

Fig. 1. Hysteresis loops for $\mathrm{BaTiO}_{3}$ at low temperatures.
polarization. This stays almost constant from the $180^{\circ} \mathrm{K}$ transition down to $4.2^{\circ} \mathrm{K}$; its value is of the order of about $P_{s} \simeq 8.0 \times 10^{-6}$ coulomb $/ \mathrm{cm}^{2}$. The absolute value changes from sample to sample because of domain formation at the two transition points ${ }^{3}$ near $270^{\circ}$ and $180^{\circ} \mathrm{K}$. In interpreting $P_{s}$ it must be kept in mind that below the $180^{\circ} \mathrm{K}$ transition the polar axis points in the (111) direction. ${ }^{3}$ The coercive field strength increases strongly from about $250 \mathrm{v} / \mathrm{cm}$ at $180^{\circ} \mathrm{K}$ to almost $10,000 \mathrm{v} / \mathrm{cm}$ at $4.2^{\circ} \mathrm{K}$ (Fig. 2).

These results show that $\mathrm{BaTiO}_{3}$ stays ferroelectric down to liquid helium temperature, and there is no indication that the crystal will change to a nonferroelectric modification. The fact


Fig. 2. Spontaneous polarization and coercive field strength os temperature.
that the coercive field strength increases strongly at low temperatures and that the hysteresis loops become wider and more and more rectangular indicates freezing-in of the domain boundaries. It is of interest to note that $\mathrm{BaTiO}_{3}$, while becoming ferroelectric by a displacive transition similarly to rochelle salt (small specific heat anomaly) and, in contrast to the $\mathrm{KH}_{2} \mathrm{PO}_{4}$ group, shows no second curie point at which the ferroelectricity disappears as in the case of rochelle salt.
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## The Intensity of the Total and of the Hard Component of the Cosmic Radiation as a Function of Altitude at Geomagnetic Latitudes of $28^{\circ} \mathrm{N}$ and $55^{\circ} \mathrm{N}^{*}$

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THE equipment used in this experiment consisted of four fourfold coincidence counter telescopes $(A B C D, B C D E$, $B C D F, B C D G)$. The counter geometry of the individual telescopes is shown in Fig. 1. Each counter had an outside diameter of 2.54


Fig. 1. Measurements at $28^{\circ} \mathrm{N}$ geomagnetic latitude (Cuba).
cm , a length of 13 cm , and brass walls of 0.078 cm in thickness. A block of 12 cm of lead was interposed between counter $D$ and counters $E F G$. Each fourfold coincidence was separately registered on photographic film moved by a clock mechanism.

Through the courtesy of the Office of Naval Research, the first of the two balloon experiments was carried out with the aid of a General Mills balloon launched south of Cuba at a geomagnetic latitude of $28^{\circ} \mathrm{N}$. The launching took place from the deck of an
aircraft carrier of the U. S. Navy on November 19, 1949. The balloon rose at the rate of $680 \mathrm{ft} / \mathrm{min}$ to a maximum altitude of 93,000 feet corresponding to an atmospheric pressure of 1.2 cm Hg and remained between a pressure of 1.2 and 1.4 cm Hg for several hours. The equipment landed in the ocean and was promptly recovered by a helicopter.

Curves A and B in Fig. 1 give the counting rates obtained during this flight for the fourfold coincidences $A B C D$ and $B C D F$, respectively. Curve A is assumed to give the total vertical flux of the cosmic radiation at $28^{\circ} \mathrm{N}$. The maximum in the total radiation occurs at a pressure of 9 cm Hg and represents an increase by a factor of 18 from its value at sea level. Curve B gives the intensity of the cosmic radiation capable of traversing 12 cm of lead (penetrating component). As is evident in Fig. 1, curve B flattens out appreciably above 18 cm Hg pressure and shows a definite drop between 5 and 1 cm Hg . The penetrating component increases by a factor of 8 from its value at sea level to that at the very flat maximum ( $\sim 9 \mathrm{~cm} \mathrm{Hg}$ ).
The second experiment, in which an identical piece of apparatus was used, was carried out at a geomagnetic latitude of $55^{\circ} \mathrm{N}$. For this purpose a General Mills balloon was launched from Minneapolis, Minnesota, on June 6, 1950. The balloon rose at the very slow rate of $311 \mathrm{ft} / \mathrm{min}$ to an altitude of $70,000 \mathrm{ft}$ and continued to rise even more slowly ( $\sim 150 \mathrm{ft} / \mathrm{min}$ ) until it leveled off for several hours at an altitude of $86,000 \mathrm{ft}$. This type of flight is particularly valuable for securing accurate counter coincidence data at various atmospheric depths.

