ERRATA

Relativistic Treatment of the Shielding of the Electron and Proton Magnetic Dipole Moments in Atomic HY drogen. Roger A. Hegstrom [Phys. Rev. 184, 17 (1969)]. The author is very grateful to H. Grotch, who has independently performed similar calculations (H. Grotch, private communication), for kindly pointing out an error in the transformed Hamiltonian [Eqs. (7)]. This error affects subsequent equations and numerical results in order α^3 . The first of Eqs. (7) should read

$$
\mathcal{R}_1 = -\frac{e_1\hbar}{2m_1c} \left[(\vec{\sigma}_1 \cdot \vec{\mathbf{H}} + \vec{\mathbf{L}}_1 \cdot \vec{\mathbf{H}}) \mathbf{1} - \left(\frac{p_1^2}{2m_1^2c^2} \right) + a_1(h) \left(\vec{\sigma}_1 \cdot \vec{\mathbf{H}} - \frac{\vec{\sigma}_1 \cdot \vec{\mathbf{p}}_1 \vec{\mathbf{p}}_1 \cdot \vec{\mathbf{H}}}{2m_1^2c^2} \right) \right] - \frac{e_2\hbar}{2m_2c} \left[(\vec{\sigma}_2 \cdot \vec{\mathbf{H}} + \vec{\mathbf{L}}_2 \cdot \vec{\mathbf{H}}) \left(1 - \frac{p_2^2}{2m_2^2c^2} \right) + a_2(h) \left(\vec{\sigma}_2 \cdot \vec{\mathbf{H}} - \frac{\vec{\sigma}_2 \cdot \vec{\mathbf{p}}_2 \vec{\mathbf{p}}_2 \cdot \vec{\mathbf{H}}}{2m_2^2c^2} \right) \right].
$$

Equations (12) then become

$$
\Delta E_e = 2[1 + a_e(h)] \mu_0 H m_s \left[1 - \frac{1}{2} \alpha^2 Z^2 \left(\frac{M}{M+m} \right)^2 \frac{1 + \frac{1}{3} a_e(h)}{1 + a_e(h)} + \frac{1}{6} \alpha^2 Z^2 \left(\frac{M}{M+m} \right)^2 \frac{1 + 2a_e(h)}{1 + a_e(h)} - \frac{1}{3} \alpha^2 Z^3 \left(\frac{m}{M+m} \right)^2 \right],
$$

\n
$$
-\Delta E_p = 2[1 + a_p(h)] Z \mu_n H m_I \left[1 - \frac{1}{2} \alpha^2 Z^2 \left(\frac{m}{M+m} \right)^2 \frac{1 + \frac{1}{3} a_p(h)}{1 + a_p(h)} + \frac{1}{6} \alpha^2 Z^2 \left(\frac{m}{M+m} \right)^2 \frac{1 + 2a_p(h)}{1 + a_p(h)} - \frac{1}{3} \alpha^2 Z \left(\frac{M}{M+m} \right)^2 \right] - \frac{1}{3} \alpha^2 Z \left(\frac{M}{M+m} \right)^2 \right].
$$

Equations (14) become

$$
g_e(h)\!=\!g_e\!\left[1-\mbox{$\textstyle \frac{1}{3}$}\alpha^2 Z^2 \!\left(\mbox{$\frac{M}{M+m}$}\right)^{\!2}\right.\nonumber\\ \left.+\frac{1}{4\pi}\;\alpha^3 Z^2 \!\left(\mbox{$\frac{M}{M+m}$}\right)^2\!\!\!\!\right.\nonumber\\ +\left.\delta_e\right]\;, \qquad g_p(h)\!=\!g_p\!\left[1-\mbox{$\textstyle \frac{1}{3}$}\alpha^2 Z \!\left(\mbox{$\frac{M}{M+m}$}\right)^2\right]\,, \label{eq:1.1}
$$

which are good to orders α^3 and $\alpha^2 m/M$ if δ_{ρ} is of order α^3 or higher and if δ_{ρ} is of order α^4 or higher. Equation (15) becomes

$$
\frac{g_e(h)}{g_p(h)} = \frac{g_e}{g_p} \left[1 + \frac{1}{4\pi} \alpha^3 \left(\frac{M}{M+m} \right)^2 + \delta_e \right] .
$$

To obtain numerical results to order α^3 , we must find a value for the bound-state correction δ_e to the Pauli moment of the electron. This correction is formally of order α^3 and must be added to the $(\alpha^3/4\pi)[M/$ $(M + m)^2$ term in our Eq. (14) to obtain the total correction in this order.

Previously, to find $\delta_{\bm{\varrho}}$ we used the radiative bound-state correction to the electron g factor calculated by Lieb to be – $(26/15\pi)\alpha^3$. The correct relationship of this quantity to δ_e would be

$$
\delta_e = -\frac{26}{15\pi} \alpha^3 - \frac{1}{12\pi} \alpha^3 \left(\frac{M}{M+m}\right)^2,
$$

because of our definition of δ_e and our expression of $g_e(h)$ as g_e times a correction. Recently, however, Grotch has recalculated the lowest-order radiative bound-state corrections to the electron g factor by another method [H. Grotch (private communication); and Phys. Rev. Letters (to be published)], and ob-

tains a value + $(1/12\pi)\alpha^3$, which differs from that of Lieb. Grotch's result indicates that in our Eq. (14) $\delta_e = 0$, and thus that our $(\alpha^3/4\pi)[M/(M+m)]^2$ correction represents the *entire* lowest-order radiative correction to the electron g factor. The reason for this, according to Grotch's calculation, is that (i) the Pauli term in the Breit equation gives the entire lowest-order self-energy correction when the free-electron anomalous moment is taken for the Pauli moment, and (ii) the lowest-order vacuum polarization diagram gives zero contribution. Thus, in retrospect, we have found a very simple method for calculating the lowest-order radiative bound-state correction and its mass dependence. Finally, Eq. (16) must be corrected. If we accepted Lieb's result we would have

$$
g_e(h)/g_p(h) = (g_e/g_p)[1 - (47/30\pi)\alpha^3] = (g_e/g_p)(1 - 1.94 \times 10^{-7}) .
$$

Accepting Grotch's recent calculation we find

$$
g_e(h)/g_p(h)\!=\!(g_e/g_p)\big[1+(1/4\pi)\,\alpha^3\big] =\!(g_e/g_p)\big[1+0,31\times10^{-7}\big]\;.
$$