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Fizeau effect for neutrons

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It is shown that local-field effects in the neutron optical potential produce a nonvanishing Fizeau effect and that the magnitude of the effect is particularly large for cold neutrons with wavelengths $\lambda \geq 13 \text{ \AA}$.

In a recent article Arif *et al.*¹ have attempted to measure the Fizeau effect for thermal neutrons transmitted through a moving slab using a perfect-crystal neutron interferometer. When the slab is moving in a direction parallel to its boundaries (so that there is no motion of the boundaries themselves) the phase shift due to the motion of the slab is¹ proportional to dV/dE , where V is the optical potential that describes the effective interaction of the neutrons with the slab and E is the incident-neutron energy. In the elementary theory^{2,3} V is given by the mean value of the Fermi pseudopotential,

$$V = \left(\frac{2\pi\hbar^2}{m} \right) \rho b \quad (1)$$

where m is the neutron mass, ρ the number of nuclei per unit volume, and b the average bound coherent scattering length per nucleus. If b is independent of E , then $dV/dE = 0$ and the Fizeau phase shift vanishes.⁴⁻⁷ This was found¹ to be consistent with the experimental results for a fused-quartz slab and an incident-neutron wavelength $\lambda = 1.268 \text{ \AA}$, which gave an upper bound

$$\left| \frac{dV}{dE} \right| \leq 2.1 \times 10^{-8} \quad (2)$$

It has been suggested^{5,6} that a nonvanishing Fizeau effect will occur near a neutron resonance, where b has a significant energy dependence. The purpose of the present paper is to call attention to the fact that V also becomes energy dependent when one takes into account the local-field effects that are neglected in the approximate result (1). In this case one finds⁸⁻¹⁰ that the real part of the optical potential is given by

$$V = \left(\frac{2\pi\hbar^2}{m} \right) \rho b (1 + J) \quad (3)$$

in which the local-field correction is given for an isotropic medium by

$$J = \frac{2\pi\rho b}{k} \int_0^\infty [1 - g(r)] \sin(2kr) dr \quad (4)$$

where k is the incident-neutron wave vector and $g(r)$ the pair-correlation function for the atoms that comprise the medium. Since $E = (\hbar k)^2/2m$, we see that

$$\frac{dV}{dE} = \left(\frac{2\pi\rho b}{k} \right) \frac{dJ}{dk} \quad (5)$$

The qualitative behavior of J is dominated by the repulsive core of the interatomic potential so that a rough estimate can be obtained by putting

$$g(r) = \begin{cases} 0, & r < a, \\ 1, & r > a, \end{cases} \quad (6)$$

in which case⁹

$$J = J_0 \sin^2(ka)/(ka)^2, \quad (7)$$

where

$$J_0 = 2\pi\rho b a^2 \quad (8)$$

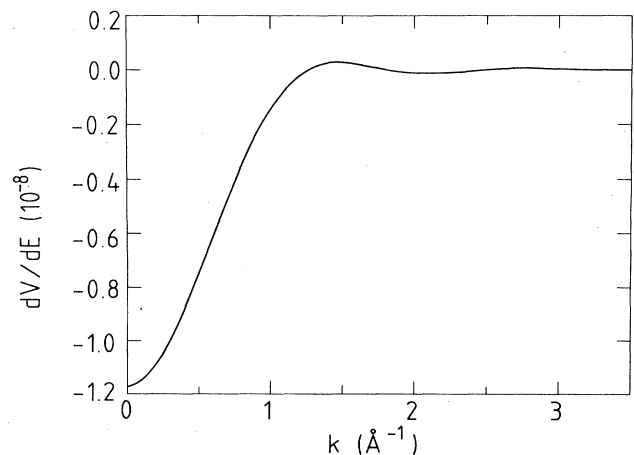


FIG. 1. Energy derivative of the neutron optical potential as a function of the neutron wave vector as calculated from (9) for fused quartz.

Hence, it follows from (5) that

$$\frac{dV}{dE} = 2[J_0^2/(ka)^4] \sin(ka) [ka \cos(ka) - \sin(ka)] \quad (9)$$

For fused quartz (amorphous SiO_2) $\rho = 0.022$ molecules/ \AA^3 , $b = 15.76$ fm/molecule, and $a = 2.5$ \AA , so that $J_0 = 1.3 \times 10^{-4}$. The corresponding results for dV/dE calculated from (9) are shown in Fig. 1. It is seen that dV/dE has a pronounced negative peak at $k = 0$ and an oscillating tail for $k \geq 1$ \AA^{-1} . These features are qualitatively consistent with our earlier results⁹ calculated directly from (4) for various monatomic liquids using experimentally determined pair-correlation functions. The value of dV/dE at

$k = 0$ in Fig. 1 is -1.2×10^{-8} , which is of the same order of magnitude as the experimental upper bound (2). However, for the wave vector $k = 2\pi/\lambda = 4.955$ \AA^{-1} , at which this upper bound was measured,¹ the expression (9) gives $dV/dE = -3.3 \times 10^{-12}$.

In conclusion, the above analysis shows that the Fizeau effect provides a direct method for the experimental investigation of local-field effects in the neutron optical potential and that these effects will be most easily observed using cold neutrons with $k \leq 0.5$ \AA^{-1} and $\lambda = 2\pi/k \geq 13$ \AA . Such measurements would be of particular importance, since it has recently been shown¹¹ that local-field effects often produce a significant correction to the values of b measured by neutron gravity refractometry.

¹M. Arif, H. Kaiser, S. A. Werner, A. Cimmino, W. A. Hamilton, A. G. Klein, and G. I. Opat, *Phys. Rev. A* **31**, 1203 (1985).

²E. Fermi, *Nuovo Cimento* **11**, 157 (1934), reprinted in *Enrico Fermi, Collected Papers*, edited by E. Segrè (University of Chicago Press, Chicago, 1962), Vol. I, p. 706.

³V. F. Sears, *Can. J. Phys.* **56**, 1261 (1978).

⁴M. A. Horne and A. Zeilinger, in *Neutron Interferometry*, edited by U. Bonse and H. Rauch (Clarendon, Oxford, 1979), p. 350.

⁵A. G. Klein, G. I. Opat, A. Cimmino, A. Zeilinger, W. Treimer,

and R. Gähler, *Phys. Rev. Lett.* **46**, 1551 (1981).

⁶M. A. Horne, A. Zeilinger, A. G. Klein, and G. I. Opat, *Phys. Rev. A* **28**, 1 (1983).

⁷A. Peres, *Am. J. Phys.* **51**, 947 (1983).

⁸M. Lax, *Phys. Rev.* **85**, 621 (1952).

⁹V. F. Sears, *Phys. Rep.* **82**, 1 (1982).

¹⁰M. Warner and J. E. Gubernatis (unpublished).

¹¹V. F. Sears, *Z. Phys. A* **321**, 443 (1985).