

calculations are made with an assumed line half-width of 29.2 Mc/sec at 1 mm Hg. The transitions are paired for all transitions found, except in the case of the  $5_{08}-5_{15}$  line, where the lower frequency component lies in the 3800-Mc/sec region. Corresponding inversion splittings are listed in Table VIII. The  $5_{08}-5_{15}$  inversion is not truly an observed splitting, since it is calculated from the position of the pure rotation transition.

The agreement between observed and calculated values is good at low values of  $J$ . Above  $J=11$  the calculated value rapidly falls below the observed value, indicating the need of higher order terms.

In the case of high  $J$  transitions, the rotational transition frequency is so slight that the lines are close doublets. Above  $J=11$  measurements were made at the maximum intensity of the unresolved doublet. For all transitions above  $J=7$ , the centrifugal distortion is so slight that the five-place tables of the reduced energy caused errors in excess of the distortion. Table IX gives

TABLE IX. Rotational transitions (Mc/sec).

Transition	Frequency		Distortion	
	Rigid rotor	Observed	Observed	Calculated
$5_{08} \rightarrow 5_{15}$	1314.31			-2.40
$6_{06} \rightarrow 6_{16}$	361.56	360.44	1.12	0.70
$7_{07} \rightarrow 7_{17}$	95.50	97.25	-1.75	-2.56
$8_{08} \rightarrow 8_{18}$	24.29	25.06	-0.77	
$9_{09} \rightarrow 9_{19}$	6.07	6.37	-0.30	
$10_{0,10} \rightarrow 10_{1,10}$	1.66	1.60	0.06	
$11_{0,11} \rightarrow 11_{1,11}$		0.36		

the calculated and observed rotational energy differences. The centrifugal distortion was calculated, using equation for Table VI above.

## Centrifugal Distortion in Asymmetric Molecules. II. HDS

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A perturbation method for relating the theory of centrifugal distortion in asymmetric top molecules to observed microwave Q branch,  $a$  or  $c$  type transitions, is presented. The formula for the distortion correction is expressed in terms of the total angular momentum,  $J$ , the symmetry axis momentum of the nearest symmetric top,  $K$ , and five distortion constants. The formula yields a satisfactory fit to the observed spectrum of HDS ( $\kappa \approx -0.5$ ). The electric dipole moment is determined as  $1.02 \pm 0.02$  debye. The inertia defect and distortion constants are calculated. The effective structure for the HDS molecule so determined is in agreement with infrared determinations.

### INTRODUCTION

THE vibration-rotation energies of polyatomic molecules have been derived to a second order of approximation for special cases by a number of authors. In general, the method followed has been equivalent to that of Wilson and Howard.<sup>1</sup> This work may be summarized by stating that the rotational term values of a molecule depend upon these quantities: the equilibrium structure of the molecule, the normal frequencies of vibration, and the nonvanishing coefficients of the anharmonic (cubic) terms in the expression for the binding potential energy of the molecule. Nielsen<sup>2</sup> has calculated the matrix elements for the rotational energy including, to second order, vibration-rotation energies.

Unfortunately the correction of centrifugal distortion in asymmetric top rotational term values is, in general, an inconvenient process. The explicit diagonalization of the reduced energy matrix may be accomplished for only a relatively few low  $J$  terms, when the asymmetry coefficient and the six centrifugal distortion coefficients

are included as unknowns. Since the experimental data are best used to evaluate the centrifugal distortion coefficients, the diagonalization of the reduced energy matrix for each trial value of these coefficients is awkward for comparison with experiment.

A method of arriving at an explicit form of the formula for centrifugal distortion correction terms to the rigid rotor frequencies has been demonstrated.<sup>3</sup> This paper will discuss a method of determining centrifugal distortion correction factors in terms of the general centrifugal distortion matrix elements.

### I. THEORY

The method used here is that of determining the centrifugal distortion correction in terms of rigid rotor frequencies specifically for Q branch " $a$ " type or " $c$ " type transitions. This dependence may be determined explicitly in terms of quantum or pseudo-quantum numbers and the rigid rotor absorption frequency, by a modified application of conventional perturbation theory.

To review briefly, the reduced energy matrix is diagonal in  $J$  and  $M$ ; the nonvanishing elements in  $K$  (i.e.,  $P_z$ ) are  $(K|H|K)$ ,  $(K|H|K \pm 2)$ , and  $(K|H|K \pm 4)$ .

\* R. B. Lawrance and M. W. P. Strandberg, Phys. Rev. **83**, 363 (1951).

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<sup>1</sup> E. B. Wilson, Jr., and J. B. Howard, J. Chem. Phys. **4**, 262 (1936).

<sup>2</sup> H. H. Nielsen, Phys. Rev. **60**, 794 (1941).

TABLE I. Identification of parameters.

	$I^r$	$II^r$	$III^r$
$x$	$b$	$c$	$a$
$y$	$c$	$a$	$b$
$z$	$a$	$b$	$c$
$F$	$\frac{1}{2}(\kappa-1)$	0	$\frac{1}{2}(\kappa+1)$
$G$	1	$\kappa$	-1
$H$	$-\frac{1}{2}(\kappa+1)$	1	$\frac{1}{2}(\kappa-1)$
$\kappa$		$\frac{2b-a-c}{a-c}$	

These elements have been given by Nielsen<sup>2</sup> explicitly in terms of the molecular parameters.

