

**Erratum: Family of oscillatory electromagnetic pulses [Phys. Rev. A **108**, 063502 (2023)]**

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I have been made aware of an error in Sec. IV of my paper, relating to backflow.

I attempted to prove the absence of backflow in transverse electric (TE) and transverse magnetic (TM) pulses based on a general waveform. The electric and magnetic fields are calculated from a vector potential, which is the curl of the vector  $[0, 0, \psi]$ , where  $\psi$  is a solution of the wave equation  $(\nabla^2 - \partial_{ct}^2)\psi = 0$ .

In cylindrical coordinates  $\rho, \phi, z$  the solutions of the wave equation may be expressed in the form

$$\psi(\rho, \phi, z, t) = e^{im\phi} \int_0^\infty dk e^{-ikt} \int_0^k dq w(k, q) e^{iqz} J_m(\kappa\rho), \quad k^2 = q^2 + \kappa^2.$$

These scalar solutions are manifestly forward propagating since the longitudinal component of the wave vector  $q$  is non-negative. However, the vector electromagnetic fields are obtained from  $\psi$  by operations that include curl, and the curl of a vector has, in general, a different direction from the that of the original vector. It is not true in general that the momentum density  $\mathbf{p}$ , proportional to  $\mathbf{E} \times \mathbf{B}$ , has no regions with negative  $p_z$  (assuming net propagation in the positive  $z$  direction).

My attempted method of proof of positive-definite  $p_z$  was to examine the sign of the integrand in the  $z$  component of the momentum. The integral is over four wave-number variables  $k, q, k', q'$ , and two angle variables  $\chi, \chi'$  (which appear when the Bessel function is expressed as an integral over angle), but my “proof” considered the sign only in the dominant diagonal case  $k' = k, q' = q, \chi' = \chi$ .

In fact, there is a small amount of backflow. Figure 1 below shows the energy density and momentum density at  $t = 0$  for a TE or TM pulse based on the imaginary part of the simplest waveform,  $G = \frac{a^2}{R(R-iz)}$ ,  $R^2 = \rho^2 + (a + ict)^2$ . This solution of the wave equation is of the form given above, with  $m = 0$  and  $w = a^2 e^{-ka}$ . The regions of negative  $p_z$  are those outside the hyperboloids of revolution  $\rho^2 + a^2 = (1 + 2/\sqrt{3})z^2$ , shown in the figure as red curves. The momentum density is largely transverse in these regions of backflow. The maximum negative  $p_z$  is located at the points indicated by the diamonds. It is about 2.4% of the maximum positive value located at  $\rho = a\sqrt{2}/3, z = 0$ . More details may be found in Appendix 3C of [1]; see also [2–4], for other examples of backflow.

Electromagnetic pulses localized in space-time may be Lorentz transformed, as discussed in Chapter 5 of [1]. Pulses with net propagation along the  $z$  direction always have  $cP_z < U$  ( $P_z$  and  $U$  are the total pulse momentum along  $z$  and the total pulse energy) and a Lorentz boost along  $z$  with speed  $c^2 P_z / U$  will bring the pulses to their zero-momentum frame, in which  $P'_z = \int d^3 r' p'_z = 0$ . In the zero-momentum frame there are equal amounts of positive and negative momentum density, and net backflow equals net

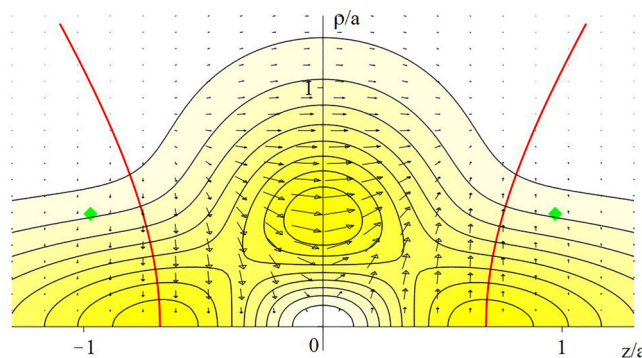


FIG. 1. Energy density (contours) and momentum density (arrows) for a TE or TM pulse based on the imaginary part of the waveform  $G$ , at  $t = 0$ . The red curves (hyperboloids of revolution, see text) are the boundaries of the backflow regions where  $p_z < 0$ . The location of the most negative  $p_z$  is indicated by the diamonds.

forward flow. In general, the amounts of forward and backward flow of momentum and energy (the momentum density and energy flux vectors  $\mathbf{p}$  and  $\mathbf{S}$  are related by  $\mathbf{S} = c^2\mathbf{p}$ ) depends on which inertial frame the electromagnetic pulse is observed.

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