

### Heavy Particles and the Neutrino

In a previous letter the writer<sup>1</sup> has suggested that the  $x$  particle may be produced by the absorption of a photon by an ordinary electron. Corson and Brode<sup>2</sup> give the rest mass of a heavy particle as  $350 m_0$ , while Street and Stevenson<sup>3</sup> give  $130 m_0$ . ( $m_0$  is the rest mass of the ordinary electron. See previous letter for definitions.) To produce a rest mass of  $p_2 m_0$  a photon of energy  $(p_2^2 - 1)m_0 c^2 / 2$  is needed. When  $p_2 = 350$ , this energy is 31,000 Mev. The values 130 and 350 for  $p_2$  and still other values are consistent with the writer's idea.

If a photon of  $\alpha_1$  greater than  $(p_2^2 - 1)/2$  strikes an ordinary electron ( $p_1 = 1$ ) at rest ( $\beta_1 = 0$ ), an energy  $\alpha' m_0 c^2$ , where  $\alpha' = (p_2^2 - 1)/2$ , is absorbed from the photon, giving rise to a particle of mass  $p_2 m_0$  with a velocity  $\beta' c = c(p_2^2 - 1)/(p_2^2 + 1)$  in the direction of the photon. The remaining energy  $(\alpha_1 - \alpha') m_0 c^2$  is treated as that of a primary photon which is to be scattered in the direction  $\theta$  according to the principles of the Compton effect by an electron of rest mass  $p_2 m_0$  traveling with an original speed  $\beta' c$  in the direction of the primary photon.

Perhaps heavy particles other than the heavy electron may be produced according to Eqs. (1) and (2) of the previous letter. To produce a heavy particle of mass  $p_2 m_0$  from a particle of mass  $p_1 m_0$  which is at rest a photon

$$h\nu_1 = m_0 c^2 (p_2^2 - p_1^2) / 2p_1 \quad (1)$$

is needed. The kinetic energy produced by complete absorption of this photon is

$$T_2 = m_0 c^2 (p_2 - p_1)^2 / 2p_1. \quad (2)$$

So, to produce a heavy proton with kinetic energy 940 Mev, a photon of energy 2260 Mev is needed. From (2),  $p_2 = 2.414 p_1$ .

In a  $\beta$ -ray radioactive transformation the  $\beta$ -rays are found to have a continuous spectrum of speeds up to a

certain maximum speed  $\beta_0 c$ , such that the energy relation

$$M_1 c^2 = M_2 c^2 / (1 - \beta_2^2)^{1/2} + m_0 c^2 / (1 - \beta_0^2)^{1/2} \quad (3)$$

holds, where  $M_1$  and  $M_2$  are the rest masses of the parent and product atoms, respectively, and  $\beta_2 c$  is the recoil velocity of  $M_2$ . The writer suggests that the energies of all  $\beta$ -rays in the spectrum are the same but that different  $\beta$ -rays have different rest masses  $p m_0$  and different velocities  $\beta c$ , such that

$$m_0 c^2 / (1 - \beta_0^2)^{1/2} = p m_0 c^2 / (1 - \beta^2)^{1/2}. \quad (4)$$

From (4), it is seen that, as  $p$  increases from unity to  $1/(1 - \beta_0^2)^{1/2}$ ,  $\beta$  decreases from  $\beta_0$  to zero. Allowing the expelled  $\beta$ -particle to have a variable rest mass thus removes the need for postulating the neutrino. The kinetic energy distribution curves for  $\beta$ -rays have been calculated on the basis of the rest mass always being  $m_0$ . According to the writer's suggestion, variation of  $H\rho$  implies variation of  $p$ . From the distribution curve with respect to  $H\rho$  for the  $\beta$ -particles from RaE, it is estimated that more than 90 percent of the electrons have a rest mass greater than  $2m_0$  and a kinetic energy less than half the maximum kinetic energy,  $2.06 m_0 c^2$ . The average kinetic energy for all the electrons is less than one-third of the maximum. This is in agreement with Meitner and Ortheim's calorimetric experiment<sup>4</sup> on RaE, provided that the life of the heavy electron is long compared with the duration of the experiment, or that the  $\gamma$ -radiation which the heavy electron ultimately emits in returning to  $m_0$  was not recorded in Meitner and Ortheim's experiments.

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<sup>1</sup> Jauncey, Phys. Rev. 52, 1256(L) (1937).

<sup>2</sup> Corson and Brode, Bull. Am. Phys. Soc., December 2 (1937).

<sup>3</sup> Street and Stevenson, Phys. Rev. 52, 1004 (1937).

<sup>4</sup> Meitner and Ortheim, Zeits. f. Physik 60, 143 (1930).