

tungsten "cat whisker" was also made, and the point was etched electro-chemically in our own laboratory. It is possible that a sharper contact point has been obtained than that of commercially available crystals. The results shown were obtained with  $K$  band klystrons. Millimeter wave tubes have also been used, but because of their lower powers they give less harmonic power at the higher frequencies.

The present developments seem not only to have made spectroscopy in the "no man's land" between optical and radio waves a practical reality but at the same time to have improved by an order of magnitude the performance in the previously worked 2- to 5-mm region. The components of the system are so broad-banded that one can obtain fourth or fifth harmonic power from the full tuning range of a  $K$  band klystron without retuning either the multiplier or detector. This reduces to insignificance the previously difficult task of "finding" fourth or fifth harmonic power.

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† The present work does not represent the first generation and detection of radiation in the region of one to two millimeters. As early as 1923 [Phys. Rev. 21, 378 (1923); 21, 587 (1923)], E. F. Nichols and J. D. Tear closed the gap between optical and radio waves in the sense that energy was generated and detected.

<sup>2</sup> Johnson, Trambarulo, and Gordy, Phys. Rev. 84, 1178 (1951).

<sup>3</sup> Gilliam, Johnson, and Gordy, Phys. Rev. 78, 140 (1950).

The secondary particle from the decay undergoes a large angle scattering ( $30^\circ$ ) in traversing a lead plate. It seems very probable that this is a nuclear scattering and that the secondary particle is a  $\pi$ -meson.

A search was made for close pairs of penetrating particles coming from interactions in carbon. Such pairs might be interpreted as being the product of the decay of a short-lived, low  $Q$ , neutral particle. Twenty-two pairs of penetrating particles with angles of separation of the order of  $5^\circ$  or less were found in 100 showers in carbon. The probability of a chance coincidence was calculated from the average angular distribution of shower particles in the showers and the angle between the shower particles in question. The result of the calculations indicated that fewer pairs were observed than were expected by chance. An angular correlation has been observed by photographic plate workers.<sup>5</sup> From our data it appears that less than  $\frac{1}{2}$  percent of the penetrating secondaries arise from the decay of a  $\zeta^0$  (if such exist).

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<sup>1</sup> Walker, Duller, and Sorrels, Phys. Rev. 86, 865 (1952).

<sup>2</sup> Fretter, May, and Nakada, Phys. Rev. 89, 168 (1953).

<sup>3</sup> Leighton, Wanlass, and Anderson, Phys. Rev. 89, 148 (1953).

<sup>4</sup> Barker, Butler, and York, Phil. Mag. 43, 1201 (1952).

<sup>5</sup> Danysz, Lock, and Vekutieli, Nature 169, 364 (1952).

## Unstable Particles from Penetrating Showers\*

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IN the course of an experiment on penetrating showers<sup>1</sup> in carbon, cloud-chamber pictures of 500 showers originating in carbon plates inside of the cloud chamber were obtained. Among these pictures only two examples of decays of neutral  $V$  particles were found. Thus only about one shower per two hundred and fifty gives rise to a  $V^0$  decay. This may be compared to the rate obtained by Fretter *et al.*<sup>2</sup> with a similar apparatus of about one  $V^0$  decay per 18 penetrating showers. This comparison is perhaps not quite correct, as Fretter pictures were on the average of higher quality than those obtained with this apparatus. On the other hand, the lack of cascade development in the carbon plates should have made the search efficiency relatively higher than in Fretter's experiment. It is assumed that the search efficiency is high only if the  $V^0$  is emitted into a less populous part of the shower, that is, at an angle greater than the median angle for the shower. This means an efficiency of about 50 percent. This assumption gives about one  $V^0$  decay per 125 showers in carbon. This is still more than a factor five lower than the rate obtained by Fretter. Since two cases were found when about 14 were expected, it seems unlikely that this results from a statistical fluctuation.

The major difference between the two experiments seems to be that these showers were generated in carbon while those of Fretter occurred in lead. It is possible also that the median energy of the primaries triggering this apparatus was somewhat higher than those triggering Fretter's apparatus because of a more complex triggering arrangement.

It seems possible that  $V_1^0$ 's (i.e., the majority of the  $V^0$ 's<sup>3</sup>) are formed primarily in interactions of lower energy than those selected here ( $\bar{E} \sim 20$  Bev). If this were the case, showers in a heavy nucleus where secondary multiplication occurs would give rise to more  $V_1^0$ 's than showers from a light nucleus such as carbon. In particular, it might be that  $\pi$ -meson nucleon collisions are the only sources of  $V_1^0$  particles.

Of the 1000 shower particles observed traversing an average distance of 15 cm, one was found to decay in flight. This may be compared with the Manchester result<sup>4</sup> of 14  $V^\pm$  decays in 8000 tracks traversing 15 cm each.

## The Short-Lived Radioisotopes $P^{28}$ and $Cl^{32}$ †

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THE series of radioactive isotopes  $B^8$ ,  $N^{12}$ ,  $Na^{20}$ , and  $Al^{24}$  has all decay by positron emission, at least some of the branches of which lead to excited states that decay by alpha-particle emission.

We have observed two new activities which result from the bombardment of silicon and sulfur by 20-Mev protons from the UCLA cyclotron. From threshold and energy considerations, we ascribe these activities to  $P^{28}$  and  $Cl^{32}$ , additional members of the above series. The method of detection involved the use of a scintillation gamma-ray spectrometer with a NaI crystal.

The half-life of the  $Cl^{32}$  activity is  $0.306 \pm 0.004$  second, and in addition to positrons it emits gamma-radiation of energy  $4.8 \pm 0.2$  Mev. The half-life of the  $P^{28}$  was found to be  $0.280 \pm 0.010$  second. It emits positrons and gamma-radiation up to an energy of 7 Mev.

The thresholds for the excitation of these activities has been measured relative to the threshold for  $Mg^{24}(p,n)Al^{24}$  at  $15.4 \pm 0.3$  Mev.<sup>2</sup> In the case of  $Si^{28}(p,n)P^{28}$  the threshold is  $15.6 \pm 0.5$  Mev, and for  $S^{32}(p,n)Cl^{32}$  the threshold is  $14.3 \pm 0.5$  Mev.

From these threshold values one calculates a mass for the  $P^{28}$  of  $28.0012 \pm 0.0007$  amu, and for the  $Cl^{32}$  a mass of  $31.9963 \pm 0.0007$  amu, using the values of  $Li^3$  for the masses of  $Si^{28}$  and  $S^{32}$ . These nuclei are apparently just barely stable to proton emission.

We have searched for alpha-particles from these activities, using a ZnS screen and photomultiplier, but without positive results. The sensitivity of our arrangement can be indicated by the following statement: Either (a) the alpha-particles have energy less than 1 Mev or (b) the transition which results in their emission has a probability less than 10 percent of those transitions which result in gamma-ray emission.

We have also obtained some results on  $Al^{24}$  from the reaction  $Mg^{24}(p,n)Al^{24}$ . We observe gamma-radiations of energy  $7.1 \pm 0.2$  Mev,  $5.3 \pm 0.2$  Mev,  $4.3 \pm 0.2$  Mev, and  $2.9 \pm 0.2$  Mev. Our value for the half-life is  $2.10 \pm 0.04$  seconds which agrees within experimental error with the value obtained by Birge.<sup>3</sup>

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<sup>1</sup> L. W. Alvarez, Phys. Rev. 80, 519 (1950).

<sup>2</sup> A. C. Birge, Phys. Rev. 85, 753 (1952).

<sup>3</sup> C. W. Li, Phys. Rev. 88, 1038 (1952).