## ERRATA

Relativistic Treatment of the Shielding of the Electron and Proton Magnetic Dipole Moments in Atomic Hydrogen. Roger A. Hegstrom [Phys. Rev. <u>184</u>, 17 (1969)]. The author is very grateful to H. Grotch, who has independently performed similar calculations (H. Grotch, private communication), for kindly pointing out an error in the transformed Hamiltonian [Eqs. (7)]. This error affects subsequent equations and numerical results in order  $\alpha^3$ . The first of Eqs. (7) should read

$$\begin{aligned} \Im c_{1} &= -\frac{e_{1}\hbar}{2m_{1}c} \left[ \left(\vec{\sigma}_{1} \cdot \vec{\mathbf{H}} + \vec{\mathbf{L}}_{1} \cdot \vec{\mathbf{H}}\right) \ 1 - \left(\frac{p_{1}^{2}}{2m_{1}^{2}c^{2}}\right) + a_{1}(h) \left(\vec{\sigma}_{1} \cdot \vec{\mathbf{H}} - \frac{\vec{\sigma}_{1} \cdot \vec{\mathbf{p}}_{1}\vec{\mathbf{p}}_{1} \cdot \vec{\mathbf{H}}}{2m_{1}^{2}c^{2}}\right) \right] \\ &- \frac{e_{2}\hbar}{2m_{2}c} \left[ \left(\vec{\sigma}_{2} \cdot \vec{\mathbf{H}} + \vec{\mathbf{L}}_{2} \cdot \vec{\mathbf{H}}\right) \left(1 - \frac{p_{2}^{2}}{2m_{2}^{2}c^{2}}\right) + a_{2}(h) \left(\vec{\sigma}_{2} \cdot \vec{\mathbf{H}} - \frac{\vec{\sigma}_{2} \cdot \vec{\mathbf{p}}_{2}\vec{\mathbf{p}}_{2} \cdot \vec{\mathbf{H}}}{2m_{2}^{2}c^{2}}\right) \right]. \end{aligned}$$

Equations (12) then become

$$\begin{split} \Delta E_e &= 2 \left[ 1 + a_e(h) \right] \mu_0 H m_s \left[ 1 - \frac{1}{2} \alpha^2 Z^2 \left( \frac{M}{M + m} \right)^2 \frac{1 + \frac{1}{3} a_e(h)}{1 + a_e(h)} + \frac{1}{6} \alpha^2 Z^2 \left( \frac{M}{M + m} \right)^2 \frac{1 + 2a_e(h)}{1 + a_e(h)} - \frac{1}{3} \alpha^2 Z^3 \left( \frac{m}{M + m} \right)^2 \right], \\ &- \Delta E_p = 2 \left[ 1 + a_p(h) \right] Z \mu_n H m_I \left[ 1 - \frac{1}{2} \alpha^2 Z^2 \left( \frac{m}{M + m} \right)^2 \frac{1 + \frac{1}{3} a_p(h)}{1 + a_p(h)} + \frac{1}{6} \alpha^2 Z^2 \left( \frac{m}{M + m} \right)^2 \frac{1 + 2a_e(h)}{1 + a_e(h)} - \frac{1}{3} \alpha^2 Z \left( \frac{m}{M + m} \right)^2 \right], \\ &- \frac{1}{3} \alpha^2 Z \left( \frac{M}{M + m} \right)^2 \right] \,. \end{split}$$

Equations (14) become

$$g_{e}(h) = g_{e} \left[ 1 - \frac{1}{3} \alpha^{2} Z^{2} \left( \frac{M}{M+m} \right)^{2} + \frac{1}{4\pi} \alpha^{3} Z^{2} \left( \frac{M}{M+m} \right)^{2} + \delta_{e} \right] , \qquad g_{p}(h) = g_{p} \left[ 1 - \frac{1}{3} \alpha^{2} Z \left( \frac{M}{M+m} \right)^{2} \right] ,$$

which are good to orders  $\alpha^3$  and  $\alpha^2 m/M$  if  $\delta_e$  is of order  $\alpha^3$  or higher and if  $\delta_p$  is of order  $\alpha^4$  or higher. Equation (15) becomes

$$\frac{g_e(h)}{g_p(h)} = \frac{g_e}{g_p} \left[ 1 + \frac{1}{4\pi} \alpha^3 \left( \frac{M}{M+m} \right)^2 + \delta_e \right] .$$

To obtain numerical results to order  $\alpha^3$ , we must find a value for the bound-state correction  $\delta_e$  to the Pauli moment of the electron. This correction is formally of order  $\alpha^3$  and must be added to the  $(\alpha^3/4\pi)[M/(M+m)]^2$  term in our Eq. (14) to obtain the total correction in this order.

Previously, to find  $\delta_e$  we used the radiative bound-state correction to the electron g factor calculated by Lieb to be  $-(26/15\pi)\alpha^3$ . The correct relationship of this quantity to  $\delta_e$  would be

$$\delta_{e} = -\frac{26}{15\pi} \alpha^{3} - \frac{1}{12\pi} \alpha^{3} \left(\frac{M}{M+m}\right)^{2} ,$$

because of our definition of  $\delta_e$  and our expression of  $g_e(h)$  as  $g_e$  times a correction. Recently, however, Grotch has recalculated the lowest-order radiative bound-state corrections to the electron g factor by another method [H. Grotch (private communication); and Phys. Rev. Letters (to be published)], and ob-

tains a value  $+(1/12\pi)\alpha^3$ , which differs from that of Lieb. Grotch's result indicates that in our Eq. (14)  $\delta_e = 0$ , and thus that our  $(\alpha^3/4\pi)[M/(M+m)]^2$  correction represents the *entire* lowest-order radiative correction to the electron g factor. The reason for this, according to Grotch's calculation, is that (i) the Pauli term in the Breit equation gives the entire lowest-order self-energy correction when the free-electron anomalous moment is taken for the Pauli moment, and (ii) the lowest-order vacuum polarization diagram gives zero contribution. Thus, in retrospect, we have found a very simple method for calculating the lowest-order radiative bound-state correction and its mass dependence. Finally, Eq. (16) must be corrected. If we accepted Lieb's result we would have

$$g_e(h)/g_p(h) = (g_e/g_p) [1 - (47/30\pi)\alpha^3] = (g_e/g_p) (1 - 1.94 \times 10^{-7}).$$

Accepting Grotch's recent calculation we find

$$g_e(h)/g_p(h) = (g_e/g_p)[1 + (1/4\pi)\alpha^3] = (g_e/g_p)[1 + 0.31 \times 10^{-7}].$$