## Letters to the Editor

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## A Method of Approximation for **Cooperative Phenomena**

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NEW method for approximating the entropy of an order-A disorder system having an assigned energy has been developed. By a suitable choice of variables it gives the Bethe-Takagi approximation<sup>1,2</sup> in one case, and the "variation method" introduced by Kramers and Wannier<sup>3</sup> in another. The new method is more intuitive than the latter and is more effective than the quasi-chemical method by Yin-Yuan Li.<sup>4</sup> It can be improved and extended successively.

Two-dimensional square net.- The probability of appearance of a configuration of a cell (a bond, a spin) given on the left column of each part of Fig. 1 is denoted by  $z_{\nu}$ ,  $y_{\nu}$ , or  $x_{\nu}$  on the corresponding second column.  $\gamma_{\nu}(\beta_{\nu}, \alpha_{\nu})$  denotes the number of configurations with the same probability  $z_{\nu}$ ,  $y_{\nu}$ , or  $x_{\nu}$ . Modifying the method of Takagi<sup>2</sup> for approximating the number of ways of constructing a crystal, the entropy of the system composed of M spins each located at a square net point can be approximated as:

$$S = kM \left[ 2 \sum_{\nu=1}^{3} \beta_{\nu} y_{\nu} \ln y_{\nu} - \sum_{\nu=1}^{2} \alpha_{\nu} x_{\nu} \ln x_{\nu} - \sum_{\nu=1}^{6} \gamma_{\nu} z_{\nu} \ln z_{\nu} \right].$$

The energy E is rigorously expressed as

$$E = \frac{1}{2} 4M \sum_{\nu=1}^{3} \epsilon_{\nu} \beta_{\nu} y_{\nu},$$

where  $\epsilon_{\nu}$  is the energy of the bond. Minimizing the free energy F = E - TS, we get results identical with those obtained by the "variation method" by Kramers and Wannier.3

As a by-product it was found that the Kramers-Wannier result<sup>3</sup> for the ordered phase can be simplified as follows: Their Eq. (91.a) can be factorized as

Cell	Prob.	$\chi_{\nu}$	
⊕⊕ ⊕⊕	Ζ,	1	
⊕⊕ e+c ⊕⊖ e+c	$Z_{z}$	4	
⊕⊖ ⊕⊖ etc	Z,	4	
⊖⊕ ⊕-⊖ etc	Z,	2	
Ф ен ен	$Z_{s}$	4	
⊖0 ⊝0	Z	1	
(0)			

Bond	Prob.	$\beta_{\nu}$	ε,
••	У,	1	-8
⊕© ⊙€	У2	<i>t</i> 2	+8
ΘΘ	У3	1	-£
	(6)		

Prob. Spin  $\alpha_{\nu}$ 1 x,  $\bigcirc$ 1 X2 (C)

FIG. 1. Configuration scheme.

$$\{k^2z^2+(k^4+1)z+k^6\}\{k^4z^2-k^2(k^4-3)z+1\}=0.$$

The true solution is contained in the second factor, which gives  $\lambda = \{k^{8} + 8k^{4} - 11 + (k^{4} - 5)\frac{1}{2}(k^{4} - 1)\frac{1}{2}\}/(2k^{6})$ 

in place of their Eq. (91.b).

C

Taking  $y_{\nu}$  and  $x_{\nu}$  as variables, we can approximate the entropy as

$$S = kM [3 \sum_{\nu=1}^{2} \alpha_{\nu} x_{\nu} \ln x_{\nu} - 2 \sum_{\nu=1}^{3} \beta_{\nu} y_{\nu} \ln y_{\nu}]$$

which gives a result identical with Takagi's<sup>2</sup> and Bethe's.<sup>1</sup>

Simple cubic lattice.—Taking  $z_{\nu}$ ,  $y_{\nu}$ , and  $x_{\nu}$  as variables, the entropy of a system with M spins, each located at a simple cubic lattice point, can be approximated as

$$S = kM \left[9 \sum_{\nu=1}^{3} \beta_{\nu} y_{\nu} \ln y_{\nu} - 7 \sum_{\nu=1}^{2} \alpha_{\nu} x_{\nu} \ln x_{\nu} - 3 \sum_{\nu=1}^{6} \gamma_{\nu} z_{\nu} \ln z_{\nu}\right]$$

Minimizing the free energy, we get the following results:

(a). The transition is of the second order and the transition temperature  $T_t$  is

$$kT_t/\epsilon = 4.6097$$

(b). The specific heat  $c_{t+}$  and  $c_{t-}$  on the higher and the lower temperature side respectively at  $T_t$  is

$$_{+}/k = 0.389$$
 and  $c_{t-}/k = 2.902$ .

(c). Extrapolated, the specific heat for the disordered phase becomes infinite at

$$kT/\epsilon = 4.2221$$

Therefore it might not be in error to infer that the correct transisition point  $T_c$  would be in the range

$$4.2221 < kT_c/\epsilon < 4.6097$$

When we take the probabilities of configurations of a cubic cell as variables, we shall have a better approximation, which will be identical with the extension of the "variation method" to the simple cubic case tried by ter Haar and Martin.<sup>5</sup>

Face-centered cubic lattice (Ising model).-Taking as variables the probabilities of configurations of a tetrahedron, we reach the following results:

(a). The transition is of the second order and

$$kT_t/\epsilon = 10.0260.$$

(b). 
$$c_{t+}/\epsilon = 0.336$$
 and  $c_{t-}/\epsilon = 2.789$ 

(c). By the same inference as above we would obtain the following range for the correct transition point,  $T_c$ ,

 $9.2388 < kT_c/\epsilon < 10.0260.$ 

A more detailed account will be published elsewhere.

<sup>1</sup> H. A. Bethe, Proc. Roy. Soc. A150, 552 (1935).
<sup>2</sup> Y. Takagi, Proc. Phys. Math. Soc. Japan 23, 44 (1941).
<sup>3</sup> H. A. Kramers and G. H. Wannier, Phys. Rev. 60, 252 and 263 (1941).
<sup>4</sup> Yin-Yuan Li, J. Chem. Phys. 17, 447 (1949); Phys. Rev. 76, 972 (1949).
<sup>5</sup> D. ter Haar and B. Martin, Phys. Rev. 77, 721 (1950).

## Cosmic Rays Underground

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HE variation of cosmic-ray intensity underground has been studied recently by several authors.<sup>1,2</sup> In this note are reported results of calculations of the absorption of  $\pi$ - and  $\mu$ -mesons underground, performed taking into account the process of creation of  $\pi$ -mesons by direct electromagnetic interaction of charged