*Work supported in part by the U. S. National Science Foundation and the R. A. Welch Foundation.

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[13]The background dielectric constant of a representative Li-NH_3 solution is calculated as follows. The density of NH_3 at 195 °K is 0.7343 g/cm^3. Assume that the 20-MPM Li-NH_3 solution can be represented by Li$(NH_3)_4$. The molar volume of $(NH_3)_4$ is 92.6 cm^3/mole. The density of a 20-MPM Li-NH_3 solution at 195 °K is 0.5098 g/cm^3, making the molar volume of Li$(NH_3)_4$ equal to 147 cm^3/mole. The optical dielectric constant of NH_3 is 1.9 and we will assume that the optical dielectric constant of the other constituents of Li-NH_3 solutions is 1.0. The fraction of volume occupied by $(NH_3)_4$ is $92.6/147 = \alpha$. The other constituents occupy a volume fraction $1-\alpha$; therefore, we assert $K = 1.9\alpha + (1-\alpha)\,1.0 = 1.57$, where K is the total background dielectric constant. Concentrations down to 8 MPM were studied and K scaled between 1.57 and 1.61.

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ERRATA

Fredholm Method. II. A Numerical Procedure for Inelastic Scattering, William P. Reinhardt [Phys. Rev. A 2, 1767 (1970)]. Equation (2.14a) should read

$$\int_0^\infty \frac{k^\alpha dk^\alpha f(k^\alpha)}{\frac{1}{2}(k_0^\alpha)^2 + i\epsilon\,\mathrm{sgn}(k_0^\alpha) - \frac{1}{2}(k^\alpha)^2} = \mathcal{P}\int_0^\infty \frac{k^\alpha dk^\alpha f(k^\alpha)}{\frac{1}{2}(k_0^\alpha)^2 - \frac{1}{2}(k^\alpha)^2} - i\pi\,\mathrm{sgn}(k_0^\alpha)\Theta(E-E_\alpha)f(k_0^\alpha)\ .$$

Equation (2.12) is correct as written for $k_0^\alpha = +[2(E-E_\alpha)]^{1/2}$; the extension to $k_0^\alpha = -[2(E-E_\alpha)]^{1/2}$, necessary for application of the substitution rules, Eq. (2.19), is obtained by including $\mathrm{sgn}(k_0^\alpha)$ factors as in Eq. (2.14a). That is, all denominators of the form

$$\tfrac{1}{2}(k_0^\alpha)^2 + i\epsilon - \tfrac{1}{2}(k_1^\alpha)^2$$

must be replaced by

$$\tfrac{1}{2}(k_0^\alpha)^2 + i\epsilon\,\mathrm{sgn}(k_0^\alpha) - \tfrac{1}{2}(k_1^\alpha)^2$$

in Eq. (2.12).

Equation (2.17b) should read

$$\tilde{V}_{i0}^{\alpha\beta} = \mathrm{sgn}(k_0^\beta)(k_i^\alpha\omega_i^\alpha)^{1/2} V_{i0}^{\alpha\beta}(k_0^\beta\omega_0^\beta)^{1/2}\ ;$$

the "sgn" was omitted by the printer.

The notation of Eqs. (2.17) and (2.19) may cause some confusion. In direct application of the substitution rules of Eq. (2.19) the actual potential matrix element $V_{i0}^{\alpha\beta}$ is to be considered a function of $|k_0^\beta|$ only; thus the only effect of changing the sign of k_0^β $(\beta = 1, \ldots, n)$ in applying the substitution rules is to change the sign of the matrix element of Eq. (2.17b). That this is the case is immediately clear from the derivation.

None of these notational corrections affects the results or conclusions presented.

Atomic Radiationless Transition Probabilities to the 1s State and Theoretical K-Shell Fluorescence Yields, Vaclav O. Kostroun, Mau Hsiung Chen, and Bernd Crasemann [Phys. Rev. A 3, 533 (1971)]. The title of the article "Atomic Radiation*less* Transition Probabilities..." was misprinted "Atomic Radiation Transition Probabilities...."

Magnetic Moments of Five Levels in the Ground-State Configuration of N I and O II, V. Beltrán-López and Teodoro González E. [Phys. Rev. A 2, 1651 (1970)]. The paper was published without the Introduction. It appears below.

I. INTRODUCTION

In two previous articles[1,2] we calculated and com-