The reduced energy matrix  $E(\kappa)$  may be set up for a given value of  $J$ , the order of which is  $(2J+1)$  since  $-J \leq K \leq J$ . This matrix may immediately be factored to four submatrices as in the case of the rigid rotor.<sup>4</sup> Explicitly these submatrices in terms of the original energy elements are:

$$E^{\pm} = \begin{pmatrix} a_{00} & \sqrt{2}a_{02} & \sqrt{2}a_{04} & 0 & \cdots & \cdots & \cdots \\ \sqrt{2}a_{02} & a_{22} \pm a_{2,-2} & a_{24} & a_{26} & 0 & \cdots & \cdots \\ \sqrt{2}a_{04} & a_{24} & a_{44} & a_{46} & a_{48} & 0 & \cdots \\ 0 & a_{26} & a_{46} & a_{66} & a_{68} & a_{6,10} & \cdots \\ \cdots & \cdots & \cdots & \cdots & \cdots & \cdots & \cdots \end{pmatrix}, \quad (1)$$

where  $E^-$  has the same elements as  $E^+$  after the first row and column have been deleted, and

$$0^{\pm} = \begin{pmatrix} a_{11} \pm a_{1,-1} & a_{13} \pm a_{1,-3} & a_{15} & 0 & \cdots & \cdots & \cdots \\ a_{13} \pm a_{1,-3} & a_{33} & a_{35} & a_{37} & 0 & \cdots & \cdots \\ a_{15} & a_{35} & a_{55} & a_{57} & a_{59} & 0 & \cdots \\ 0 & a_{37} & a_{57} & a_{77} & a_{79} & a_{7,11} & \cdots \\ \cdots & \cdots & \cdots & \cdots & \cdots & \cdots & \cdots \end{pmatrix}. \quad (2)$$

In deriving these, use has been made of the relations

$$a_{K,K} = a_{-K,-K}$$

and

$$a_{K,K+2} = a_{K+2,K} = a_{-K,-K-2} = a_{-K-2,-K}$$

$$a_{K,K+4} = a_{K+4,K} = a_{-K,-K-4} = a_{-K-4,-K}. \quad (3)$$

The matrices are presented in the general notation of King, Hainer, and Cross<sup>5</sup> (hereinafter referred to as KHC) so that the identification of explicit energy levels is immediately obtained from their work, and hence will not be duplicated here. The difference between KHC's matrices and those given above is that the latter give the entire energy of the molecule including distortion corrections.

<sup>4</sup> Obviously there is no loss in generality in considering only the reduced energy matrix since it contains the entire complexity of the asymmetric rotor problem. The addition of  $\frac{1}{2}(a+c)J(J+1)$  to  $E(\kappa)$  yield the total energy.

<sup>5</sup> King, Hainer, and Cross, *J. Chem. Phys.* **11**, 27 (1943).

The matrix elements for the rotational reduced energy are given explicitly by Nielsen<sup>6</sup> as

$$(K|H|K) = (R_0 \mp R_1 K + R_2 K^2 + R_3 K^4), \quad (4)$$

$$(K|H|K \pm 2) = \{f(J, K \pm 1)\}^{\frac{1}{2}} \times \{R_4 - [K^2 + (K \pm 2)^2]R_5(a-c)/2\}, \quad (5)$$

$$(K|H|K \pm 4) = \{[f(J, K \pm 1)] \times [f(J, K \pm 3)]\}^{\frac{1}{2}} R_6(a-c)/2, \quad (6)$$

where

$$R_0 = \frac{1}{2}F(a-c)J(J+1) - \frac{1}{2}D_J(1-c)J^2(J+1)^2,$$

$$R_1 = 2Z_e \sum_r \xi_r e_r, \quad R_3 = -\frac{1}{2}D_K(a-c),$$

$$R_2 = \frac{1}{2}(G-F)(a-c) - \frac{1}{2}D_{JK}(a-c)J(J+1),$$

$$R_4 = \frac{1}{4}H(a-c) + \frac{1}{2}\delta_J(a-c)J(J+1),$$

$$f(J, n) = [J(J+1) - n(n-1)][J(J+1) - n(n+1)], \quad (7)$$

and  $D_J$ ,  $D_{JK}$ ,  $D_K$ ,  $\delta_J$ ,  $R_5$ , and  $R_6$  are the six centrifugal distortion terms dependent upon the vibration frequencies and molecular structure parameters.  $R_1$  arises from internal angular momentum and is hereafter neglected.  $a$ ,  $b$ , and  $c$  are defined as the effective reciprocal moments of inertia of the molecule which when corrected for zero point vibration yield the equilibrium values,  $a_e$ ,  $b_e$ , and  $c_e$ .<sup>7</sup> The latter are defined as the reciprocal moments of inertia the molecule would have if the atoms were located at the exact minima of the molecular binding potential, and may be expressed

$$a_e = h/(8\pi^2 I_{aa^e}) \text{ cps},$$

$$b_e = h/(8\pi^2 I_{bb^e}) \text{ cps},$$

$$c_e = h/(8\pi^2 I_{cc^e}) \text{ cps},$$

$$a \geq b \geq c. \quad (8)$$

Planck's constant has been factored from all elements to reduce them to frequency units. The  $I^r$  representation is used for  $a$ -type Q branch transitions and the  $III^r$  representation for  $c$ -type transitions. The actual definition of H, F, and G in terms of  $a$ ,  $b$ , and  $c$  are determined by the type of representation used. Table I has been prepared from KHC for this purpose. Further discussion of this point is given later. In the  $I^r$  representation  $a$ -type transitions are between term levels of  $E^+$  and  $E^-$ , or  $0^+$  and  $0^-$ . Similarly for the  $III^r$  representation,  $c$ -type transitions are given by the same selection rules.

