

excited by fast electrons. Presumably a specifically nuclear interaction is not involved, which suggests that there may be similarities between electron excitation of levels and Coulomb excitation by heavy particles.

Figures 1 and 2 and the data at 60° permit a rough, limited measurement of the angular distribution of the cross section for the inelastic scattered electrons. These data are plotted in Fig. 3 along with the elastic cross section. The flatter angular distributions of the inelastic electrons appear to be in agreement with a preliminary theory.<sup>5</sup>

Inelastic processes have also been observed in copper.

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<sup>1</sup> Hofstadter, Fechter, and McIntyre, *Phys. Rev.* **92**, 978 (1953).

<sup>2</sup> R. Britten, *Phys. Rev.* **88**, 283 (1952).

<sup>3</sup> K. E. Davis and E. M. Hafner, *Phys. Rev.* **73**, 1473 (1948).

<sup>4</sup> E. H. Roderick, *Proc. Roy. Soc. (London)* **A201**, 348 (1950).

<sup>5</sup> L. I. Schiff (private communication).

### Pion Production Ratios\*

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MEASUREMENTS have recently been made<sup>1</sup> on the production of charged pions in the reaction  $\text{Be}^9 + p$ . The  $\pi^+/\pi^-$  ratio  $\rho$  shows a striking energy dependence,<sup>2</sup> decreasing from  $\rho \approx 6$  at  $E_p = 1$  Bev to  $\rho = 1.8$  at  $E_p = 2.3$  Bev. The present note interprets this behavior in terms of an "excited state" of the nucleon with isotopic spin  $\frac{3}{2}$ .

Pion production is envisioned as follows: two nucleons collide, forming a "compound state"; on emerging, one or both of the nucleons are in an excited state, which subsequently decays by  $\pi$  emission. This separation of the process into two distinct stages linked by a quasi-stable excited nucleon state is without justification; it should not, however, vitiate conclusions based on charge independence arguments and forms a convenient physical picture. For the sake of simplicity the nucleon is assumed to have only one excited state, which can decay to the ground state by emission of a single  $\pi$  meson. The excited state must therefore have  $T' = \frac{1}{2}$  or  $\frac{3}{2}$ ; meson production can be characterized by the indices 1 or 2, according as one or both nucleons emerge from the collision in an excited state. Because of the threshold for excitation, the production cross section satisfies  $\sigma_1 > \sigma_2$  for low  $E_p$ , while  $\sigma_2 \gtrsim \sigma_1$  for high  $E_p$ . The variation of  $\rho$  with  $E_p$  can be related to that of  $\sigma_1/\sigma_2$ , provided that  $\rho_1 \neq \rho_2$ . This immediately excludes the case  $T' = \frac{1}{2}$ , for then the charge symmetries of the system are identical whether 0, 1, or 2 nucleons are excited, so that  $\rho = \rho_1 = \rho_2 = \text{const}$ .

The case  $T' = \frac{3}{2}$  remains to be examined. For this purpose let  $\sigma^{xy}(T^c)$  be the cross section for production of a compound state of isotopic spin  $T^c$  by the collision of incident nucleons  $x, y$ . Charge independence requires that  $\sigma^{np}(1) = \frac{1}{2}\sigma^{pp}(1) = \sigma^s$ , but does not relate  $\sigma^s = \sigma^{np}(0)$  to  $\sigma^s$ ; of course  $\sigma^{pp}(0) = 0$ . Consider a  $p$ - $p$  collision with excitation of both nucleons: the compound state has  $T^c = 1$ ,  $T_z^c = 1$ , and the Clebsch-Gordon coefficients for decomposition into two  $T' = \frac{3}{2}$  nucleons yield respective fractional weights of  $\frac{2}{3}$  and  $\frac{1}{3}$  for the combinations  $T_z' = (\frac{1}{2}, \frac{1}{2})$  and  $T_z' = (\frac{3}{2}, -\frac{1}{2})$ . The fractional weights for decay of the excited state by  $\pi^+$ ,  $\pi^0$ ,  $\pi^-$  emission are (100),  $(\frac{1}{3}, \frac{2}{3}, 0)$ ,  $(0, \frac{2}{3}, \frac{1}{3})$  according as  $T_z' = \frac{3}{2}, \frac{1}{2}, -\frac{1}{2}$ . Combining these factors leads to over-all fractional weights for  $\pi$  production of  $(13/15, 14/15, 3/15)$ . These last are normalized to give a total production of two  $\pi$  mesons.

