	-0.95716
		0.90733	2.11585	-0.95661 -0.95723
		0.90741 0.90756	2.11834 2.12385	-0.95853
Archimedean	$\langle s \rangle/b$	0.94876	1.11839	-0.96837
(Sub 2)		0.944 29	1.09020	-0.94189
		0.94387	1.08703	-0.93918
		0.943 65	1.08507	-0.93765
		0.943 53	1.08369	-0.93666
		0.94376 0.94409	1.08671 1.09156	-0.93869 -0.94173
	$\langle t \rangle/b$	1.34903	2.87945	-0.92760
		1.35545	2.91647	-0.94117
		1.35715	2.92945	-0.94530
		1.35874	2.944 67	-0.94964
		1.35993	2.95838	-0.95322
		1.36031	2.963 52	-0.95446 -0.95453
		1.36032	2.96381	
(Sub 1)	$\langle s \rangle/b$	0.94149	1.033 67	-0.92845
		0.943 15	1.043 17	-0.93818
		0.943 60	1.046 61	-0.94125
		0.944 16	1.05179	-0.94540
		0.944 63	1.05718	-0.94935
		0.94482 0.94486	1.05985 1.06036	-0.95117 -0.95149
	$\langle t \rangle/b$	1.37680	2.36526	-1.00313
		1.36375	2.27302	-0.96311
		1.36156	2.255 11	-0.95589
		1.360 53	2.24503	-0.95216
		1.35987	2.23736	-0.94952
		1.36030 1.36099	2.243 22 2.25396	-0.95140 -0.95464
		0.444 50	1.66640	-1.00256
Honeycomb trees	$\langle t \rangle/b$	0.445 44	1.71862	-1.01709
		0.445 05	1.69132	-1.00971
		0.44520	1.700 10	-1.01206
		0.445 16	1.69785	-1.01147
		0.44509	1.69299	-1.01021
		0.445 14	1.696 62	-1.01117

TABLE V. Three point fits, A, B , and ϕ .

duced here for convenience.) As we might expect, the growth parameters for trees are lower for all lattices than those for the unrestricted animals. The asymptotic limit for central θ is subject to some curvature in all cases, although this is less marked for the longer series. Trees in all cases tend to underestimate the exponent value with respect to the exact result predicted by Dhar [2]. Biasing the sequences with respect to exact θ leads, however, to generally good agreement with respect to the limiting estimates of λ , and this is reflected in the uncertainties quoted. It would appear, therefore, on the evidence of the tree and animal data, that the property of nonregular periodicity does not modify the hierarchy of the exponent structure.

Given the known shortcomings of ratio procedures with respect to confluence singularities of the type typically found for animal and fixed-cycle series, we have also looked at the correction for the subdominant singularity as described in Eq. (4) together with the limiting behavior of the averages, $\langle t \rangle/b$. Results for the modified asymptotic form for the averages $A + Bb^{-\phi}$ (Eq. (6) and [15]) are given in Figs. 2(a) and 2(b) for ϕ , together with three point fits on A, B, and ϕ (Table V). Here, ϕ does not necessarily relate to the sum of $\theta_c + w$ [Eq. (4)] but depends on the amptitudes $A_{n,c}$. The numerical evidence appears good for a central estimate of $\phi=1.0$ (honeycomb animals and trees), although the Archimedean results (based on shorter sequences) would suggest a lower

value, $\phi \sim 0.955$. The Kagomé results fall fairly smoothly between those for the alternating lattices. An overall estimate, taking all sequences into account, supports $\phi = 0.985 \pm 0.025$. Results obtained from applying the second-log-derivative scheme directly to estimate w (Table VI and Fig. 3) afford rather poor determination of

the subdominant animal exponent. The irregularities of the sequences obtained are clearly seen from the table, with only the Kagomé lattice (site values included for illustration) showing some degree of stability and suggesting $w = 0.8 \pm 0.1$, which excludes the $1/n$ correction, suggesting that the leading one is indeed a scaling correction.

FIG. 2. (a) Three-point fit sequences for ϕ (modified exponent) honeycomb and Kagomé. \times , honeycomb animals $\langle t \rangle /b$; \circ , honeycomb animals $\langle s \rangle / b$; \bullet , honeycomb trees $\langle t \rangle / b$; $+$, Kagomé animals $\langle s \rangle / b$. (b) Three-point fit sequences for ϕ (modified exponent) Archimedean (3,3,3,4,4). Sublattice 1: \bullet , $\langle t \rangle / b$; \circ , $\langle s \rangle / b$. Sublattice 2: \bullet , $\langle t \rangle / b$; \triangle , $\langle s \rangle / b$.

FIG. 3. Kagomé sequences for w (second log-derivative scheme). \bullet , bond; \circ , site.

While the precision is not particularly good, it is comparable with that attained for studies on the Lee-Yang edge singularities and other allied problems [16], with w in broad agreement with subsequent estimates [3,11].

