values. There are several reasons, however, why such information is of little use. First, at laboratory temperatures, the transport properties of gases depend upon molecular interactions which must be represented as a combination of attractive and repulsive forces. Second, it is possible to obtain essentially the same $\rho D/\eta$ values for different molecular-force models which are not demonstrably equivalent. This is illustrated by recent calculations of low temperature transport properties of gaseous helium by use of an interaction potential consisting of two repulsive and three attractive terms.⁴ These calculations predict a value of 1.18 for $\rho D/\eta$ at 300°K in close agreement with the value of 1.20 for classical hard spheres. And yet, a classical hardsphere model has on many occasions been shown to be incapable of representing the transport properties of gaseous helium as a function of temperature, in agreement with quantum-mechanical calculations which show the need for both repulsive and attractive terms in the interaction potential.

Table I, from a paper on the self-diffusion coefficient of neon by Groth and Sussner,⁵ shows that other gases besides helium have $\rho D/\eta$ values which lead to incorrect laws of force if interpreted on the basis of point-repelling molecules, since hydrogen has the "hardest" and xenon the "softest" force field of the group.

TABLE I.

Gas	$\rho D/\eta$ at 20°C
 H ₂	1.370 ± 0.003
Ne	1.275 ± 0.006
Kr	1.30 ± 0.06
Xe	1.24 ± 0.06

It should be pointed out that, for the reasons already given, the temperature variation of the coefficient of viscosity $(\eta \propto T^{\frac{1}{2}+2/(\nu-1)}$ for point-repelling molecules) provides no better information about the law of force than the ratio $\rho D/\eta$ or the temperature dependence of D if it were known. Chapman and Cowling⁶ summarize the difficulty with respect to both D and η in their statement that "viscosity data throw light on the appropriateness of a formula for the molecular field derived by other means, but by themselves they do not give unambiguous information as to the field.'

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Diamond as a Gamma-Ray Counter

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NONSIDERABLE interest has been evinced in the use • of crystals as counters for nuclear radiations since the report by van Heerden¹ in 1945 on the successful use of



silver chloride crystals for this purpose. Silver chloride crystals suffer the serious disadvantage that they must be cooled to liquid-air temperatures to reduce the spontaneous random conduction in these crystals at room temperatures.

Although van Heerden tried a single specimen of diamond he was unsuccessful in obtaining counting. However Woolridge, Ahearn, and Burton² have subsequently found that diamonds will count alpha-particles when electrodes are applied to a surface with a small separation. This counting effect, of course, arises in the surface layer of the crystal.

Diamond does not require cooling to operate as a counter and therefore it can well prove to be a convenient detector. It is of considerable interest to determine whether it responds to gamma-rays. Since the absorption per unit volume for this type of radiation is much greater in diamond than in a Geiger-Müller counter, considerable increase in volume sensitivity for gamma-rays can be anticipated. To survey the possibilities of diamonds as gamma-ray counters approximately 100 industrial diamonds of roughly ¹/₄-carat size have been tested. These diamonds were in their natural state. They were held in a simple compression mounting shown schematically in Fig. 1 in which the high voltage electrode pressed against one side of the crystal and the electrode leading to the grid of the first amplifier tube was in contact with the opposite face. From the total number tested two were found to give excellent response to radium gamma-rays and several others showed some response. The diamonds which gave good response were single crystals free from obvious flaws and water-white. This substantiates the general ideas of crystal conduction and those regarding properties required for counting.

The pulses delivered by the diamond are quite small, requiring considerable amplification for convenient observation. However many of the pulses are at least 10 times the background noise of the crystal plus that of the amplifying system so that no difficulty is encountered in recording them. The variation in pulse size indicates that they originate at random throughout the effective volume of the crystal.

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