<sup>6</sup> Nielsen's choice of phase factor for  $P_x$  and  $P_y$  differs from KHC. We have modified Nielsen's results to agree with KHC by changing the sign of the  $(K|H|K \pm 2)$  elements. Either choice of phase yields identical results for the  $(K|H|K)$  and  $(K|H|K \pm 4)$  elements. In any case the modification is purely formal. Note that Nielsen's  $D_J$ ,  $D_{JK}$ ,  $D_K$ ,  $\delta_J$ ,  $R_5$ , and  $R_6$  have been normalized to  $\frac{1}{2}(a-c)$ .

<sup>7</sup> This is not quite correct, for there still remains a small correction due to vibration-rotation interaction which is dependent neither on vibration quantum number (zero point vibration), nor rotational quantum number (centrifugal distortion) i.e., Nielsen's  $x_0$ ,  $y_0$ , and  $z_0$ .

The energy matrices (1) and (2) show clearly that the symmetric top  $K$ -degeneracy is removed in the  $K$ th level by a perturbation of order  $K$ . To illustrate, consider the  $E^\pm$  matrix. The  $0^+$  level is a singlet level (nondegenerate in zero order). The  $2^+$  level is split from the  $2^-$  level because of the second-order coupling with the  $0^+$  level through the off-diagonal elements  $\sqrt{2}a_{02}$ . The term  $a_{2,-2}$  is neglected for the moment since it is only a centrifugal distortion correction and small compared to the absorption energy. Both  $K=2$  levels are perturbed almost equally by the  $4^\pm$  levels. So the  $K=2$  degeneracy is removed in second order. The  $4^\pm$  degeneracy is removed due to a second-order coupling with the  $2^\pm$  levels. Since the  $2^\pm$  degeneracy is removed in second order, the  $4^\pm$  degeneracy is removed in fourth order, and so on. The degeneracy of the levels  $K^+$  and  $K^-$  is thus eliminated in  $K$ th order due to the progressively greater splitting of the  $(K-2)^\pm$  and  $(K-4)^\pm$ , etc., levels.

Observation of this fact simplifies the calculation of the effect of including centrifugal distortion terms in the matrix. We postulate that the matrix has been progressively diagonalized by second-order perturbation theory beginning with the next to the lowest  $K$  level in the matrix. Thus the new diagonal elements of the  $K-2$  levels,  $(a_{K-2, K-2})_{p^\pm}$ , are assumed nondegenerate. The amount by which the levels are split is of course closely related to the absorption frequency  $\nu_{K-2}$ . It differs slightly from the absorption frequency since the  $K^\pm$  levels, which will eventually have their degeneracy removed, will further perturb the diagonal  $K-2$  elements. The latter effect is taken into account by writing:<sup>8</sup>

$$(a_{K-2, K-2})_{p^+} - (a_{K-2, K-2})_{p^-} \cong \nu_{K-2} [1 - a_{K, K-2}^2 / (a_{K, K} - a_{K-2, K-2})^2].$$

This of course reduces to  $\nu_{K-2}$  in the region where second-order perturbation is applicable, i.e., where

$$a_{K, K-2}^2 / (a_{K, K} - a_{K-2, K-2})^2 = H^2 f(J, K-1) / 64(G-F)(K-1)^2 \ll 1. \quad (10)$$

By similar reasoning the interaction of the  $K$  level with the  $K+2$  level is negligibly small, thus

$$E^\pm \cong a_{K, K-2}^2 / a_{K, K} \pm (a_{K-2, K-2}^\pm)_p + a_{K, K-4}^2 / a_{K, K} \pm (a_{K-4, K-4}^\pm)_p \quad (11)$$

and hence

$$\nu_K \cong \nu_{K-2} a_{K, K-2}^2 / (a_{K, K} - a_{K-2, K-2})^2 + \nu_{K-4} a_{K, K-4}^2 / (a_{K, K} - a_{K-4, K-4})^2. \quad (12)$$

The term involving  $a_{K, K-4}$  is second order in a cen-

<sup>8</sup> The rigid rotor frequency is not calculated so crudely, of course. Such an approximation is only justified since we are interested in the *additional* contribution due to centrifugal distortion to an accuracy of a few percent.

trifugal distortion correction, and hence may be neglected.

The terms above are then written separating the rigid rotor contributions from the centrifugal distortion contributions.

$$\nu_K = (\nu_K)_{\text{rigid}} \left\{ 1 + \frac{2\delta_{K, K-2}}{a_{K, K-2}'} - \frac{2(\delta_{KK} - \delta_{K-2, K-2})}{(a_{KK}' - a_{K-2, K-2}')} + D_{K-2} \right\} \quad (13)$$

where:

$$a_{K, K-2} = a'_{K, K-2} (\text{rigid}) + \delta_{K, K-2} (\text{distortion}), \text{ etc.}$$

$$\nu_{K-2} = (\nu_{K-2})_{\text{rigid}} (1 + D_{K-2})$$

and  $D_{K-2}$  is the centrifugal distortion contribution to  $\nu_{K-2}$ .

