## Observation of laser-induced associative ionization in crossed-beam Na+Li collisions

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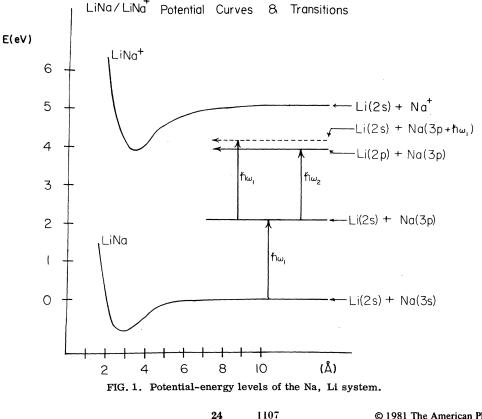
We report the observation of NaLi<sup>+</sup> produced by crossed-beam collisions of Na and Li in the presence of a single laser field. The magnitude of the cross section and its dependence on laser intensity are in accord with a simple second-order perturbation theory.

During the past three years a growing body of experimental evidence  $^{1-7}$  has demonstrated that collisional dynamics may be dramatically altered in the presence of intense laser radiation. In these phenomena the light fields play an active role in the overall interaction and are not simply the means by which initial and final states are prepared or detected.

In previous papers<sup>3,4</sup> we reported laser-induced production of  $Li_2^+$  and  $Na_2^+$  in collisions between identical particles. We report here the observation of NaLi<sup>\*</sup> when separate beams of Na and Li atoms are crossed in the presence of a laser field. The collisional equation is

$$Na^{*}(3p) + Li(2s) + \hbar \omega - NaLi^{*} + e.$$
 (1)

This new observation differs from previous reports in two important ways. First, the associative process is heteronuclear. It demonstrates that a laser field can induce bonding between chemically distinct particles in a two-body collision. Second, the effect is simpler to analyze because



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## NaLi<sup>+</sup>INTENSITY vs LASER FREQUENCY

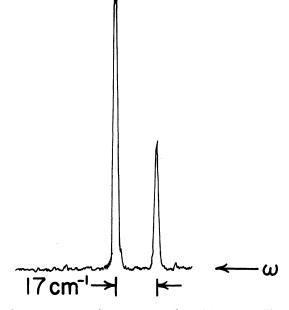


FIG. 2. Mass spectrum of NaLi<sup>+</sup> as a function of laser frequency. Peak positions are within 1 cm<sup>-1</sup> of Na $({}^{2}P_{3/2}, {}^{2}P_{1/2})$  transitions.

only one of the partners (sodium) interacts strongly with the light field. In previous experiments field excitation of colliding partners was necessarily symmetrical.

The experimental apparatus is quite similar to that described in earlier reports.<sup>4,5</sup> The lithium atomic-beam source was designed as a "double oven" similar to the sodium source but modified to permit continuous operation at temperatures near 1000 °C. The atomic beams were maintained at right angles in the horizontal scattering plane but moved radially somewhat closer together so as to increase particle-number density in the collision zone. Typical operating conditions were (a) alkali atom density  $\simeq 10^{12}$  atoms cm<sup>-3</sup>, (b) unfocused laser intensity  $\simeq 5 \times 10^4$  W cm<sup>-2</sup>. Invoking the usual assumptions of instrumental detection sensitivity and efficiency,<sup>4</sup> we estimate the cross section for the laser-induced processes  ${\sim}10^{-19}\,{\rm cm}^2$ which is in agreement with a perturbation calculation of collisional ionization.<sup>8</sup>

Figure 1 depicts the relevant potential curves and the optical-collisional interaction sequence. In the first step the laser field tuned to the Na(3s - 3p) transition excites a *virtual* state of sodium [approximately midway between Na\*(5s) and Na\*(4d) levels] at *twice* the resonant frequency. The second step, associative ionization, then takes place between the "dressed" sodium atom and ground-state lithium.

Figure 2 shows two peaks in the NaLi<sup>\*</sup> spectrum at laser frequencies within 1 cm<sup>-1</sup> of the Na\*( ${}^{2}P_{3/2}$ ,  ${}^{2}P_{1/2}$ ) levels. No other NaLi<sup>\*</sup> peaks were observed.

Theory<sup>8,9</sup> predicts that in the weak-field limit, intensity of the NaLi<sup>\*</sup> should be linear with laser intensity. Figure 3 shows the results of experiments in which the laser beam was systematically

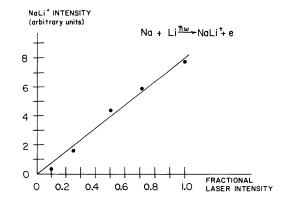


FIG. 3. Ion intensity vs laser intensity for laser-induced process. Unattenuated laser peak power  $\simeq 5 \times 10^4$  W cm<sup>2</sup>.

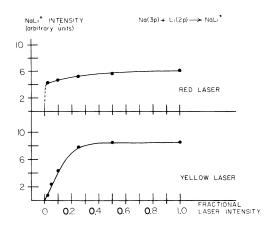


FIG. 4. Ion intensity vs laser intensity for two-laser experiment in which each laser resonantly populates an excited atomic state.

attenuated by a set of calibrated neutral density filters. Linearity of the plot is self-evident. This result contrasts sharply with those obtained from conventional associative ionization between excited states of Na and Li. In a separate set of experiments with two lasers, tuned, respectively, to the Na\*(3p) and Li\*(2p) transitions, we also observe NaLi<sup>\*</sup> production. In this case, however, we speculate that the process may proceed via a two-step mechanism in which the initial step is an

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inverse predissociation,

$$Na(3p) + Li(2p) \rightarrow NaLi^*$$
, (2a)

followed by photoionization of the quasimolecule,

$$NaLi^* + h\omega - NaLi^* + e.$$
 (2b)

Saturation of NaLi<sup>\*</sup> intensity at high laser power is a necessary consequence of such a two-step mechanism.<sup>10</sup> Attenuation experiments on each laser were repeated with the result presented in Fig. 4. The saturation characteristic is clearly evident.

It is interesting to speculate on two-laser experiments in which the first laser is tuned to the Na(3s - 3p) frequency but the second is nonresonant with respect to any atomic transition. If the first Na(3s - 3p) transition is "uncorrelated" with the collisional interaction,<sup>6</sup> then we may expect enhancement of the ion signal by a second nonresonant laser, the light field of which has no definite phase relationship with the first. The results of the other two-laser experiments, both resonant and nonresonant, will be reported shortly.

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- <sup>9</sup>S. E. Harris, unpublished lecture notes available from Stanford University, Department of Electrical Engineering, Palo Alto, California 94305.
- <sup>10</sup>The model-potential calculation of the NaLi<sup>+</sup> binding energy by P. Habitz and W. H. E. Schwartz [Chem. Phys. Lett. <u>34</u>, 248 (1975)] implies that the overall process (2) is endoergic by 0.2 eV, and therefore *conventional* associative ionization cannot explain the results. An independent *ab initio* calculation of the ground  $^{2}\Sigma^{+}$  state of NaLi<sup>+</sup> by W. Stevens of the National Bureau of Standards confirms this conclusion.

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