The results obtained are shown in Fig. 2. The maximum in the total component occurs at an atmospheric pressure of 5 to 6 cm


FIG. 2. Measurements at $55^{\circ} \mathrm{N}$ geomagnetic latitude (Minneapolis, Minnesota).

Hg and represents an increase by a factor of 35 from its value at sea level. The hard component, as measured through 12 cm of lead, flattens out above 5 cm Hg , exhibiting a small maximum at a pressure of about 3 cm Hg . The increase from sea level to the maximum corresponds to a factor of 19.

Comparison of the results obtained at the two latitudes gives a latitude effect of $2.75 \pm 0.21$ for the hard component at an atmospheric pressure of 1.6 cm Hg . For the total radiation the latitude effect is 1.85 at 9 cm Hg (maximum at $28^{\circ} \mathrm{N}$ ), 3.1 at 5 cm Hg (maximum at $55^{\circ} \mathrm{N}$ ), and $3.85 \pm 0.13$ at 1.6 cm Hg .

It is of interest to note that the latitude effect of the hard component is smaller than that of the total radiation at the maximum altitude of $85,000 \mathrm{ft}$.
Winckler ${ }^{1}$ has measured the latitude effect of the total radiation
in a series of balloon flights and obtains results in good agreement with the ones presented here. However, an accurate comparison of the data is not possible owing to the different counter geometry and the thickness of lead in the counter telescope used in his equipment.
A more detailed discussion of the results will be published later.
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* Assisted by the joint program of the ONR and AEC.
${ }^{1}$ Winckler, Stix, Dwight, and Sabin, Phys. Rev. 79, $656^{\circ}$ (1950).


## An $8 \times 10^{-10}-$ Sec Isomeric State in ${ }_{76} \mathrm{Os}^{186}$ <br> F. K. McGowan <br> Oak Ridge National Laboratory,* Oak Ridge, Tennessee January 19, 1951

AN excited state in $\mathrm{rsO}^{186}$ with a half-life $(8 \pm 1) \times 10^{-10} \mathrm{sec}$ has been observed with a delayed coincidence scintillation spectrometer. The data have been analyzed by a graphical method which is analogous to a method of analysis recently reported by Newton. ${ }^{1}$
$\operatorname{Re}^{186}(90 \mathrm{hr})$ is known ${ }^{2}$ to decay in 90 percent of its disintegrations by $\beta^{-}$emission into $\mathrm{Os}^{186}$. A 25 percent partial $\beta^{-}$ spectrum of $930-\mathrm{kev}$ end point is followed by a $137-\mathrm{kev} \gamma$-ray, and the remaining $\beta^{-}$disintegrations lead directly to the ground state. For the $137-\mathrm{kev} \gamma$-ray, Metzger and Hill measured $\alpha_{K}=0.35$ $\pm 0.1$ and $N_{K} / N_{L}=0.6 \pm 0.1$.
Curve (1) of Fig. 1 shows the number of coincidences as a function of delay time obtained with a source of $\mathrm{Re}^{186}$. This delayed


Fig. 1. The number of delayed coincidences as a function of delay time.
coincidence resolution curve was recorded by exciting one channel of the delayed coincidence apparatus by 250 - to $450-\mathrm{kev}$ nuclear beta-rays and the other channel by the $L$ and $M$ internal conversion electrons of the $137-\mathrm{kev}$ transition. Without a change in the apparatus, a resolution curve for prompt events was obtained with a source of $\mathrm{Au}^{198}$. The prompt coincidences are (a) between 110 - to 150 -kev nuclear beta-rays and conversion electrons of the $411-\mathrm{kev}$ transition and (b) between 110 - to $150-\mathrm{kev}$ Compton recoil electrons from $411-\mathrm{kev} \gamma$-radiation and 250 - to $450-\mathrm{kev}$ nuclear beta-rays. Curve (2) shows the result of such a measurement. Since curves (1) and (2) overlap appreciably, the half-life of

