

energy (<100 kev) or of low intensity (< ten percent of disintegrations if gamma-ray were of 1 Mev). Further comparison of the He⁶ curve with the P³² and UX₂ curves gives ranges of 1.85 and 1.88 g/cm² of aluminum, respectively. This corresponds to a maximum energy of 3.7 Mev, using the range-energy relationship $E_{\max} = 1.85R + 0.245$,⁶ with an estimated possible error of ± 0.2 Mev. This value for the energy is in close agreement with previously reported values.^{4,5}

Appreciation is expressed to Professor E. M. McMillan for suggesting this problem. This paper is based on work performed under the auspices of the Atomic Energy Commission in connection with the Radiation Laboratory, University of California, Berkeley, California.

- ¹ Delsasso, White, Barkas, and Creutz, Phys. Rev. **58**, 586 (1940).
² Haxby, Shoupp, Stephens, and Wells, Phys. Rev. **58**, 1035 (1940).
³ T. M. Putnam, to be published.
⁴ Bjerger and Brøstrom, Kgl. Danske Vid. Sels. Math.-Fys. Medd. **16**, No. 8 (1938).
⁵ H. S. Sommers, Jr. and R. Sherr, Phys. Rev. **69**, 21 (1946).
⁶ L. E. Glendenin, Nucleonics **2**, No. 1 (Jan. 1948).

The Nuclear Spin and Magnetic Moment of Na²² *

LUTHER DAVIS, JR.

Research Laboratory of Electronics, Massachusetts Institute of Technology, Cambridge, Massachusetts
 September 13, 1948

THE atomic beam magnetic resonance method,¹ as used by Zacharias² to determine the nuclear spin and magnetic moment of K⁴⁰, has been modified to permit similar observations to be made with radioactive nuclei having considerably shorter half-lives. These modifications involve the use of a mass spectrometer for separating the Na²² peaks from the large background of Na²³ and a beryllium copper multistage electron multiplier^{3,4} for detecting the Na²². The source of atoms for the beam is a monel metal oven with an exit canal long compared with its width, designed so as to conserve radioactive material. A sample of NaN₃ containing 330 microcuries of Na²² (3×10^{-9} mole) in a dilution of one part in 10^4 of Na²³N₃ was placed in the oven. Decomposition occurs at 320°C, providing an atomic sodium beam. A 16-hour run was sufficient for all present determinations and required the evaporation of only 4×10^{-10} mole of Na²² out of the oven. The beam intensity of Na²² was naturally small, a few thousand atoms/sec., and the radiofrequency transition intensities were less than 100 atoms/sec.

Two types of transitions were observed, first at low frequency with $\Delta m_j = \pm 1$ and $\Delta F = 0$. These permit direct comparison of g_F for Na²², with g_F for Na²³. For an unexcited alkali atom, $F = I + \frac{1}{2}$, and the spin of the Na²² is obtained as soon as a transition is observed at the mass position 22 of the mass spectrograph.

Search was made for resonances with the homogeneous field (as determined by the Na²³ transition) and radiofrequency set for transitions corresponding to $I = 0, 1, 2, 3, 4, 5$, and 6, and with the mass spectrometer set for 22. Resonances were observed only for spin 3 at mass 22.

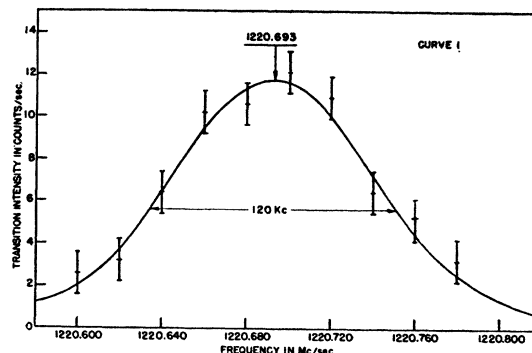


FIG. 1. Transition intensity vs. frequency of oscillating field for the two Na²² lines $F = 7/2$, $M_f = \pm 1/2 \rightarrow F = 5/2$, $M_f = \pm 1/2$ at a field of 6.17 gauss. One count/sec. is equivalent to about 7 atoms/sec.

These peaks faded out for mass numbers 22 $\frac{1}{2}$ and 21 $\frac{1}{2}$ and could not possibly have been a spurious effect from the more abundant Na²³. These low frequency resonances were followed as the magnetic field was increased to permit an evaluation of the h.f.s. $\Delta\nu$ of \bullet Na²². The second type of transition can then be made in which $\Delta F = \pm 1$. Curve 1 shows the result of an experiment in which a $\Delta F = \pm 1$ transition was made in almost zero field.

The results show that the spin of Na²² is 3 in units of $h/2\pi$, and its h.f.s. $\Delta\nu$ is 1220.64 ± 0.04 mgc/sec. By comparison with the value of the nuclear magnetic moment of Na²³, given by Millman and Kusch,⁵ we obtain for Na²², 1.746 ± 0.003 nuclear magnetons. The sign of this nuclear moment was also shown to be positive in just the way that was used by Zacharias on K⁴⁰.

Good, Peaslee, and Deutsch⁶ report that the disintegration of Na²² goes by 100 percent positron emission to an excited state of Ne²², and they conclude also that the spin of this excited state differs from that of Na²² by at most one unit. The spin of Ne^{22*} is therefore either 2, 3, or 4. Further unpublished data by Professor Deutsch indicate that the effective half-life for the possible transition to the ground state of Ne²² is at least 5000 years, consistent with a spin change of three units.

I am greatly indebted to Dr. E. T. Clarke who prepared the Na²² by a $Mg^{24}(d,\alpha)Na^{22}$ reaction in the M. I. T. cyclotron, to Professor J. W. Irvine, Jr., for the separation⁷ of the Na²² from the cyclotron target and for the conversion of the Na²²Cl to Na²²N₃, to Dr. Darragh Nagle for the design and construction of the molecular beam apparatus adapted for use in this experiment, and to Professor J. R. Zacharias who suggested the problem and some of the methods, and with whom constant discussions of the problems involved led to the successful conclusion of this experiment.

* This work has been supported in part by the Signal Corps, the Air Materiel Command, and the Office of Naval Research.

¹ P. Kusch, S. Millman, and I. I. Rabi, Phys. Rev. **57**, 765 (1940).

² J. R. Zacharias, Phys. Rev. **61**, 270 (1942).

³ J. S. Allen, Phys. Rev. **55**, 967 (1939).

⁴ Hin Lew, *Hyperfine Structure of Aluminum by Atomic Beam Method* (M. I. T. thesis, 1948).

⁵ S. Millman and P. Kusch, Phys. Rev. **58**, 438 (1940).

⁶ W. M. Good, D. Peaslee, and M. Deutsch, Phys. Rev. **69**, 813 (1946).

⁷ J. W. Irvine, Jr. and E. T. Clarke, J. Chem. Phys. **16**, 686 (1948).