

TABLE I.

Altitude (feet)	Pressure (g cm ⁻²)	Time (hour)	A ₁ B ₂ C ₂ (counts)	A ₁ B ₂ C ₂ F ₁ (counts)	A ₂ B ₂ C ₂ (counts)	A ₂ B ₂ C ₂ F ₁ (counts)	A ₂ B ₂ C ₂ F ₂ (counts)	A ₂ B ₂ C ₂ F ₃ (counts)	A ₂ B ₂ C ₂ F ₄ (counts)
250	1030	308 113.5	309	16	25	3			
9500	725	58			174	17	12	10	6
14,000	625	41.5			.293	15	10	9	8

density. The area of each counter is 67.5 cm². If one assumed that all of the showers striking tray *F* had the same average density, one would calculate this density as being about 220 particles per square meter.

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¹ J. Tinlot, *Phys. Rev.* **73**, 1476 (1948).

² D. Broadbent and L. Janossy, *Proc. Roy. Soc.* **A192**, 364 (1948).

Development of Thick Emulsions by a Two-Bath Method*

M. BLAU

Columbia University, New York, New York

AND

J. A. DE FELICE

Brookhaven National Laboratory, Upton, Long Island, New York

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THE increased use of thick emulsions in nuclear physics has made it desirable to find a satisfactory technique of uniformly developing them. Dilworth, Occhialini, and Payne¹ have described the so-called temperature development method. We are using an alternate method of development on Ilford, C2, 200 u plates, which is probably applicable to even thicker emulsions. The results we have obtained to date might be helpful to others who are using these emulsions.

The method we adapted for our purpose is essentially that described by Crabtree *et al.*² which was used for the uniform development of large quantities of motion picture film. In this method the developer is divided into two baths. The first bath contains the developing agent, part of the sodium sulfite and the potassium bromide, but no alkali. The second bath contains all the necessary constituents of an ordinary developer plus an additional amount of alkali. In the first bath the developer diffuses into the emulsion. However, the rate of development is very low because of the lack of alkali. In the second bath the actual development takes place because of the presence of the alkali. It was necessary to add developing agent to the second bath, because not enough can be absorbed from the first bath.

After trying various combinations of the constituents and different times of development we find the following procedure to give the best results.

Step 1: Soak in water for 10 min.

Step 2: Solution *A* for 30 min. (slight agitation).

Step 3: Solution *B* for 30 min. (no agitation).

Step 4: 2 percent acetic acid 15 min. (agitation).

Step 5: Fix in *F-5* at 74°F with constant agitations 6–8 hours.

Step 6: Wash in running water 2 hours.

Solution *A*:

Elon	1.1 g.
Na ₂ SO ₃	24.0 g.
Hydroquinone	4.4 g.
KBr	2.0 g.
H ₂ O to make	2000 cc

Solution *B*:

Stock <i>D-19</i>	400 cc
H ₂ O	1600 cc
Additional Na ₂ CO ₃	16 g.

For different batches of the same emulsion, slight adjustments of the developing times and the composition of the solutions may be necessary. The temperatures of the solutions in the Steps 1–4 were all kept constant at 68°F. The temperature of the fixer could also be kept at 68°F. However, it was increased to 74°F to shorten the fixing time.

Because the temperature is kept constant the danger of reticulation is avoided. None of our plates showed any sign of reticulation. Proton tracks in the emulsion had their normal grain density while the background fog was very low. The plates appeared to be uniformly developed throughout the emulsion.

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¹ C. C. Dilworth, G. P. S. Occhialini, and R. H. Payne, *Nature* **162**, 102 (1948).

² Crabtree, Parker, and Russel, *Soc. Mot. Pic. ENO, VO RR* **21**, 21 (1933).

Thermonuclear Reactions in the Expanding Universe

R. A. ALPHER AND R. HERMAN

Applied Physics Laboratory, The Johns Hopkins University,
Silver Spring, Maryland*

AND

G. A. GAMOW

The George Washington University, Washington, D. C.

September 15, 1948

IT has been shown in previous work¹⁻³ that the observed relative abundances of the elements can be explained satisfactorily by consideration of the building up of nuclei by successive neutron captures during the early stages of the expanding universe. Because of the radioactivity of the neutron, and also because neutrons are used in forming the elements, the building up process must have been completed essentially in a time of the order of several neutron decay periods, i.e., about 10³–10⁴ sec. It should be noted that following the essential completion of the main element forming process, the temperature prevailing

in the expanding universe was still high. Consequently, one should expect changes in the relative abundances of the lighter elements as a result of thermonuclear reactions of these nuclei with protons. While these thermonuclear reactions must have proceeded during the element forming stage, their rates were negligible by comparison with those of neutron capture reactions. Examination of the relative abundance data suggests that the abnormally low abundances of certain of the light elements, such as lithium, beryllium, and boron, resulted from thermonuclear destruction of these nuclei with the concomitant formation of additional helium. In contradistinction to the other light elements, these are found to have large cross sections for proton reactions.⁴

