

The agreement between experiment and the calculation of Koppe, as shown in Figs. 1 and 2, is on the whole striking, although the departure from the curve in Fig. 1 of the low temperature points for niobium and vanadium is beyond the limit of experimental error. Nevertheless, the general fit of Koppe's curve suggests a real regularity displayed by all superconductors, a regularity which may be regarded as an approximate "law of corresponding states."

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† City College of New York, New York, New York.

‡ Barnard College, Columbia University, New York, New York.

<sup>1</sup> W. Heisenberg, *Z. Naturforsch.* **IIa**, 185 (1937).

<sup>2</sup> H. Koppe, *Ann. phys.* **1**, 405 (1947).

<sup>3</sup> W. H. Keesom and P. H. van Laer, *Physica* **5**, 193 (1938).

<sup>4</sup> W. H. Keesom and M. Désirant, *Physica* **8**, 273 (1941).

<sup>5</sup> P. L. Bender and C. J. Gorter, *Physica* **18**, 1 (1952). We are indebted to Professor Gorter for a copy of their paper in advance of publication.

<sup>6</sup> Brown, Zemansky, and Boorse, *Phys. Rev.* **86**, 134 (1952).

<sup>7</sup> J. R. Clement and E. H. Quinell, *Phys. Rev.* **79**, 1028 (1952). ( $\gamma$  was

calculated from magnetic data and taken to be  $8.92 \times 10^{-4}$  cal/mole-deg<sup>2</sup>.)

<sup>8</sup> Parkinson, Simon, and Spedding, *Proc. Roy. Soc. (London)* **A207**, 177 (1951).

<sup>9</sup> Worley, Zemansky, and Boorse, *Phys. Rev.* (to be published). See also the Report of the Low Temperature Conference of the ONR, General Electric Research Laboratories, Schenectady, New York, October 1952.

<sup>10</sup> J. R. Clement and E. H. Quinell, *Phys. Rev.* **85**, 502 (1952). ( $\gamma$  was calculated from magnetic data and taken to be  $8.92 \times 10^{-4}$  cal/mole-deg<sup>2</sup>.)

<sup>11</sup> W. H. Keesom and J. A. Kok, *Physica* **1**, 175, 595 (1934).

## Evidence for Domain Structure in Antiferromagnetic CoO from Elasticity Measurements

M. E. FINE

Bell Telephone Laboratories, Murray Hill, New Jersey

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YOUNG'S modulus in CoO on cooling undergoes a large decrease near the Néel point.<sup>1</sup> Street and Lewis<sup>1</sup> suggest that this is due to the influence of an applied stress on the preferred orientation of the spin vectors in antiferromagnetic domains. A more detailed study and explanation of the elasticity of CoO is herein reported.

A rod of CoO was prepared by pressing powdered CoO and sintering at 1130°C in N<sub>2</sub>. X-rays revealed only the CoO structure; the apparent density after sintering was 86 percent of the density of CoO.

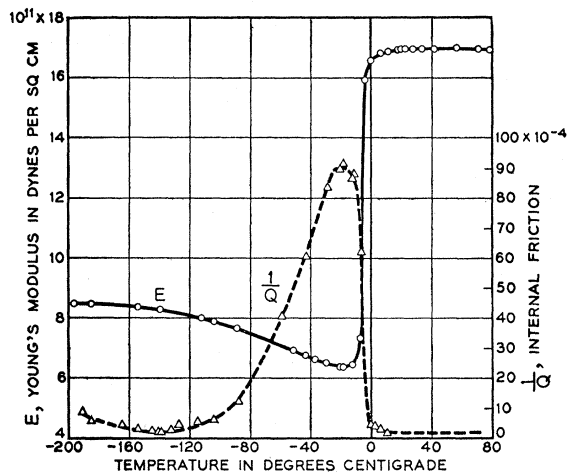


FIG. 1. Young's modulus and internal friction of CoO.

Young's modulus ( $E$ ) and the internal friction ( $1/Q$ ) were measured from 80 to  $-196^\circ\text{C}$  in the pressed and sintered rod, employing both electromagnetic and piezoelectric excitation and detection of longitudinal vibrations (35 to 60 kc sec<sup>-1</sup>). On cooling from 20 to  $-20^\circ\text{C}$ , as shown in Fig. 1,  $E$  drops from 17 to  $6.3 \times 10^{11}$  dynes cm<sup>-2</sup>, the region of maximum change being 0 to  $-8^\circ\text{C}$ . (The Néel temperature, determined from paramagnetic susceptibilities, is  $-2^\circ\text{C}$ .)<sup>2</sup> The effect persists to low temperature, for at