The term  $D_{K-2}$  is of course a similar term to that given above, but now in terms of  $D_{K-4}$ , and so on. Finally, also, for low  $K$  values (1, 2, 3) the fourth off-diagonal terms,  $(K|H|K' \pm 4)$ , must be considered. When Eq. (13) is properly summed, using Eqs. (4), (5), (6), and (7), it may be written for odd and even  $K$ 's as follows:

$$\begin{aligned} \nu_o(J_K) = \nu_r(J_K) \{ & 1 + 2KJ(J+1)\delta_J/H \\ & + (K-1)J(J+1)D_{JK}/(G-F) \\ & + \frac{2}{3}K(K^2+2)[D_K/(G-F) - 2R_6/H] \\ & + [16(G-F)/H^2R_6 - 4D_K/(G-F)]_{K(\text{even})} \\ & + [64(G-F)/H^2R_6]_{K(\text{odd}) \geq 3} \\ & - [2D_K/(G-F)]_{K(\text{odd})} \}. \quad (14) \end{aligned}$$

This is our desired equation.

It is well to reconsider the restrictions of this formula. First, one must stay away from the diagonal of perturbation divergence, i.e.,

$$\frac{H^2 f(J, K-1)}{64(G-F)^2(K-1)^2} \ll 1. \quad (15)$$

Further discussion of this point is given in reference 3. Equation (14) is used to correct for the distortion in the observed microwave spectrum of HDS. Though this molecule is quite asymmetric ( $\kappa \cong -0.5$ ) the convergence criterion given above is satisfied for the observable levels. It is to be noted, however, that though the convergence criterion is satisfied even for the observed high  $J$  high  $K$  lines, it may not be satisfied for the low  $K$  terms of the same  $J$  value. In other words, the convergence of a level  $J_K$  may be assured, but the convergence of the splittings of the  $J_{K-2}$ ,  $J_{K-4}$ , etc., may not be satisfactory. Such convergence is postulated in the derivation of Eq. (14). Since the distortion of the  $J_K$  transition is composed of its own

terms, plus those of the lower order transitions, i.e.,  $J_{K-2}, J_{K-4}$ , etc., each of these lower order contributions must be in turn calculable by this method to achieve Eq. (14). This is certainly not the case in general. We only note that the contributions from the lower order  $K$ 's diminish rapidly with decreasing  $K$  so that little error is caused as long as the main terms, the  $K-2$ , and possibly  $K-4$  terms, meet the convergence requirement.

Second, it may be noted again that the above analysis is appropriate to Q-branch "a" and "c" type transitions. The type of transition covered is governed through  $F, G$  and  $H$ , by the representation used:  $I'$  for "a" type and  $III'$  for "c" type transitions. The value of  $K$  in Eq. (14) corresponds to the KHC index  $K_{-1}$  or  $K_{+1}$  for "a" and "c" type transitions respectively.

The restriction given above is not as drastic as may be first surmised, for molecules having significant distortion ("light" molecules) and yet a microwave spectrum will, in general, satisfy it.

Finally, this analysis could be extended with more difficulty to  $P$  and  $R$  branches. At present such an extension seems of little general importance, since molecules with  $P$  or  $R$  branches in the microwave region will generally have large moments of inertia, and consequently small centrifugal distortion correction; or they will be nearly symmetric tops (only one small moment of inertia), and hence symmetric top formulas may be reasonably applied; or the qualitative analysis of reference 3 may also be used.

## II. CENTRIFUGAL DISTORTION OF HDS

The experimental values of the rotational absorption frequencies of HDS<sup>32</sup>, observed for the first time in the microwave region, and identified by the Stark effect, are given in Table III. This data is now to be fitted to the distortion formula Eq. (14). The rigid rotor frequency is given by:

$$\nu_r(J_K) = \frac{1}{2}(a-c)[\Delta E(\kappa)]. \quad (16)$$

The process of fitting will be described briefly. A study of Eqs. (14) and (16) shows that there are seven adjustable parameters, i.e., the five distortion constants and  $\frac{1}{2}(a-c)$  and  $\kappa$ . Fairly accurate values of the latter two constants are obtained from low  $J$  transitions. Assuming these are correct, the distortion constants are determined using the experimental frequencies and the

TABLE II.  $\kappa$ ,  $\frac{1}{2}(a-c)$  and distortion constants of hydrogen deuterium sulfide.

$\frac{1}{2}(a-c) = 3.2664 \text{ cm}^{-1}$
$\kappa = -0.47767$
$R_0 = -2.808 \times 10^{-4} \text{ cm}^{-1}$
$\delta_J = +2.531 \times 10^{-5} \text{ cm}^{-1}$
$D_{JK} = +9.133 \times 10^{-4} \text{ cm}^{-1}$
$D_K = (-1.64 \times 10^{-4} \text{ cm}^{-1})^*$
$R_6 = +0.695 \times 10^{-5} \text{ cm}^{-1}$

\* Assumed.

method of least squares. The process is then reversed;  $\nu(J_K)$  is calculated from the distortion constants just evaluated and compared with the observed  $\nu_o(J_K)$ . Discrepancies are attributed to the fact that the  $\nu_r(J_K)$ , as determined from  $\frac{1}{2}(a-c)$  and  $\kappa$ , are not exactly correct. Therefore, slight variations in  $\frac{1}{2}(a-c)$  and  $\kappa$  are made and the computations repeated until an entirely consistent set of results is obtained.

