The Angular Momenta of the Excited

States of Pt196†

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HE angular correlation of the Pt196 gamma-rays has been measured before¹ and was interpreted as characteristic for two quadrupole transitions between levels of angular momenta I of 2, 2, and 0. Internal conversion data² indicate electric quadrupole radiation for both transitions.

Recently, a number of parity unfavored gamma-transitions were found to be a mixture of magnetic dipole and electric quadrupole radiation.³⁻⁵ A similar mixed transition may be expected for the first gamma-ray of the Pt¹⁹⁶ cascade, being due to a parity unfavored transition between states of the same angular momentum (I=2). The accuracy of the former Pt¹⁹⁶ angular correlation measurements was not sufficient to allow a determination of the dipole admixture intensity; the Pt196 gamma-gamma correlation was therefore remeasured using much improved experimental techniques and evaluation methods.6

To avoid the detection of scattered radiation two techniques were employed which yielded the same angular correlation (within 1.5 percent): (1) the two NaI(Tl) crystals which were used to detect the Pt196 gamma-rays of 0.330 Mey and 0.358 Mey were covered with 0.2-cm lead on the front and 1.5-cm lead on the side, or (2) pulse-height discrimination was employed making the use of any lead absorber unnecessary. The discriminator bias was set at the lower edge of the 0.330-Mev photopeak. The data for the Pt¹⁹⁶ gamma-gamma angular correlation are shown in Fig. 1.

The Au¹⁹⁶ sources used were in the form of dilute solutions of AuCl₃ in order to obtain the maximum angular correlation;^{7,8} however, sources of dry AuCl₃ and Au¹⁹⁶ imbedded in gold (Au¹⁹⁶ produced by the $Au^{197}(n, 2n)Au^{196}$ reaction) showed the same angular correlation as above (within 3 percent).

The data, fitted by least squares to an angular correlation function $W(\theta) = 1 + a_2 \cos^2\theta$, $+a_4 \cos^4\theta$, taking the finite angular resolution of the apparatus into account,⁹ give for the correlation coeffi-







FIG. 2. Decay of Au^{196} and the angular momenta of the excited states of Pt^{196} .

cients: $a_2 = -1.02 \pm 0.05$ and $a_4 = +1.32 \pm 0.05$. The maximum anisotropy, calculated from the coincidence rate at $\theta = 90^{\circ}$, $\theta = 270^{\circ}$, and $\theta = 180^{\circ}$ after correction for the angular resolution was determined as $(a_2+a_4) = +0.28 \pm 0.03$. This correlation function is consistent with the assignment 2, 2, 0 for the angular momenta of the excited states and the ground state of Pt196, the first transition being a mixture of 95 percent quadrupole and 5 percent dipole radiation and the two components being out of phase (in Lloyd's¹⁰ notation; same sign of the reduced electric and magnetic matrix elements in the notation of Biedenharn and Rose¹¹). The data on internal conversion are consistent with both transitions being mainly electric quadrupole (+5 percent magnetic dipole in the first transition), indicating the same parity for all three levels involved (Fig. 2).

These assignments would suggest a strong cross-over transition of 0.688-Mev energy to the ground state of Pt196. The search for such a gamma-radiation with a scintillation spectrometer resulted in an upper limit for the intensity of the cross-over transition of less than 1 percent.

It is interesting to note that, in the very similar case of the Cd¹¹⁴ gamma-rays,^{4,5} the magnetic dipole (96 percent) and electric quadupole (4 percent) components are in phase, whereas the case of Pt^{196} there is a phase difference of π .

Recently, an angular correlation of the Pt196 gamma-rays in agreement with these measurements has been observed by the Zurich group.12

[†] Supported by the U. S. Atomic Energy Commission.
[†] R. M. Steffen and D. M. Roberts, Phys. Rev. 82, 332 (1951).
^{*} Steffen, Huber, and Humbel, Helv. Phys. Acta 22, 167 (1949).
^{*} Aceppil, Frauenfelder, and Walter, Helv. Phys. Acta 24, 335 (1951).
^{*} E. D. Klema and F. K. McGowan, Phys. Rev. 87, 524 (1952).
^{*} Rolf M. Steffen (to be published).
^{*} J. C. Kluyver and M. Deutsch, Phys. Rev. 87, 203 (1952).
^{*} Rolf M. Steffen (to be published).
^{*} Walter, Huber, and Zunti, Helv. Phys. Acta 23, 697 (1952).
^{*} Brolf M. Steffen and W. Zobel, Phys. Rev. 87, 203 (1952).
^{*} Rolf M. Steffen (to be published).
^{*} Walter, Huber, and Zunti, Helv. Phys. Acta 23, 697 (1952); see also
^{*} H. Frauenfelder, Annual Review of Nuclear Science, Vol. 2, (to be published).
¹⁰ S. P. Lloyd, Phys. Rev. 85, 904 (1952).
^{*} L. C. Biedenharn and M. E. Rose, Revs. Modern Phys. (to be published).

lished). ¹² E. Heer *et al.* (private communication by H. Frauenfelder).

States of S^{33} and Si^{29} from (d, p) Stripping

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WE have used a proportional counter telescope to examine the angular distributions within the range 0° to 90° of several groups of protons from the reactions $S^{32}(d, p)S^{33}$ and $Si^{28}(d, p)Si^{29}$ with 8-Mev deuterons. By comparing these distributions with the theories of the stripping process,^{1,2} the appropriate value of the orbital angular momentum l of the captured neutron has been determined in each case. These values are given in Tables I and II, together with the excitation energies of the final states. The spins of these states have one of the values $l \pm \frac{1}{2}$. Some of the

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