On the basis of the foregoing discussion it may be shown that for certain of the light elements the difference between the present relative abundance and the abundance computed according to the neutron capture process is consistent with known thermonuclear reaction rates and cosmological information afforded by the neutron capture process. The well-known formula for the rate of thermonuclear reactions gives the number of processes per gram per second as^{4,5}

$$p = 5.3 \times 10^{26} \rho x_1 x_2 \Gamma \phi \tau e^{-\tau} \text{ g}^{-1} \text{ sec}^{-1}, \quad (1)$$

where ρ is the total density of matter, x_1 and x_2 the concentrations by weight of the nuclei involved in the reaction, Γ/\hbar the probability of the reaction in sec^{-1} after penetration, ϕ a function of the nuclear parameters, and

$$\tau = \alpha T^{-1/3}, \quad (1a)$$

where T is given in units of 10^6K . The time dependence of temperature and density⁶ of matter in the expanding universe, consistent with the conditions required by the neutron capture process, has been shown to be^{7,8}

$$T = 2.14 \times 10^{10} t^{-1/2} \text{K}, \quad (2a)$$

and

$$\rho = 7.2 \times 10^{-3} t^{-3/2} \text{ g cm}^{-3}. \quad (2b)$$

Using Eqs. (1) and (2) in

$$\frac{d[\ln x_2]}{dt} = -p m_2, \quad (3)$$

where m_2 is the mass of the nuclear species 2, with the concentration by weight of protons assumed constant, $x_1 = 0.5$, we obtain

$$\ln R = A [I(t_0) - I(t)], \quad (4)$$

where R is the ratio of the observed relative abundance to that computed by the neutron capture theory,¹⁻³

$$A = -2.96 \times 10^{21} m_2 x_1 \Gamma \phi \alpha^2, \quad (5a)$$

$$I(t) = t^{-1/6} \exp(-Bt^{1/6}) + BEi(-Bt^{1/6}), \quad (5b)$$

and $B = 3.60 \times 10^{-2} \alpha$. In Eqs. (5), t_0 represents the time at which the proton reactions became important, and t represents the current epoch. The quantities Γ , ϕ , and α are given by Bethe⁴ for a number of reactions. The quantity α is $(20)^{1/3}$ times the value of τ tabulated by Bethe for $T = 20 \times 10^6 \text{K}$.

Applying Eqs. (5) to the reactions of Li, Be, and B with protons, we find that in order to explain the present extremely low relative abundances of these elements by

thermonuclear destruction one must assume that the process became important at about $t_0 \sim 10^3 \text{ sec}$. This value of t_0 is of the correct order of magnitude for the time at which the neutron capture process became unimportant.

The same analysis applied to other light elements, such as F^{19} , which also have rather low abundances compared to their computed values, yields a similar result. On the other hand, the elements involved in the carbon-nitrogen cycle reactions do not show any appreciable depletion because of thermonuclear reactions in the expanding universe. For example, in the case of N^{14} the assumption that $t_0 \sim 10^3 \text{ sec}$ leads to the result that up to the present epoch the abundance of this nuclear species would have been reduced by only one part per million.

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¹ R. A. Alpher, H. A. Bethe, and G. A. Gamow, *Phys. Rev.* **73**, 803 (1948).

² R. A. Alpher, *Phys. Rev.* (in press).

³ R. A. Alpher and R. C. Herman, *Phys. Rev.* (in press).

⁴ H. A. Bethe, *Phys. Rev.* **55**, 434 (1939).

⁵ G. A. Gamow and E. Teller, *Phys. Rev.* **53**, 608 (1938).

⁶ We wish to point out that Eq. (2b) gives the time dependence of matter density in an expanding universe controlled by radiation. It is believed that the error involved in employing this relation up to the present epoch is not large.

⁷ G. A. Gamow, *Phys. Rev.* **74**, 505 (1948).

⁸ G. A. Gamow, *Nature*, in press.

On the Total Half-Life Period of K^{40}

T. GRÄF

L. Meitner Laboratory of Nuclear Physics, Stockholm, Sweden
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IN connection with the recent determination¹ of the β -decay constant of K^{40} , which yielded $\lambda_\beta = (0.505 \pm 0.055) \times 10^{-9} \text{ year}^{-1}$, with a corresponding β -decay period $T_\beta = (13.7 \pm 1.5) \times 10^8 \text{ years}$, an attempt was made to revise the branching ratio λ_e/λ_β , where λ_e is the decay constant for electron capture, as obtained by Bleuler and Gabriel,² in order to allow for the effect of back-scattering. With the corrected branching ratio (1.8 ± 0.5) an approximate value of $(5 \pm 1.5) \times 10^8 \text{ years}$ was computed for the total half-life T .

At present a more accurate value of T may be calculated by combining the above value of λ_β with Ahrens and Evans'³ recent $\text{Ca}^{40}/\text{K}^{40}$ ratio in lepidolites, which became available after the mailing date of the author's notes.^{1,4} Since in these minerals, of age $t = 2.1 \times 10^9 \text{ years}$, an excess

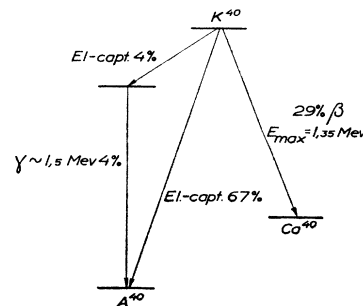


FIG. 1. Decay scheme of K^{40} with relative branching intensities.