From the central limiting values of the ratios given in Table III, we find the effective coordination number for the directed honeycomb lattice to be $\tilde{z} = 1.450 + -0.050$, which is low compared with the percolation series results [13]. These were shown to support a central value as high as $\tilde{z} = 1.75$, but with the smoother earlier sequences favoring the lower bound, $\tilde{z} = 1.5$, the large uncertainties being ascribed to the characteristic honeycomb series oscillation of sign for larger n . Unfortunately, the confluence effects present in the animal data in this instance also inhibit a more precise interpretation, although the evidence again suggests that a value close to 1.5 (roughly equivalent to that found for percolation data) is more likely to be correct.

From the central limiting value for the average perimeter-to-bond ratio following Duarte $[11]$, we predict the asymptotic value of the growth parameter for the honeycomb lattice to be $\lambda = 2.362 + -0.007$. For the Archimedean lattice, the central limiting value for $\langle t \rangle/b$ quoted by Ruskin [13] predicts an asymptotic value of $\lambda = 5.002 + -0.008$. The higher asymptotic values for both honeycomb and Archimedean lattices are predictable, or course, since the region of interest is now $a > a_c$ [Eq. (5)].

CONCLUSION

Investigation of the dominant exponent behavior for trees and animals on selected "exotic" lattices with both

regular and alternating periodicity provides evidence in support of a hierarchical exponent structure, which is unaffected by the periodicity, with θ for unrestricted animals equivalent to θ_0 for trees. Estimates of the growth parameters for the trees on selected lattices and for the honeycomb animal data are also given, and an analysis of the site-to-bond ratios for the honeycomb series suggests that the "effective" coordination number governing the lattice behavior is lower than that originally conjectured by percolation analyses. We quote $\tilde{z} = 1.450 + -0.050$. The effect of confluent singularities on the exponent estimates for the exotic lattices is investigated for the animal and tree series, but while numerical evidence supports an exponent ϕ for the averages in the region of 1.0 (particularly for longer sequences), results of estimating the animal correction term w directly are not conclusive. The best estimates for w are obtained from the Kagomé lattice and broadly support those of previous works, with $w = 0.8 \pm 0.1$.

ACKNOWLEDGMENTS

Thanks are due to Professor D. Weaire of Trinity College, University of Dublin for first drawing "exotic" lattices to our attention and to Dr. J. Duarte for useful discussions and generous provision of copies of his work prior to publication.

APPENDIX: RECURSION RELATIONSHIP

In terms of the enumeration of clusters, the linkages of the different site types can be vitally important, and the establishment of recursion relationships for cluster

growth on the different sublattices can provide a substantial reduction of effort required. Taking the Kagomé lattice as an example, it is well known that the isotropic lattice has three sublattices defined by three different rootsite types. For the directed (anisotropic) Kagome lattice, however, there is redundancy in site types, with two of the three being equivalent, owing to preferred directions of linkage. Hence, the directed Kagomé lattice has two sublattices, dependent on the two distinct site types, which we can denote I and II. As the Kagomé lattice is built up from hexagons and triangles, we can distinguish between these, with type I on the triangle having two incoming bonds, and type II having one incoming and one outgoing bond. Consequently, type-I sites can be exactly counted for a tree configuration of given size s. Further, since the Kagomé trees are bond-to-site transformations from the honeycomb lattice (i.e., self-avoiding walks), we can establish a recursion relationship as follows. From any type-I site, we can progress to either a type-I or a type-II site but not both, since we are interested in selfavoiding walks (SAW's). Thus, the number of SAW's for a type-I site is given by

$$
N_s^1 = N_{s-1}^1 + N_{s-1}^{11}
$$
 (A1)

Further, since the type-II progression represents a branch

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from a type-I site en route to another type-I site, we have

$$
N_{s-1}^{\rm II}\!=\!N_{s-2}^{\rm I}\;\;,\;\;
$$

i.e.,

$$
N_s^{\rm I} = N_{s-1}^{\rm I} + N_{s-2}^{\rm I} \tag{A2}
$$

In other words, only one site type is relevant for enumeration purposes, with initial conditions given by enumeration purposes, with finitial conditions given by $N_1^1 = 1$, $N_2^1 = 2$. For type-II sites as defined above, trees are not equivalent to SAW's, since the root site has two outgoing bonds, i.e., two branches. Clearly, since we can regard type-II sites as intermediary points on the type-I sublattice, SAW's are one branch trees in the type-II sublattice, and we have directly for SAW's

$$
N_s^{\rm II} = 2N_{s-1}^{\rm I} \tag{A3}
$$

For trees with more than one branch, an exact relationship is not straightforward, because as branches grow there is the possibility of interaction leading to a configuration that is not a tree. Growth parameters on the two sublattices will be equivalent, but in this case also, we have from Eqs. (A 1) and (A2) the result that only type-I sites need to be enumerated and that $\lambda_s^{\text{II}} = \lambda_s^{\text{I}} \forall s > 3.$

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