

photoproton energy distribution<sup>11</sup> which shows the ground state to be favored by a factor of 4 or more.<sup>12</sup>

We conclude that the experimental evidence on  $C^{12}(\gamma, p)B^{11}$  supports an independent-particle description of the giant resonance and suggests  $LS$  coupling as the more valid approximation.

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<sup>1</sup> We do not imply that only one excited state is involved. It is, however, convenient to discuss one resonance state corresponding to the gross resonance in the cross section around 22 Mev; in practice this gross state may well appear shared between several "fine structure" states of largely common parentage.

<sup>2</sup> At the giant resonance the respective proton energies are 5.5 Mev and 3.4 Mev. We may ignore transitions to other states of  $B^{11}$ ; they are eliminated by the Coulomb barrier and also effectively by experimental considerations such as the target thickness.

<sup>3</sup> J. Halpern and A. K. Mann, *Phys. Rev.* **83**, 370 (1951).

<sup>4</sup> E. D. Courant, *Phys. Rev.* **82**, 703 (1951).

<sup>5</sup> J. L. Burkhardt, *Phys. Rev.* **91**, 420 (1953).

<sup>6</sup> A. Reifman, *Z. Naturforsch.* **8a**, 505 (1953).

<sup>7</sup> J. S. Levinger and D. C. Kent, *Phys. Rev.* **95**, 418 (1954).

<sup>8</sup> D. H. Wilkinson, Proceedings of University of Pennsylvania Photoneuclear Conference, 1954 (unpublished); Proceedings of the University of Glasgow Nuclear Physics Conference, 1954 (to be published); *Phil. Mag.* (to be published).

<sup>9</sup> R. F. Christy, *Phys. Rev.* **89**, 839 (1953).

<sup>10</sup> A. M. Lane and D. H. Wilkinson, *Phys. Rev.* (to be published).

<sup>11</sup> W. E. Stephens and A. K. Mann, *Bull. Am. Phys. Soc.* **29**, No. 7, 26 (1954).

<sup>12</sup> It is only a limit that may be given from these experiments because identification of protons in the low-energy "tail" of the observed spectrum with a particular transition cannot be made with certainty. The assumptions involved in identification of the "tail" protons which resulted in the factor 4 were such as to make that value a lower limit.

## $(\gamma, p)$ and $(\gamma, n)$ Yield Ratios from Self-Conjugate Nuclei

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IT has been pointed out that the observed  $(\gamma, n)$  cross sections for light nuclei do not attain the sum rule limit.<sup>1,2</sup> Various interpretations have been given,<sup>1,2</sup> but in only a few cases are both  $(\gamma, p)$  and  $(\gamma, n)$  cross sections data available for comparison with the sum rule limit. In the cases of self-conjugate nuclei like  $Mg^{24}$  or  $Ca^{40}$ , however, calculation of the proton-to-neutron yield ratio on the basis of compound nucleus formation, can be made relatively safely since product nuclei from  $(\gamma, p)$  and  $(\gamma, n)$  reactions are mirrors of each other whose level structures are identical.

Calculation was made on the following basis: (1) The level density for the  $4n+3$  residual nuclei was taken as  $\omega(E) = C \exp(aE)$ , where  $a = d \log \omega / dE$  is taken to be a constant over the range of the residual energies

involved. This assumption fits well with  $(p, p')$  data<sup>3</sup> on Al from 3- to 10-Mev excitation; it can be extrapolated to zero in good agreement with the density of known levels in Al<sup>27</sup>. Values of  $a$  for different nuclear mass numbers  $A$  were obtained from the known levels of  $B^{11}$ ,  $Na^{23}$ , and  $Al^{27}$  and also from  $(p, p')$  data<sup>3</sup> on Al, Ni, and Ag. A smooth curve was drawn through these points and values of  $a$  interpolated for arbitrary  $A$ . (2) For the barrier penetration factor, a formula  $(1 - kB/E_p)$  was used. Here  $E_p$  is the energy of the outgoing protons,  $B$  is the classical barrier height for  $r_0 = 1.5 \times 10^{-13}$ , and  $k$  is adjusted so that the penetration factor makes the best fit to the quantum-mechanically calculated value.<sup>4</sup> Then,

$$r_{pn} = \frac{\exp\{a(E_e - b_p)\} - 1}{\exp\{a(E_e - b_n)\} - 1} \rightarrow \exp\{a(b_n - b_p)\}$$

when  $E_e$  is large.

Here  $r_{pn}$  is proton-neutron yield ratio,  $E_e$  is the excitation energy of the compound nucleus,  $b_n$  is the neutron binding energy and  $b_p$  is the proton binding energy plus  $kB$ . The asymptotic values of  $r_{pn}$  are given in Table I. If we estimate the sum of integrated  $(\gamma, p)$  and

TABLE I. Asymptotic value of  $r_{pn}$  and the sum of  $(\gamma, p)$  and  $(\gamma, n)$  integrated cross sections for certain self-conjugate nuclei.

	$C^{12}$	$O^{16}$	$Mg^{24}$	$Si^{28}$	$S^{32}$	$Ca^{40}$	Dipole sum rule limit by Levinger and Bethe	
							$x=0$	$x=1$
$r_{pn}$	1.6	1.9	3.6	4.4	5.6	9.7		
$\frac{A}{NZ} \int_{\text{giant resonance}} \{\sigma(\gamma, p) + \sigma(\gamma, n)\} dE$	0.046	0.054	0.047	0.049	0.074	0.078	0.060	0.108

$(\gamma, n)$ <sup>5</sup> cross sections we obtain the results in Table I. Here the dipole sum rule is resumed except for very light nuclei.

Preliminary measurements of  $(\gamma, p)$  yields on  $Mg^{24}$ ,  $Si^{28}$ ,  $S$ , and  $Ca$  were made by Johansson<sup>6</sup> and the results show good qualitative agreement, but with rather lower yields. The difference may be due to oversimplification in the above arguments or to contributions from direct photoelectric processes. Measurements of  $r_{pn}$  using alpha particles are now under way here which might shed more light on this problem.

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<sup>1</sup> Montalbetti, Katz, and Goldemberg, *Phys. Rev.* **91**, 659 (1953).

<sup>2</sup> R. Nathans and J. Halpern, *Phys. Rev.* **93**, 437 (1954).

<sup>3</sup> P. C. Gugelot, *Phys. Rev.* **93**, 425 (1954).

<sup>4</sup> J. M. Blatt and V. F. Weisskopf, *Theoretical Nuclear Physics* (John Wiley and Sons, Inc., New York, 1952), pp. 352.

<sup>5</sup> Taken from Summers-Gill, Haslam, and Katz, *Can. J. Phys.* **31**, 70 (1953).

<sup>6</sup> S. A. E. Johansson (private communication); see following Letter [*Phys. Rev.* **97**, 1186 (1955)].