The actual computations would seem to be quite laborious, but fortunately, Eq. (14) lends itself to an analytical simplification which reduces the computations involved by a large factor. Taking values of  $\frac{1}{2}(a-c)$  and  $\kappa$ ,  $\nu_r(J_K)$  is calculated for each value of  $J_K$ .  $K$  in the present case is the KHC  $K_{-1}$  index for the asymmetric rotor. The differences  $\nu_o(J_K) - \nu_r(J_K)$  are taken and when finally the correct values of  $\nu_r(J_K)$  are found, these differences will be the centrifugal distortion corrections. The following ratios are formed for each line

$$a(J_K) = [\nu_o(J_K) - \nu_r(J_K)] / \nu_r(J_K). \quad (17)$$

Now, if the differences  $a((J+1)_K) - a(J_K)$  are taken between lines of equal  $K_{-1}$ , the only distortion terms remaining are those which are a function of  $J$ . With slight rearrangement, a series of equations are obtained

$$\frac{a((J+1)_K) - a(J_K)}{2(J+1)} = \left[ \frac{D_{JK}}{(G-F)} + 2\frac{\delta_J}{H} \right] K - \frac{D_{JK}}{(G-F)} = AK + B, \quad (18)$$

which represents a straight line as a function of  $K$ . The values of the constants  $A$  and  $B$  are taken as those which best fit this linear relationship and are found by the method of least squares. It turns out that if  $\frac{1}{2}(a-c)$  and  $\kappa$  as chosen are almost correct, then the values of  $A$  and  $B$  are essentially independent of  $\nu_r(J_K)$ . Since  $\frac{1}{2}(a-c)$  and  $\kappa$  are known to three figure accuracy from low  $J$  transitions, the above computation need only be carried out once or twice during the process of fitting.

$A$  and  $B$  thus determine the  $J$  dependence of distortion and when substituted back into Eq. (14), a new series of equations are obtained involving only  $K$  dependent and constant terms. These can be written

$$\begin{aligned} a_K &= a(J_K) - J(J+1)(AK+B) \\ &= \frac{2}{3}K(K^2+2) \left[ \frac{D_K}{(G-F)} - \frac{2R_6}{H} \right] \\ &\quad + \left[ \frac{16(G-F)}{H^2} R_6 - \frac{4D_K}{(G-F)} \right]_{K(\text{even})} \\ &\quad + \left[ \frac{64(G-F)}{H^2} R_6 \right]_{K(\text{odd}) \geq 3} - \left[ \frac{2D_K}{(G-F)} \right]_{K(\text{odd})} \end{aligned} \quad (19)$$

where the "even" and "odd" notations indicate the terms which are present for even and odd values of  $K$  respectively. By taking the differences,  $a_{K+2} - a_K$ , the

TABLE III. Comparison of calculated and observed frequencies for hydrogen deuterium sulfide (in Mc/sec).

Initial state	Final state	Distortion correction	Calculated frequency	Observed frequency $\pm 0.05$ Mc/sec	Obs.-Calc. frequency
1 <sub>1,1</sub>	1 <sub>1,0</sub>	-73.43	51,073.30	51,073.27	-0.03
2 <sub>2,1</sub>	2 <sub>2,0</sub>	-49.61	11,283.73	11,283.83	+0.10
3 <sub>2,2</sub>	3 <sub>2,1</sub>	-220.38	53,204.97	53,200.93	-4.04
4 <sub>3,2</sub>	4 <sub>3,1</sub>	-125.09	10,862.97	10,861.07	-1.90
5 <sub>3,3</sub>	5 <sub>3,2</sub>	-411.22	40,917.57	40,929.20	+11.63
6 <sub>4,3</sub>	6 <sub>4,2</sub>	-177.27	7939.65	7936.74	-2.91
7 <sub>4,4</sub>	7 <sub>4,3</sub>	-516.94	27,557.68	27,566.31	+8.63
8 <sub>4,5</sub>	8 <sub>4,4</sub>	-1111.14	75,545.79	75,551.73	+5.94
9 <sub>5,5</sub>	9 <sub>5,4</sub>	-529.15	17,209.60	17,212.61	+3.01
10 <sub>5,6</sub>	10 <sub>5,5</sub>	-1123.07	47,906.66	47,905.36	-1.30
11 <sub>6,6</sub>	11 <sub>6,5</sub>	-476.79	10,233.30	10,235.81	+2.51
12 <sub>6,7</sub>	12 <sub>6,6</sub>	-1008.25	28,843.69	28,842.84	-0.85

coefficient of the  $K$  dependent term and the correct value of  $\kappa$  may be determined. There is insufficient data to evaluate all the constants unambiguously, therefore the value of  $D_K$  is assumed to be known (see Section IV) since, in the case of HDS, it involves the smallest centrifugal distortion correction. Equation (19) for  $K=1$  then determines the value of  $\frac{1}{2}(a-c)$ ,  $R_6$  is calculable from the  $K$ -dependent term, and  $R_6$  results from the final remainder. It is at this point that the experimental data do not agree with the theory. The theory indicates that the remaining terms in  $R_6$  should show a four-to-one alternation with  $K$  odd or even respectively, since  $D_K/(G-F)$  is practically negligible. Experimentally, the remainder is a constant for any value of  $K \geq 2$  for HDS. Since this disagreement has not been resolved, the remainder, for purposes of fitting the experimental data, has been arbitrarily chosen to be  $16(G-F)R_6/H^2$ . Obviously, the constant terms become less and less important as  $J$  and  $K$  become large, i.e., for  $J=12, K=6$ , the three  $J$  and  $K$  dependent terms account for over 99 percent of the centrifugal distortion correction.

