viously used for the mass determination of active isotopes.4 The instrument and techniques employed in this mass analysis were the same as those used in the measurement of the isotopic constitution of lanthanum and cerium.⁵ Two different samples of Yb₂O₃ were analyzed. The first was secured from Adam Hilger, Ltd. (laboratory number 1071), while the second was an exchange column purified sample obtained from Dr. D. H. Harris of the Clinton Laboratories. There was no detectable difference in the isotopic constitution of the two samples. The values quoted in the third row of Table I are the averages of thirty separate determinations on each sample. The mean deviation of each percentage was about 1 percent of that percentage. Because of the large number of readings the precision should be better than this by a factor of one over the square root of the number of determinations. However, since we have not as yet been able to rule out systematic mass discriminations because of selective emission from the surface ionization source, or because of non-linearity of the 1010-ohm resistor over the wide current range used, to better than 1 percent we do not wish to quote any limits closer than this.

The following upper limits for the natural occurrence of other isotopes of ytterbium were set: $Yb^{166} < 0.002$ percent, Yb167 < 0.002 percent, Yb169 < 0.01 percent, Yb175 < 0.02 percent, Yb¹⁷⁷<0.01 percent, Yb¹⁷⁸<0.002 percent.

With the assumption of zero packing fraction for the ytterbium isotopes the chemical atomic weight of ytterbium calculated from the new abundances is 173.046 ± 0.006 . The international value, chemically determined, is 173.04.

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A Measurement of the Half-Life of Double Beta-Decay from 50 Sn¹²⁴ *

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 \mathbf{I}^{F} two isobars differ by two units in atomic number, the heavier may decay into the lighter by double betadecay.1,2 This is the simultaneous emission of two negatrons if the heavier has lower atomic number or the simultaneous emission of two positons, 1 positon +1K capture, or 2Kcaptures if the heavier has higher atomic number. The half-life depends markedly upon whether or not two neutrinos are emitted in the process. If no neutrinos are

TABLE I. Theoretical half-life for allowed double negaton emission.

Atomic mass difference	0	0.52 Mev	1.04 Mev	1.56 Mev	2.08 Mev	2.60 Mev
2 neutrinos	œ	2.6 · 10 ²⁷ yr.	2.4 · 10 ²⁵ yr.	1.3 · 10 ²⁴ yr.	2.1 ·10 ²³ yr.	4.3 · 10 ²² yr
No neutrinos	8	2.1 · 10 ¹⁶ yr.	2.7 ·10 ¹⁵ yr.	6.5·10 ¹⁴ yr.	2.2·10 ¹⁴ yr.	8.3 · 10 ¹³ yr
-						



FIG. 1. Experimental arrangement.

emitted, the half-life is of the order of 1010 times shorter than if two neutrinos are emitted (see Table I). The reason for this large difference in half-life arises from the number of cells in phase space available for the transition.

In the present work an experimental investigation of $_{50}$ Sn¹²⁴ which belongs to isobaric triplet $_{50}$ Sn¹²⁴ - $_{52}$ Te¹²⁴ - 54Xe¹²⁴ has been carried out. The experimental arrangement consists of four thin window counters (3 mg/cm² mica) connected in pairs to coincidence circuits L and R. These counters are shielded by anticoincidence counters and by an Fe and Pb box. Two specimens** of Sn (25 g) identical in all respects except for isotopic constitution are placed between the thin window counters. These specimens are called A and B. Specimen A contains 54 percent of Sn¹²⁴; specimen B contains 0.4 percent of Sn¹²⁴. The position of the specimens between the counters is interchanged by rotating the specimen holder through 180° (see Fig. 1).

Coincidences and single counts from both specimens are recorded simultaneously. The specimen holder is rotated through 180° every other hour and the positions of the specimens in the holder are interchanged every 20 hours. These data are summarized in Table II.