$-196^\circ\text{C}$   $E$  is  $8.5 \times 10^{11}$  dynes cm<sup>-2</sup>. The internal friction ( $1/Q$ ) for a constant maximum strain amplitude of  $10^{-7}$ , also shown in Fig. 1, rises on cooling from  $2 \times 10^{-4}$  (the background) to a peak value of  $91 \times 10^{-4}$  at  $-20^\circ\text{C}$ , then decreases to a minimum at  $-150^\circ\text{C}$  and rises again to  $8.5 \times 10^{-4}$  at  $190^\circ\text{C}$ . At room temperature and  $-20^\circ\text{C}$ , increasing the maximum strain amplitude from  $10^{-8}$  to  $10^{-5}$  has little effect on  $E$  or  $1/Q$ , but at  $-190^\circ\text{C}$   $E$  decreases 1 percent and  $1/Q$  increases from 5 to  $25 \times 10^{-4}$ .

The general course of our modulus-temperature curve and that reported by Street and Lewis<sup>1</sup> are in substantial agreement; however, our modulus values are larger by approximately a factor of 5, and the Néel temperature appears to be  $-2^\circ\text{C}$  in our sample rather than  $16^\circ\text{C}$ . The discrepancy among modulus values and Néel points may be due to composition and porosity variations.

In considering possible explanations the following properties of CoO seem important: The formation of antiferromagnetism causes an increase in volume; the coefficient of linear expansion is high in the temperature range 5 to  $-75^\circ\text{C}$ , with a peak of  $22 \times 10^{-6}$  deg C<sup>-1</sup> at  $-2^\circ\text{C}$ . Antiferromagnetism distorts CoO from cubic (NaCl structure) to tetragonal;<sup>3,4</sup>  $c/a$  becomes increasingly less than one as the temperature is decreased, being 0.9884 at  $-180^\circ\text{C}$ .<sup>3</sup> The antiferromagnetically aligned atomic magnetic moments are parallel to  $\langle 100 \rangle$ .<sup>5</sup> Presumably it is this direction that becomes shortened.

The large decrease in  $E$  implies that an applied stress produces strains in addition to, and in some cases even larger than, the normal elastic strain. In antiferromagnets an effect analogous to the stress-induced change in domain structure in ferromagnets is suggested.<sup>6</sup> Neighboring domains within a crystal are postulated, differing as to which of the three possible orientations of the  $c$  axis is present. Then, for example, tension would increase the area of those domains whose  $c$  axes are nearly perpendicular to the stress axis, thereby increasing the sample length. The amount of strain from a given amount of domain boundary movement would increase on cooling since  $c/a$  decreases, but domain boundary movement becomes more difficult owing to the lattice strain produced. Thus,  $E$  from this effect would not continuously decrease on cooling. Stress-induced antiferromagnetic domain boundary movement, like domain movement in ferromagnets, is expected to be hysteretic; consequently, the proposed mechanism seems especially reasonable for low temperatures where  $1/E$  and  $1/Q$  increase with strain amplitude and, furthermore, well-defined domains are expected.

Near the Néel temperature, where the magnetic moments are loosely coupled, another mechanism which may yield a strain and internal friction is a stress-induced change in the degree of antiferromagnetic order. The strain would come from the resulting change in volume and tetragonality.  $E$  and  $1/Q$  are more conceivably independent of strain amplitude by this mechanism.

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<sup>1</sup> R. Street and B. Lewis, *Nature* **168**, 1036 (1951).

<sup>2</sup> H. Bizette, *Ann. phys.* **1**, 295 (1946).

<sup>3</sup> N. C. Tombs and H. P. Rooksby, *Nature* **165**, 442 (1950).

<sup>4</sup> S. Greenwald and J. S. Smart, *Nature* **166**, 523 (1950).

<sup>5</sup> Shull, Strauser, and Wollan, *Phys. Rev.* **83**, 333 (1951).

<sup>6</sup> R. M. Bozorth, *Ferromagnetism* (D. Van Nostrand, Inc., New York, 1951), pp. 472, 595.

## Errata

**The Symmetrical Pseudoscalar Meson Theory of Nuclear Forces**, MAURICE M. LÉVY [*Phys. Rev.* **86**, 806 (1952)]. In the last portion of Eq. (3), replace  $[(2/\pi)K_0(\mu r)]^2$  with  $(\mu/2M)[(2/\pi)K_1(\mu r)]^2$ .

**Angular Distribution of Shower Particles as a Function of Depth**, L. EYGES AND S. FERNBACH [*Phys. Rev.* **82**, 287 (1951)]. Equation (1) should read:  $\pi(E_0, E, \theta, t)2\pi\theta d\theta dE = \pi_{\log}(E_0, E, t)dE \cdot P(x, s)\theta d\theta \cdot (E/E_0)^2$ . The present addresses