

reported here, we obtain  $1.008992 \pm 0.000010$  amu for the mass of the neutron. The probable error given includes an uncertainty of 0.3 percent in the absolute energy value for the  $\text{ThC}''$   $\gamma$ -ray. This neutron mass value is 0.051 mu greater than the value quoted by Stephens<sup>1</sup> and gives a value for the  $n-H^1$  difference of  $0.804 \pm 0.009$  Mev. The theoretical lifetime of the neutron is reduced by a factor of 1.3 when this new value for the  $n-H^1$  difference is used.

<sup>1</sup> W. E. Stephens, *Rev. Mod. Phys.* **19**, 19 (1947).

<sup>2</sup> K. Kimura, *Kyoto Coll. Sci. Mem.* **22**, 237 (1940).

<sup>3</sup> F. E. Myers and L. C. Van Atta, *Phys. Rev.* **61**, 19 (1942).

<sup>4</sup> R. Cohen and W. R. Hornyak, *Phys. Rev.* **72**, 1127 (1947).

### On the Radioactivity of $\text{K}^{40}$

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SOME years ago we published measurements of the quantum energy and the upper limit of the beta-ray spectrum of  $\text{K}^{40}$ . The values found by us are, respectively,<sup>1</sup>

$$E_{\gamma} = 1.54 \pm 0.1 \text{ Mev}, \quad E(\beta_{\text{max}}) = 1.41 \pm 0.02 \text{ Mev}.$$

These data are in good agreement with later measurements by Meyer *et al.*,<sup>2</sup> Gleditsch and Gráf,<sup>3</sup> Dželepov *et al.*,<sup>4</sup> and Henderson.<sup>5</sup> Recently, Franchetti and Giovanozzi,<sup>6</sup> using the cloud-chamber method, obtained a much higher value for the maximum beta-ray energy of  $\text{K}^{40}$ , namely,  $1.7 \pm 0.1$

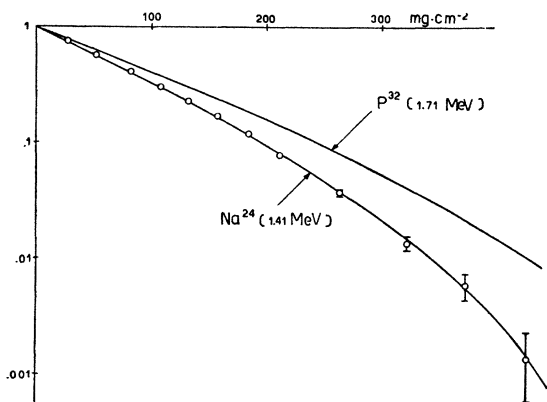


FIG. 1. Absorption curves in aluminium of  $\text{P}^{32}$  and  $\text{Na}^{24}$ . The points refer to  $\text{K}^{40}$ .

Mev. We believe that with respect to our measurements such a high beta-ray energy is rather improbable. Figure 1 gives the absorption curves obtained by us with  $\text{P}^{32}$  ( $E_{\text{max}} = 1.71$  Mev),  $\text{Na}^{24}$  ( $E_{\text{max}} = 1.41$  Mev), and  $\text{K}^{40}$  in the same geometrical arrangement. The points for  $\text{K}^{40}$  are taken up to 1/500 of the initial intensity and are all lying on the  $\text{Na}^{24}$  curve.

In addition we have determined the number  $\Gamma$  of quanta emitted per 100 beta-rays. For this purpose, the radiation from a thick KCl sample (cylindrical arrangement) was measured (a) with a thin-walled ( $27 \text{ mg/cm}^2$  Al) G-M

counter and (b) with a cylindrical absorber, thick enough to absorb all the beta-rays, between the sample and the counter. The same measurements were performed with  $\text{Al}^{28}$ , which is known to emit one quantum of 1.8 Mev per beta-ray. The ratio of the sensitivities for  $\gamma$ -rays of 1.54- and 1.8-Mev quantum energies is 0.84 for Al counters, as computed by Bleuler and Zünti<sup>7</sup> and obtained experimentally by Bradt *et al.*<sup>8</sup> Taking into account the self-absorption of the beta-rays in the samples and their absorption in the counter wall, we obtain for  $\Gamma$

$$8.7 \pm 1.2 \text{ } \gamma\text{-quanta per 100 beta-rays.}$$

Furthermore we have determined the half-life of the transition  $\text{K}^{40} \rightarrow \text{Ca}^{40}$ . The number of counts from a thin sample ( $4 \text{ mg/cm}^2$ ) of purified KCl (cylindrical arrangement) was compared with a very thin ( $< 1 \text{ mg/cm}^2$ )  $\text{U}_3\text{O}_8$  sample. The back-scattering from the holder (0.01-mm Al foil) was determined to be smaller than 1 percent. Taking into account the somewhat different absorption of the two beta-spectra in the wall of the G-M counter, as well as the very weak intensity of UX<sub>1</sub> and UY radiation passing through the counter wall, we get

$$T_{1/2}(\text{K}^{40}) = T_{1/2}(\text{U}^{238}) \times 0.246$$

$$T_{1/2}(\text{K}^{40}) = (11.1 \pm 1.9) \times 10^8 \text{ a.}$$

The greatest contribution to the error in  $T_{1/2}$  is given by the uncertainty in the relative abundance of  $\text{K}^{40}$  ( $0.011 \pm 0.001$  percent).<sup>9</sup>

We are very indebted to Professor P. Scherrer for his stimulating interest in this work.

- <sup>1</sup> O. Hirzel and H. Wäffler, *Helv. Phys. Acta* **19**, 216 (1946).  
<sup>2</sup> H. A. Meyer, G. Schwachheim, and M. D. de Souza Santos, *Phys. Rev.* **71**, 908 (1947).  
<sup>3</sup> E. Gleditsch and T. Gráf, *Phys. Rev.* **72**, 640 (1947).  
<sup>4</sup> B. Dželepov, M. Kojava, and E. Vorobjov, *Phys. Rev.* **69**, 538 (1946).  
<sup>5</sup> W. J. Henderson, *Phys. Rev.* **71**, 323 (1947).  
<sup>6</sup> S. Franchetti and M. Giovanozzi, *Phys. Rev.* **74**, 102 (1948).  
<sup>7</sup> E. Bleuler and W. Zünti, *Helv. Phys. Acta* **19**, 375 (1946).  
<sup>8</sup> Bradt, Gugelot, Huber, Medicus, Preiswerk, and Scherrer, *Helv. Phys. Acta* **19**, 77 (1946).  
<sup>9</sup> Alfred O. Nier, *Phys. Rev.* **50**, 1041 (1936).

### Observations of Naphthalene Scintillations Caused by Tritium Beta-Rays\*

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IT seems desirable to report here some preliminary observations on the scintillations produced in commercial naphthalene by the beta-rays from tritium and by the bremsstrahlung coming from tritium occluded in tantalum. Of immediate interest is the lower limit set on the conversion efficiency from beta-ray to visible light energy.

A small amount of gaseous tritium was put in direct contact with finely powdered naphthalene crystals in a 15-cm<sup>3</sup> glass Kjeldahl flask, an identical flask but without tritium being used as a control, to find the direct action of the betas. A tantalum disk containing tritium occluded throughout its volume was placed near a solid piece of