

Nucleus	Sr ⁹⁰	Sr ⁹¹	Y ⁹⁰	Y ⁹¹	Sb ¹²⁵	Cs ¹³⁷
$(W_0^2 - 1)ft \times 10^{10}$	0.4	0.7	0.3	0.8	1.1	1.8

The criterion fits Sr⁸⁹ and K⁴²; in both cases shell model considerations support the assignment $\Delta I = \pm 2$ and change in parity. An excellent measurement of K⁴² has already been made by Siegbahn.⁹ It disintegrates (Fig. 1) by negatron emission to the ground state of Ca⁴² and also to an excited state followed by a gamma-ray. We have analyzed Siegbahn's data for the high energy component with results shown in Fig. 2 both for an allowed FK plot and again as a forbidden FK plot wherein the factor G is included. The bulging curvature of the one and the near-straightness of the other confirm the above-stated criterion as well as those aspects of the shell model and of the Fermi theory of beta-decay which are involved in this example.

The shell model associates the odd configuration $(3d)^{-1}(4f)^3$ with the ground state of K⁴². Similarly the even configuration $(4f)^2$ may be associated with the ground state of Ca⁴² (probably $I=0$) and with low excited states ($I=2, 4, \text{ and } 6$). We suggest $I=2$ and even parity for the excited state of Ca⁴² observed in the K⁴² decay. The low energy transition is thus interpreted as first-forbidden with $\Delta I=0$.

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¹ E. J. Konopinski and G. E. Uhlenbeck, Phys. Rev. **60**, 308 (1941).

² L. M. Langer and H. C. Price, Jr., Phys. Rev. **75**, 1109 (1949).

³ J. S. Osoba, to be published.

⁴ A. C. G. Mitchell and C. L. Peacock, to be published in Phys. Rev. (private communication).

⁵ C. H. Braden, L. Slack, and F. B. Shull, to be published. Also L. J. Laslett and E. Jensen (private communication).

⁶ A. C. G. Mitchell (private communication); also Zaffarano, Kern, and Mitchell, Phys. Rev. **74**, 682 (1948), particularly Fig. 4.

⁷ E. J. Konopinski, Rev. Mod. Phys. **15**, 209 (1943).

⁸ E. Feenberg and K. Hammack, Phys. Rev., to be published.

⁹ K. Siegbahn, Arkiv. f. Mat., Astr. o. Fys. **34B**, No. 4 (1946).

A total of ten determinations was made. From these the following value for the ratio of the resonant frequencies in the same magnetic field was obtained:

$$\nu(\text{Be}^9)/\nu(\text{H}^1) = 0.1405187 \pm 0.000002.$$

All ten values fall well within the above limits which were determined entirely by the uncertainty in locating the exact centers of the resonance curves.

Applying the Lamb⁴ diamagnetic correction, and using the Millman and Kusch⁵ value of the proton moment (2.7896 ± 0.0008) μ_N , the Be⁹ magnetic moment is found to be

$$\mu(\text{Be}^9) = (-)(1.17619 \pm 0.00034) \mu_N.$$

If, in agreement with more recent experiments, the more accurate value⁶ of $(2.7926 \pm 0.0006) \mu_N$, is taken for the proton moment, the following value is obtained:

$$\mu(\text{Be}^9) = (-)(1.17747 \pm 0.00027) \mu_N.$$

It is to be noted that the uncertainties indicated for the Be⁹ magnetic moment are determined by the uncertainty in the proton moment and not by that of the above frequency ratio.

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¹ Pound, Purcell, and Torrey, Phys. Rev. **69**, 681 (1946).

² F. Bitter, Phys. Rev. **75**, 1326 (1949).

³ H. L. Poss, Phys. Rev. **75**, 600 (1949).

⁴ W. E. Lamb, Jr., Phys. Rev. **60**, 817 (1941).

⁵ S. Millman and P. Kusch, Phys. Rev. **60**, 91 (1941).

⁶ This value was calculated using the absolute value of the proton gyromagnetic ratio given by Thomas, Driscoll, and Hipple, Phys. Rev. **75**, 902 (1949), and the physical constants e , M_p , and c taken from J. W. Dumond and E. R. Cohen, Rev. Mod. Phys. **20**, 82 (1948). The proton moment calculated in this way agrees very well with the value given by Millman and Kusch if the latter is corrected for the magnetic moment of the electron as suggested by J. Schwinger, Phys. Rev. **73**, 416 (1948). However, the small uncertainty in the absolute value of the proton gyromagnetic ratio cannot be carried over to the proton moment in units of the nuclear magneton because of the larger uncertainty in the values of e and M_p .

The Magnetic Moment of Be⁹ *

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THE ratio of the frequency of the nuclear magnetic resonance of Be⁹ to that of the proton has been measured at 7000 gauss by the magnetic-resonance-absorption method of Pound, Purcell, and Torrey.¹ The magnet and radiofrequency bridges employed were those used by both Bitter² and Poss³ in their recent nuclear-magnetic-moment determinations. However, the accuracy of the present measurement exceeds that of previous measurements with this equipment by about a factor of ten. This was accomplished by taking advantage of the almost integral ratio of the two frequencies and heterodyning the seventh harmonic of the Be⁹ frequency with the fundamental of the proton frequency.

Both the proton and Be⁹ resonances were obtained from a single sample consisting of an aqueous solution of beryllium fluoride. The sample container was a cylindrical section $\frac{1}{2}$ -in. long and $\frac{1}{8}$ in. in diameter. The coil for the 4.2-Mc (Be⁹ resonant frequency) bridge was wound around the cylinder, and the coil for the 30-Mc (proton resonant frequency) bridge was wound at right angles over the first. This technique has previously been used by F. Bitter and assures that the external magnetic field seen by the two types of nuclei is the same. The two resonances were traced on separate recording milliammeters as the magnetic field was slowly varied. The frequencies were adjusted so that the resonances occurred nearly simultaneously. From a calibration of the rate of change of magnetic field, a correction was made, when necessary, to take into account the small field differences separating the two resonances. This correction never exceeded 0.003 percent of the observed frequency ratio.

Paramagnetic Resonance Absorption in Crystals Colored by Irradiation

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PARAMAGNETIC resonance absorption at 9350 mc produced by the irradiation of LiF with neutrons has been observed. Single crystals of LiF, which showed no resonance, were irradiated in a flux of approximately 10^{12} neutrons cm^{-2} for periods of time varying from 1 to 24 hrs. After irradiation, these crystals showed the resonance absorption described in Fig. 1, when placed at the midpoint of a resonant section of

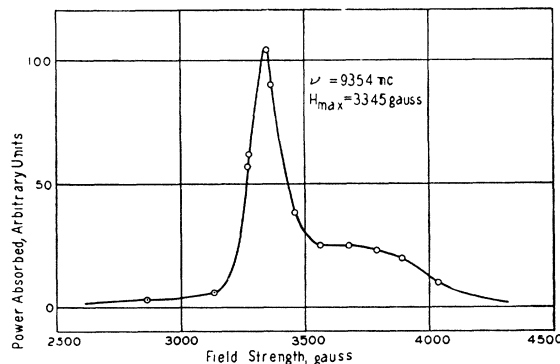


FIG. 1. Resonance absorption of LiF crystals after irradiation with neutrons.