For a  $p$ - $p$  collision with excitation of one nucleon, the compound state decomposes into  $T' = \frac{3}{2}$ ,  $T = \frac{1}{2}$  with fractional weights  $\frac{1}{2}$  and  $\frac{3}{4}$  for  $T_z' = (\frac{3}{2}, -\frac{1}{2})$  and  $(\frac{1}{2}, \frac{1}{2})$ , respectively. The resultant over-all fractional weights for  $\pi$  production are  $(\frac{1}{2}, \frac{1}{2}, 0)$ , normalized to a total of one meson. The total cross sections for  $\pi$  production in a  $p$ - $p$  collision are then

$$\begin{aligned}(\pi^+): & (1/15)[26\sigma_2^s + 15\sigma_1^s], \\(\pi^0): & (1/15)(28\sigma_2^s + 15\sigma_1^s), \\(\pi^-): & (1/15)[6\sigma_2^s].\end{aligned}\quad (1)$$

In a similar way the cross sections for meson production in an  $n$ - $p$  collision are

$$\begin{aligned}(\pi^+ \text{ or } \pi^-): & (1/15)[14\sigma_2^s + (5/2)\sigma_1^s + 10\sigma_2^a], \\(\pi^0): & (1/15)[2\sigma_2^s + 10\sigma_1^s + 10\sigma_2^a].\end{aligned}\quad (2)$$

Combining Eqs. (1) and (2) in the ratio 4 to 5 for a  $\text{Be}^9$  nucleus leads to

$$\begin{aligned}(\pi^+): & [1.28\sigma_2^s + 0.54\sigma_1^s + 0.37\sigma_2^a], \\(\pi^0): & [0.90\sigma_2^s + 0.82\sigma_1^s + 0.37\sigma_2^a], \\(\pi^-): & [0.70\sigma_2^s + 0.09\sigma_1^s + 0.37\sigma_2^a].\end{aligned}\quad (3)$$

When  $\sigma_1 \gg \sigma_2$ ,  $\rho \approx 6$ ; when  $\sigma_2 \gg \sigma_1$ ,  $1 \lesssim \rho \lesssim 1.8$ , depending on the relative magnitudes of  $\sigma_2^a$  and  $\sigma_2^s$ .

Of course these considerations are subject to many qualifications, but in terms of the model employed they indicate that (i) wide variation in  $\rho$  implies  $T' = \frac{3}{2}$  for the excited nucleon state; (ii) the observed variation of  $\rho$  suggests  $\sigma_1 \gg \sigma_2$  for  $E_p = 1$  Bev,  $\sigma_2 \gg \sigma_1$  for  $E_p = 2.3$  Bev. The variation (ii) is plausible if the  $T' = \frac{3}{2}$  state is identified with the "resonance" appearing at about 0.2 Bev in  $\pi$ -nucleon scattering, as has been suggested from energy analysis of the produced mesons.<sup>1</sup> According to (ii) and Eq. (3), the relative  $\pi^0$  production should also show a marked variation with energy: for low  $E_p$ ,  $\pi^0$  production exceeds the total charged  $\pi$  production; for high  $E_p$ ,  $\pi^0$  production is intermediate between  $\pi^+$  and  $\pi^-$ .

The same model can be applied to calculation of various other production ratios. For example, in  $n$ - $p$  collisions the ratio of  $(\pi^-, \pi^0)$  plus  $(\pi^-, \pi^+)$  production to single  $\pi^-$  production is  $(5.6\sigma_2^s + 4\sigma_2^a)/\sigma_1^s$ . For neutrons in the 1-2 Bev range, this ratio is observed<sup>3</sup> to be of order 5, suggesting that  $\sigma_2 \sim \sigma_1$  for an average  $\bar{E}_n \sim 1.5$  Bev. This is compatible with the conclusions from  $\text{Be}^9 + p$ .

It is tempting to speculate on the existence of a further "excited state" of the nucleon with  $T'' = 5/2$ , which would decay to the ground state by emission of two mesons. It could not appear in  $\pi$ -nucleon scattering but would be associated with three- and four-meson production in a nucleon-nucleon collision. If the  $T''$  state has a well-defined energy, the threshold for these processes may be high. We could, for example, compute  $\rho$  for a  $p$ - $p$  collision involving  $T'' = 5/2$  states. The cross sections concerned are  $\sigma_3^s$  (excitation of one nucleon to  $T' = \frac{3}{2}$ , the other to  $T'' = 5/2$ ) and  $\sigma_4^s$  (excitation of both to  $T'' = 5/2$ ). The meson production cross sections are

$$\begin{aligned}(\pi^+): & (1/50)(78\sigma_3^s + 76\sigma_4^s), \\(\pi^0): & (1/50)(39\sigma_3^s + 88\sigma_4^s), \\(\pi^-): & (1/50)(33\sigma_3^s + 36\sigma_4^s).\end{aligned}\quad (4)$$

Excitation of the  $T'' = 5/2$  state gives a characteristic value  $\rho \approx 2$ , more or less independent of the ratio  $\sigma_3/\sigma_4$ . To (4) must of course be added to cross sections (1), which give  $\rho > 4$ . Thus observation of  $\rho < 4$  would be indication of a  $T'' = 5/2$  state.

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<sup>1</sup> L. C. L. Yuan and S. J. Lindenbaum, *Phys. Rev.* **93**, 1431 (1954).

<sup>2</sup> L. C. L. Yuan (private communication).

<sup>3</sup> Fowler, Shutt, Thorndike, and Whittemore, *Phys. Rev.* (to be published).