



FIG. 1. Rate of change of atomic spacing with increasing copper content in eta and epsilon copper-zinc alloys.

in zinc content) causes a large decrease in atomic spacing in the directions of greatest atomic vibration in each of the two crystal lattices.

The pronounced anisotropy of zinc is demonstrated by these two alloy phases. Although both phases are hexagonal close-packed with the same coordination number, a complete series of solid solutions is not formed. There exists a gap in axial ratio and in composition between eta and epsilon within which no hexagonal close-packed copper-zinc phase exists. In eta the influence of the zinc atom is evidently in the shape of a prolate spheroid and in epsilon, in the shape of an oblate spheroid, without the intermediate spherical condition at axial ratio 1.633.

M. L. FULLER

The New Jersey Zinc Co.,
Research Division,
Palmerton, Pa.,
August 3, 1938.

¹ Robert A. Howard, Phys. Rev. 53, 966 (1938).

The Internal Friction of Metallic Crystals

Several investigators have studied the internal friction of polycrystalline metals,¹ but as far as the writer is aware the only observations on single metal crystals are those of Zacharias on nickel,² and these deal entirely with the ferromagnetic contribution. Accordingly the following results, obtained in the course of a systematic investigation of internal friction in single and polycrystalline metals, seem worth reporting at this time. The experimental method is that devised by Cooke and Brown.³

Observations have been made on single crystals of copper, lead and tin. The vibration frequency is 38.9 kilocycles in each case, and the temperature is about 25°C. The observed longitudinal decrements of the crystals are 3.6×10^{-5} , 2.8×10^{-4} , and 6.9×10^{-5} , respectively. These values may be compared with those of the corresponding polycrystalline substances obtained by Forster and Korster,¹ namely, 3.5×10^{-3} , 4.6×10^{-3} , and 5.4×10^{-3} . It thus appears that the internal friction of a single metal

crystal is substantially less than that of the polycrystalline material. The decrement of crystalline quartz is 3×10^{-6} , and it is suggested that the decrement of a pure and strain free metal crystal might approach this value.

The effect of internal strain on the internal friction of crystalline copper is noteworthy. The decrement of a crystal, on removal from the furnace in which it is grown, is of the same order as that of the polycrystalline material. The low value here reported is obtained by annealing the crystal in a vacuum. Lead and tin crystals, on the other hand, remain unaffected by annealing.

The writer is indebted to Dr. Clarence Zener for helpful suggestions made during the course of the present work.

T. A. READ

Pupin Physics Laboratories,
Columbia University,
New York, N. Y.,
August 15, 1938.

¹ See for example, Wegel and Walther, Phys. 6, 141 (1935); Forster and Korster, Zeits. f. Metallkunde 29, 116 (1937).

² Zacharias, Phys. Rev. 44, 116 (1933).

³ Cooke and Brown, Phys. Rev. 50, 1158 (1936).

Ignition Potential in a Low Pressure Neon Tube

In an article by Lindenhovius¹ an expression is derived and experimentally verified for the relation between the voltage at which ignition of a low pressure irradiated neon glow occurs and the frequency of the applied sinusoidal voltage. The expression was based on the determination of the time constant of an assumed exponential current rise,² and this time constant was shown to depend upon the magnitude of the overvoltage applied to the discharge. The time required for ignition can be computed from the relation

$$\log \frac{i_1}{i_0} = a \int_{t_0}^{t_1} (v - V_0) dt, \quad (1)$$

where v is the instantaneous electrode voltage, V_0 is the minimum ignition voltage, i_0 and t_0 are the current and time when $v = V_0$, i_1 is the current at time t_1 when ignition takes place and a is a constant depending upon the ionization coefficient, mobility of the ions, and number of electrons ejected from the cathode per ion pair formed. It is assumed that space charge does not distort the electric field.

If an exponentially increasing voltage, obtained from the plates of a condenser C charged through a resistance R by a battery V_m , is placed across the electrodes Eq. (1) may be conveniently tested. The term $[V_m(1 - e^{-t/RC}) - V_0]$ may be substituted for $(v - V_0)$ in Eq. (1) and the solution of the resulting equation may be written

$$\log \frac{i_1}{i_0} = a \left[(V_m - V_0)(t_1 - t_0) + RC V_m \left\{ \exp\left(-\frac{t_1}{RC}\right) - \exp\left(-\frac{t_0}{RC}\right) \right\} \right]. \quad (2)$$

A neon discharge tube with parallel-plane iron electrodes was used. The transient visualizer³ was used to apply the exponential voltage to the electrodes at a very low frequency, and at the same time to sweep the beam of a