

The time of slowing of the Li particles in FeB is smaller by a factor of about 0.5 than in B<sub>4</sub>C. A further experiment using this compound as a  $\gamma$ -ray source together with an improved spectrometer resolving power may make it possible to give a definite value for the mean life of this excited state in Li<sup>7</sup> rather than only an upper limit.

The angular momentum of Li<sup>7</sup> in the ground state is  $\frac{3}{2}$  and the 478.5-kev excited state is commonly supposed<sup>5</sup> to have angular momentum  $\frac{1}{2}$  forming with the ground state a doublet,  $^2P_{3/2}$  and  $^2P_{1/2}$ . On this basis the transition takes place by magnetic dipole radiation. The transition probability for this case may be calculated closely by considering the interaction of the magnetic moment ( $e/2Mc$ )( $L+5.6S$ ) of the  $2p$  proton in Li<sup>7</sup> with the radiation field. The square of the matrix element<sup>6</sup> representing this interaction is

$$\frac{8}{3} \left( \frac{2.3eh}{4\pi Mc} \right)^2,$$

which together with the energy 478.5 kev available for the transition gives a calculated mean life for the Li<sup>7</sup> excited state of  $1.5 \times 10^{-13}$  sec. This is seen to be consistent with the experimental upper limit of  $2.0 \times 10^{-13}$  sec.

<sup>1</sup> DuMond, Lind, and Watson, Phys. Rev. **73**, 1392 (1948).

<sup>2</sup> Rubin, Snyder, Lauritsen, and Fowler, Bull. Am. Phys. Soc. **23**, No. 5, 15 (1948).

<sup>3</sup> W. Hornyak and T. Lauritsen, Bull. Am. Phys. Soc. **23**, No. 5, 16 (1948).

<sup>4</sup> C. W. Gilbert, Proc. Camb. Phil. Soc. **44**, 447 (1948).

<sup>5</sup> G. Breit and J. R. Stehn, Phys. Rev. **53**, 459 (1938).

<sup>6</sup> We are indebted to Dr. J. A. Spiers for help with this calculation.

## Microwave Determination of the Molecular Structure of Chlorosilane

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THE second rotational transition,  $J=1$  to  $J=2$ , of SiH<sub>3</sub>Cl<sup>35</sup> and SiH<sub>3</sub>Cl<sup>37</sup> has been measured and analyzed. The Hughes-Wilson Stark modulation technique was used for detection of the absorption lines, and absorption frequencies were determined by comparison with a crystal-controlled secondary frequency standard.

In Fig. 1 is shown the theoretical hyperfine pattern for a nuclear spin of  $\frac{3}{2}$  for chlorine which is in satisfactory agreement with the observed spectrum. No lines caused by excited vibrational states were observed. The unperturbed transition frequency  $\nu_0$  is 26,695.24 Mc for SiH<sub>3</sub>Cl<sup>35</sup> and 26,049.60 Mc for SiH<sub>3</sub>Cl<sup>37</sup>. From these frequencies, the  $I_B$  and  $B_0$  values were calculated and listed in Table I. Since the quadrupole moments of Cl<sup>35</sup> and Cl<sup>37</sup> have been measured to be  $-7.921 \pm 0.05 \times 10^{-26}$  and  $-6.189 \pm 0.05 \times 10^{-26}$  cm<sup>2</sup>,<sup>1</sup> respectively, we may evaluate  $\partial^2 V / \partial Z^2$  at

TABLE I. Nuclear and molecular constants of chlorosilane.

	$\nu_0$ (Mc/sec.)	$I_B$ (g-cm <sup>2</sup> ) $\times 10^{40}$	$eQ \left( \frac{\partial^2 V}{\partial Z^2} \right)$ (Mc)	$B_0$ (cm <sup>-1</sup> )
SiH <sub>3</sub> Cl <sup>35</sup>	26,695.24	125.7 <sub>1</sub>	-40. <sub>0</sub>	0.2226 <sub>2</sub>
SiH <sub>3</sub> Cl <sup>37</sup>	26,049.60	128.8 <sub>2</sub>	-30. <sub>8</sub>	0.2172 <sub>2</sub>

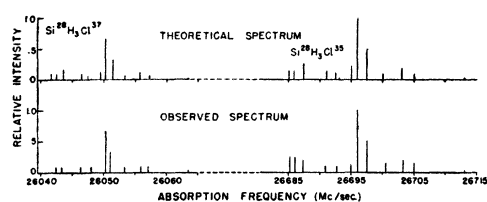


FIG. 1. Microwave absorption spectrum of monochlorosilane (SiH<sub>3</sub>Cl).

the chlorine nucleus of chlorosilane to be about  $7 \times 10^{15}$  e.s.u. This is to be compared with the value  $13 \times 10^{15}$  e.s.u. as determined by Gordy<sup>2</sup> and his co-workers for methyl chloride.

Assuming the value of 1.456Å for the Si-H distance determined from infra-red measurements on silane,<sup>3</sup> we may calculate from the  $I_B$  values in Table I, the Si-Cl internuclear distance to be 2.035Å, and the H-Si-H angle to be  $103^\circ 57'$ . These are to be compared with the value 2.16Å for Si-Cl computed from Pauling's covalent single-bond radii and the tetrahedral angle of  $109^\circ 28'$ . Electron diffraction measurements<sup>4</sup> on this compound yield a value of  $2.06 \pm 0.05$ Å for the Si-Cl distance.

The lines arising from  $K=0 \rightarrow 0$  and  $K=1 \rightarrow 1$  transitions were expected from theory to show second- and first-order Stark effect, respectively. This was strikingly observed in the experiment when only the  $K=1 \rightarrow 1$  lines appeared at field of 10-20 volts/cm, being augmented by the  $K=0 \rightarrow 0$  lines when the field was increased to several hundred volts/cm. The determination of the quadrupole coupling constants,  $eQ(\partial^2 V / \partial Z^2)$ , was based on the frequencies of the  $K=0 \rightarrow 0$  lines only, since these could be measured more accurately than the  $K=1 \rightarrow 1$  lines because of experimental considerations.

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<sup>1</sup> Davis, Field, Zabel, and Zacharias, Phys. Rev. **73**, 525 (1948). See also reference 2.

<sup>2</sup> W. Gordy, J. W. Simmons, and A. G. Smith, Phys. Rev. **74**, 243 (1948).

<sup>3</sup> C. H. Tindal, J. W. Straley, and H. H. Nielsen, Phys. Rev. **62**, 151 (1942).

<sup>4</sup> R. L. Livingston and L. O. Brockway, J. Am. Chem. Soc. **68**, 719 (1946).

## Argon and Neon ( $\beta n$ ) Thresholds

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FROM mass values<sup>1</sup> one calculates a  $Q = -1.1 \pm 1.1$  Mev for the reaction  $A^{40}(\beta n)K^{40}$ . Observation of the neutron threshold should give precisely the mass difference between the important isobars  $A^{40}$  and  $K^{40}$ . A gas target was recently available on the Wisconsin electrostatic generator so a quick search was made for this threshold. A BF<sub>3</sub> counter surrounded by paraffin served as neutron detector. The high resolution electrostatic analyzer<sup>2</sup> was