## The Absolute Strength of Radioactive Sources

M. L. WIEDENBECK Department of Physics, University of Michigan, Ann Arbor, Michigan September 15, 1947

T has been the custom to determine the strength of a radioactive material by comparing the counting rate of the material in question with a "standard" source. Usually the energies of the "standard" and test sources are not comparable and, as such, corrections for absorption are required. Thus absolute source strengths usually are not very reliable, thereby sometimes preventing the accurate comparison of work from different laboratories.

In the case of ThC", Dunworth1 has pointed out a method for obtaining the  $\beta$ - and  $\gamma$ -efficiencies of counters as well as the strength of a source. The calculations are set up for the modified-level scheme proposed by Oppenheimer.

We have found that the absolute number of disintegrations per second in a source can be found simply from the simultaneous  $\beta$ - and  $\gamma$ -counting rates and from the  $\beta$ - $\gamma$ coincidence rate, independently of the solid angle subtended by the counters and of the  $\beta$ - and  $\gamma$ -efficiencies of the counters.

The relationship is

$$\frac{N_{\beta}N_{\gamma}}{N_{\beta\gamma}} = N. \tag{1}$$

 $N_{\beta}$  and  $N_{\gamma}$  are the respective counting rates of the  $\beta$ - and  $\gamma$ -counters,  $N_{\beta\gamma}$  is the  $\beta$ - $\gamma$ -coincidence rate, and N is the number of disintegrations per second occurring in the source placed between the two counters.

For a nucleus which has a simple  $\beta$ -spectrum followed by  $\gamma$ -radiation, no assumptions need be made. The value obtained is independent of the absorption of the  $\beta$ -radiation in the window of the  $\beta$ -counter as well as the absorption of the  $\gamma$ -rays in the walls of the  $\gamma$ -counter.

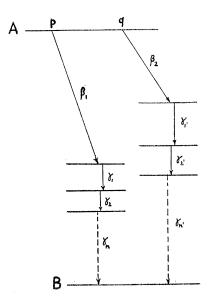


Fig. 1. Decay scheme for a complex  $\beta$ -spectrum.

For complex  $\beta$ -decay (Fig. 1)

$$\frac{N_{\beta\gamma}}{N_{\beta}} = \frac{\left[pF_1\left(\sum_{1'}^{n} \epsilon_n e^{-K_n T}\right) + qF_2\left(\sum_{1'}^{n'} \epsilon_{n'} e^{-K_{n'} T}\right)\right]\sigma_{\gamma}}{(pF_1 + qF_2)},\qquad(2)$$

$$N_{\gamma} = \left[ p(\sum_{i=1}^{n} \epsilon_{n} e^{-K_{n}T}) + q(\sum_{i=1}^{n'} \epsilon_{n'} e^{-K_{n'}T}) \right] N\sigma_{\gamma}. \tag{3}$$

p=fraction of disintegrations in which a beta-ray of maximum energy  $\beta_1$  is emitted (p+q=1).  $F_1$ =fraction of  $\beta_1$ transmitted through absorber.  $F_2$ =fraction of  $\beta_2$  transmitted through absorber.  $\epsilon_n$  = efficiency of  $\gamma$ -counter for  $\gamma_n$ .  $K_n$ =absorption coefficient for  $\gamma_n$ . T=thickness of  $\gamma$ -absorber (counter wall, etc.).

At zero beta-absorber thickness  $F_1$  and  $F_2=1$ . Then  $N_{\gamma}N_{\beta}/N_{\beta\gamma} = N$  as in Eq. 1. That is, the absolute disintegration rate can be obtained for a complex  $\beta$ -decay if the value  $N_{\beta\gamma}/N_{\beta}$  is taken at zero absorber thickness. In practice this would require that an extrapolation be made for the thickness of the counter window. This should introduce little error, however, since counter windows of only a few mg/cm<sup>2</sup> thickness are in common use.

It would seem that this method would offer a simple, accurate method of measuring long lifetimes, since no special care need be taken in reproducing counter efficiencies or geometries from one measurement to the next.

<sup>1</sup> J. V. Dunworth, Rev. Sci. Inst. 11, 167 (1940).

## Further Work on Satellites in the Microwave Spectrum of Ammonia\*

DUDLEY WILLIAMS The Ohio State University, Columbus 10, Ohio September 26, 1947

SINCE the publication of our work<sup>1</sup> on hyperfine structure in the microwave spectrum of ammonia, we have observed additional satellites near lines involving J values as high as 10. The intensities of these satellites are apparently in agreement with theoretical predictions; i.e., intensity relative to the main line decreases rapidly as J value increases. In observing line positions, incident power was varied until the satellite intensity relative to the main line was a maximum; i.e., until saturation phenomena had set in. This technique was also used in our published work.1 The separation of the satellites from the main line is also in apparent agreement with theory. Since as J values increase the satellites appear closer together and experimental measurements of  $\Delta \nu / \Delta v'$  become less accurate, this type of check of the theory becomes less exact as J values increase. As the results already published are in agreement with and merely extend other published work2 and give satisfactory values for quadrupole coupling, we shall not pursue the subject further.

This work has been done in connection with the research contract \* Ins work has been done in Contection with the research contract W28-099-ac-179 between Watson Laboratories, Air Materiel Command, Army Air Forces and The Ohio State University Research Foundation.

1 R. J. Watts and Dudley Williams, Phys. Rev. 72, 263 (1947).

2 Daily, Kyhl, Strandberg, Van, Vleck, and Wilson, Phys. Rev. 70, 984 (1946).