The results of the equation fitting are summarized in Tables II and III.

The data presented in Table IV is quite consistent with no apparent systematic errors. For this reason, it is felt that the fitting is quite good and that the constants given in Table III truly represent the centrifugal distortion parameters of HDS.

### III. CALCULATION OF THE INERTIA DEFECT AND THE EFFECTIVE ROTATIONAL CONSTANTS

The study of the rotational absorption spectrum of HDS yielded accurate values of  $(b-c)$  and  $\frac{1}{2}(a-c)$ , but to obtain the effective rotational constants  $a$ ,  $b$ , and  $c$  another equation is required. If a  $\Delta J = +1$  transition were observed and corrected for centrifugal distortion, the problem would be solved since  $(a+c)$  would then be known. In the absence of this experimental data, one must resort to other means. Fortunately, another method is available because HDS is a planar molecule

and hence the effective moments of inertia can be expressed in terms of the inertia defect  $\Delta$ , as follows:

$$I_{aa}v + I_{bb}v - I_{cc}v = \Delta. \quad (20)$$

With a knowledge of the molecular force constants and vibrational frequencies obtained from infrared data,  $\Delta$  can be computed with a fair degree of accuracy. Since  $\Delta$  is a small correction term amounting to only a few percent of the rotational constants, an error as large as ten to twenty percent may be permissible in the calculation of  $\Delta$  and not affect the final results more than one percent. In addition, Darling and Dennison<sup>9</sup> and Nielsen<sup>2</sup> have pointed out that  $\Delta$  is independent to second-order approximation of the anharmonic force constants, and therefore only the harmonic terms in the potential energy function need be considered.

The necessary normal coordinates and vibrational frequencies are determined from existing infrared and Raman spectra data. The physical parameters of HDS are taken to be identical with  $H_2S$ ,<sup>10</sup> an assumption which should be quite valid.

On the basis of Slater-Pauling valence forces, the kinetic and potential energy functions can be expressed directly in terms of the directed valence coordinates  $\delta r_1$ ,  $\delta r_2$ , and  $\delta \gamma$  which refer respectively to small displacements from equilibrium positions in the sulfur-hydrogen bond length, the sulfur-deuterium bond length, and the angle  $\gamma$  between the bonds as pictured in Fig. 1. The energy equations are assumed purely quadratic functions in the absence of knowledge concerning anharmonic terms.

The three lagrangian equations of motion can be written immediately, and when solutions of the form  $e^{i\lambda t}$  are assumed, the following secular equation is formed,

$$\begin{vmatrix} a_{11}\lambda^2 - k_{11} & a_{12}\lambda^2 - k_{12} & a_{13}\lambda^2 - k_{13} \\ a_{12}\lambda^2 - k_{12} & a_{22}\lambda^2 - k_{22} & a_{23}\lambda^2 - k_{23} \\ a_{13}\lambda^2 - k_{13} & a_{23}\lambda^2 - k_{23} & a_{33}\lambda^2 - k_{33} \end{vmatrix} = 0, \quad (21)$$

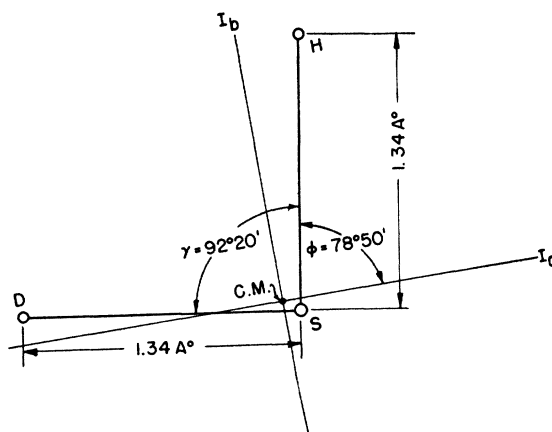


FIG. 1. HDS molecule.

<sup>9</sup> B. T. Darling and D. M. Dennison, Phys. Rev. **57**, 128 (1940).

<sup>10</sup> P. C. Cross and G. W. King, private communication.

TABLE IV. Force constants and normal vibration frequencies of HDS<sup>22</sup>.