In all situations specimen A gives 2 coincidence counts/ hr. more than specimen B. By repeating this type of measurement with Al absorbers over one side of each specimen an absorption curve is obtained. This absorption curve is similar to that of electrons from a spectrum with an energy end point between 1.0 Mev and 1.5 Mev. The single counts from specimens A and B both give 6.5 ± 0.3 counts/min. If one interprets this effect as double betadecay from Sn^{124} , one obtains a half-life between $0.4 \cdot 10^{16}$ vr. and $0.9 \cdot 10^{16}$ yr. Other alternative explanations for these observations have been considered but none have been found to be plausible. This result would indicate that double beta-decay is unaccompanied by neutrinos. A further consequence of these results pointed out to the author by Professor J. R. Oppenheimer is that the neutronproton charge difference is exactly equal to the electron charge,

TABLE II. L(A) gives the coincidences from specimen A between counters L, and R(B) gives the coincidences from B between counters R. Holder 0° and 180° are the two holder positions. Positions 1 and 2 are positions for the specimens in the holder.

Pos. 1	Holder 0° Coin. counts/hr.	L(A) 16.4 ± 0.3	R(B) 14.3 ±0.3
	Holder 180° Coin. counts/hr.	L(B) 14.4 ± 0.3	R(A) 15.9 ±0.3
Pos. 2	Holder 0° Coin. counts/hr.	L(B) 14.6 ±0.3	R(A) 16.4 ±0.3
	Holder 180° Coin. counts/hr.	L(A) 16.4 ±0.3	$\begin{array}{c} R(B) \\ 13.9 \pm 0.3 \end{array}$

A detailed report of this work is being prepared for publication in the Physical Review.

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** These isotopes were obtained from Oak Ridge.
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Internal Conversion Electrons from Metastable **Te**¹²⁵

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R ECENTLY, it has been found¹ that an isomer of the stable isotope Te^{125} with a half-life of two months stable isotope Te¹²⁵ with a half-life of two months



grows out of Sb125 (2.7 yr.). Absorption measurements of its radiations showed that it decays through a highly con-

Line Βρ gauss-cm energy kev Gamma-ray energy kev Relative intensity Assign ment 77.5 105.1 77.5 + 31.8 = 109.3105.1 + 4.35 = 109.45108.6 + 0.58 = 109.18975.2 1147 100 $\gamma_1 - L$ $\gamma_1 - M$ 67 19 1169 108.6 598.6 30.4

TABLE I. Average relative intensities in lines.

verted isomeric transition of approximately 120-kev energy.

The decay of Te¹²⁵ has now been followed for a longer period and the half-life was found to be 58 ± 4 days. For a more accurate measurement of the energy of the isomeric transition, a source of a few microcuries of Te125 was separated from Sn irradiated with slow neutrons at Oak Ridge. It was subjected to an analysis in a 180° magnetic beta-ray spectrograph previously described.² Three lines corresponding to the K, L, and M conversion electrons of a 109.3-kev gamma-ray were obtained after a 36-hr. exposure. A longer exposure of about one week yielded, in addition, a weak line produced by conversion electrons of 30.4 key, for which we cannot state the gamma-ray energy since we do not know in which shell the conversion took place. Figure 1 shows the spectrum of the conversion lines. A part of a microphotometer trace of the film is reproduced in Fig. 2. Using the density-intensity calibration recently described,3 we obtained the relative intensities of the lines from a number of microphotometer traces. The average relative intensities are given in Table I. It follows that the K/L intensity ratio is ~1.5, the L/M ratio ~3.5. From the half-life and energy, and from the lower limit of 0.99 for the total internal conversion coefficient, it was previously concluded¹ that the effective $\Delta l = 5$ for this transition. According to Drell's⁴ recent calculations, the K/Lratio obtained by us is compatible with an effective $\Delta l = 5$, if one assumes that this transition is "parity forbidden," and that 95 percent is 24-pole magnetic while the rest is 25-pole electric. The spin change between the metastable state and ground state would then be 4. The role of the very weak transition is at present in doubt. Further experiments are needed to decide whether it has to be ascribed to Te^{125*}.



FIG. 2. Microphotometer trace of the internal conversion electron lines from the 109.3-key gamma-ray of Te^{125*}.

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