Force constants	Normal frequencies
$k_{11} = 4.25 \times 10^6$ dynes/cm	$\nu_1 = 1973$ cm <sup>-1</sup>
$k_{22} = 4.33 \times 10^6$ dynes/cm	$\nu_2 = 1122$ cm <sup>-1</sup>
$k_{33} = 0.238 \times 10^6$ dynes/cm	$\nu_3 = 2723$ cm <sup>-1</sup>
$k_{13} = 0.04 \times 10^6$ dynes/cm	
$k_{12} = 0.05 \times 10^6$ dynes/cm	

where the  $\lambda$ 's are the normal frequencies ( $\lambda^2 = 4\pi^2 c^2 \nu^2$ ), the  $k$ 's are force constants and the  $a$ 's, reduced masses. The reduced masses are given explicitly by Cross and Van Vleck<sup>11</sup> for the general triatomic nonlinear molecule. The force constants were assumed to be known from the work of Bailey, Thompson, and Hale<sup>12</sup> on H<sub>2</sub>S and D<sub>2</sub>S, thus the normal frequencies are determined for HDS. This assumption is quite valid in view of the fact that a comparison of the force constants of H<sub>2</sub>S and D<sub>2</sub>S show differences of the order of magnitude of only one or two percent. The results seem quite reasonable considering the questionable infrared data on HDS. The force constants and zero-order frequencies are given in Table IV.

It is interesting to note that the off-diagonal terms of Eq. (21) are quite small, and thus the actual valence coordinates chosen are essentially the normal coordinates with minor corrections to be made because of the off-diagonal terms. The form of the normal vibrations compared with the original coordinates may be obtained directly from Eq. (21) and will be in the ratio of the first minors of any row in the determinant for each particular normal mode of oscillation. These ratios turn out to be

$$\begin{aligned} \text{for } \nu_1, \quad q_1 : q_2 : q_3 &= -1 : 49 : -2.1 \\ \text{for } \nu_2, \quad q_1 : q_2 : q_3 &= -1.24 : -1 : 150 \\ \text{for } \nu_3, \quad q_1 : q_2 : q_3 &= 109 : 1 : -2, \end{aligned} \quad (22)$$

where  $q_1 \equiv \pi r_1$ ,  $q_2 \equiv \Delta r_2$  and  $q_3 \equiv \sqrt{2} r \Delta \gamma$ . It is immediately obvious that the off-diagonal terms may be neglected with very small effect on the ultimate calculation of  $\Delta$ , and hence the original coordinates are taken to be the

TABLE V. Comparison of calculated and experimental values of the distortion coefficients of HDS<sup>22</sup>.

Distortion coefficient	Calculated value (cm <sup>-1</sup> )	Experimental value (cm <sup>-1</sup> )
$R_5$	$-2.21 \times 10^{-4}$	$-2.808 \times 10^{-4}$
$\delta_f$	$+2.69 \times 10^{-5}$	$+2.531 \times 10^{-5}$
$D_{JK}$	$+9.16 \times 10^{-4}$	$+9.133 \times 10^{-4}$
$D_K$	$-1.64 \times 10^{-4}$	$(-1.64 \times 10^{-4})^*$
$R_6$	$-1.90 \times 10^{-5}$	$+0.695 \times 10^{-5}$
$D_J$	$+8.28 \times 10^{-5}$	

\* Assumed.

<sup>11</sup> P. C. Cross and J. H. Van Vleck, J. Chem. Phys. **1**, 357 (1933).

<sup>12</sup> Bailey, Thompson, and Hale, J. Chem. Phys. **4**, 625 (1936).

normal coordinates. This is quite reasonable from a physical point of view; there are two stretching vibrations and one bending vibration. Each stretching vibration involves only one of the lighter atoms because of the large mass of the sulfur atom and the complete lack of symmetry in the molecule.

Expressing the normal coordinates in Nielsen's notation,  $\Delta$  for the ground vibrational state is computed to be

$$\Delta = -0.244 \times 10^{-40} \text{ g-cm}^2. \quad (23)$$

With  $\Delta$  and the precise values of  $(b-c)$  and  $\frac{1}{2}(a-c)$ , the effective rotational constants of HDS are determined to be

$$\begin{aligned} a &= 9.682 \text{ cm}^{-1}; \quad I_{aa}^v \equiv 2.890 \times 10^{-40} \text{ g-cm}^2 \\ b &= 4.844 \text{ cm}^{-1}; \quad I_{bb}^v \equiv 5.777 \times 10^{-40} \text{ g-cm}^2 \\ c &= 3.140 \text{ cm}^{-1}; \quad I_{cc}^v \equiv 8.910 \times 10^{-40} \text{ g-cm}^2. \end{aligned} \quad (24)$$

These values can now be used in subsequent work, primarily to predict a  $\Delta J = +1$  transition for the microwave spectroscopist and thus the complete experimental determination of the effective rotational constants; and to present considerable aid to the infrared spectroscopist in the analysis of his data. It would be

TABLE VI. Dipole moments.

$\mu(2_{2,1} \rightarrow 2_{2,0})$	$= 1.020$ debyes
$\mu(4_{3,2} \rightarrow 4_{3,1})$	$= 1.017$ debyes
$\mu(7_{4,4} \rightarrow 7_{4,3})$	$= 1.015$ debyes
$\mu(9_{5,5} \rightarrow 9_{5,4})$	$= 1.030$ debyes

eloquent to complete this work by attempting to predict or calculate the equilibrium values of the moments of inertia, but the lack of knowledge of the anharmonic force constants, which give rise to the effects of the same order of magnitude as the effect of harmonic constants, make this task impossible with any degree of accuracy.

#### IV. DISTORTION COEFFICIENTS

The experimental distortion coefficients for HDS are given in Table II. These same coefficients can be calculated utilizing exactly the same data that was needed for the computation of  $\Delta$ . Hence, the validity of the centrifugal distortion theory, of the calculated effective rotational constants and of the structure of HDS can be quantitatively checked by computing these coefficients. The results of these calculations, compared with the experimental values of the distortion constants, are given in Table V.

While the results are not perfect, the excellent agreement between the calculated and measured values of the distortion constants exceeds expectations in view of the many simplifying assumptions made in the course of the previous work, and completes the cycle providing a check on the analysis presented in this paper.

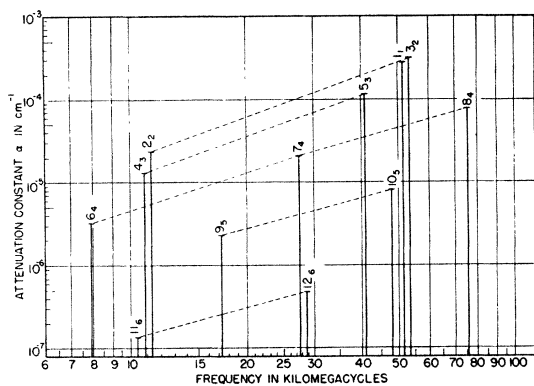


FIG. 2. HDS spectrum.

### V. STARK EFFECT ON HDS

The Stark effect was measured for four HDS lines, the  $2_{2,1} \rightarrow 2_{2,0}$ ;  $4_{3,2} \rightarrow 4_{3,1}$ ;  $7_{4,4} \rightarrow 7_{4,3}$  and the  $9_{5,5} \rightarrow 9_{5,4}$ . The dipole moment of each transition was computed, with the results shown in Table VI. In all cases the Stark splitting follows a  $BM^2$  law, where  $B$  is a constant depending on the transition. The contributions from terms involving states other than the two participating in the absorption transition cancel to less than 1 percent of the major term value for all transitions except that for  $J=2$  where the contribution amounts to 1.2 percent. The latter correction is due entirely to the state  $J=2_{1,1}$  mixing with the  $J=2_{2,0}$  state.

Within the accuracy with which the measurements were made, the electric dipole moment of HDS is quoted to be

$$\mu = 1.02 \pm 0.02 \text{ debyes.}$$

It was originally thought that a correlation would be detected between a uniform change in  $\mu$  and increasing  $J$ , i.e., increase of centrifugal distortion of the molecule. Such a change would indicate, on the naive picture of a rotating molecule, whether the hydrogen and deuterium were approaching or receding from each other. Unfortunately, the results are not accurate enough to discern such a trend. From an angular momentum point of view, i.e.,  $\langle P_i^2 \rangle = \partial E / \partial i$ , it would seem that with increasing energy the dipole moment should increase.

### VI. ABSORPTION COEFFICIENTS AND ISOTOPIC SULFUR TRANSITIONS

The absorption coefficients in  $\text{cm}^{-1}$  of  $\text{HDS}^{32}$  absorption transitions are illustrated graphically in Fig. 2 as

TABLE VII. Absorption frequencies of  $\text{HDS}^{34}$  and  $\text{HDS}^{33}$ .

Compound	Transition	Frequency (Mc/sec)
$\text{HDS}^{34}$	$1_{1,1} \rightarrow 1_{1,0}$	50,912.27
$\text{HDS}^{34}$	$2_{2,1} \rightarrow 2_{2,0}$	11,235.45
$\text{HDS}^{34}$	$3_{2,2} \rightarrow 3_{2,1}$	52,979.67
$\text{HDS}^{34}$	$4_{3,2} \rightarrow 4_{3,1}$	10,802.36
$\text{HDS}^{34}$	$7_{4,4} \rightarrow 7_{4,3}$	27,392.00
$\text{HDS}^{33}$	$2_{2,1} \rightarrow 2_{2,0}$	11,258.21

a function of the frequency, in which form it serves several useful purposes. One, it presents all pertinent experimental data on HDS to the observer at a glance; two, experimentalists in the field may find the data useful as a standard; and three, qualitative predictions can be made regarding additional absorption transitions in HDS by observing the trends of the  $K$  families as indicated by the light dotted lines on the graph.

Several absorption lines were observed and measured for the isotopic compounds  $\text{HDS}^{34}$  and  $\text{HDS}^{33}$ , and these are tabulated in Table VII for future reference. The intensities of these lines are the same as shown in Fig. 2 but corrected for the natural abundance of the sulfur isotopes. The  $\text{HDS}^{33}$  line is the weakest line observed ( $\alpha \approx 10^{-7} \text{ cm}^{-1}$ ) and was the largest of three lines dispersed over a region of approximately 10 Mc/sec. The poor sensitivity is not attributed to faulty equipment, but to a poor sample of HDS.

### VII. CONCLUSION

Within certain restrictions a formula for centrifugal distortion in asymmetric molecules has been presented. The distortion is determined by five calculable constants. As a test of the formula it has been applied to the observed microwave spectrum of HDS. The agreement is sufficiently satisfactory to confirm the general functional dependence of the formula on rotational state. Disagreement does exist, however, between the computed and observed effect both in magnitude and dependence of the contribution resulting from the fourth off-diagonal element, ( $R_6$ ). Fortunately the effect of this term is minor.

The formula given above compares with the approximate one derived in reference 3, except for a slight difference in the dependence on  $K$ . Further discussion of the comparison of the two formulas will be given in paper III of this series.