

Absorption Spectra of the Samarium Ion in Solids. III. Absorption of $\text{Sm}(\text{BrO}_3)_3 \cdot 9\text{H}_2\text{O}$ and a Partial Energy-Level Diagram for the Sm^{+++} Ion as It Exists in Crystalline $\text{Sm}(\text{BrO}_3)_3 \cdot 9\text{H}_2\text{O}$

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The absorption spectrum of $\text{Sm}(\text{BrO}_3)_3 \cdot 9\text{H}_2\text{O}$ has been studied at four temperatures between 20° and 298°K . The temperature variation in the absorption intensities is used to determine the location of excited lower levels situated at 39 and 68 cm^{-1} above the basic one. Other levels from 100 to 230 cm^{-1} above the normal state are probable. Between 78° and 20° many lines show a marked change, becoming double at the lower temperature. This is shown to be compatible with a splitting of the levels into new ones at 0, 37, 45, 67 and 82 cm^{-1} . Selection rules are indicated which show two apparently independent sets of these

levels; 0, 37, 82 and one group of high levels; and 0, 45, 67 and another group of upper levels. This low-temperature splitting may be due either (1) to actual removal of degeneracy or (2) to development of two types of Sm^{+++} ion, differing in crystal environments. Either effect could presumably be caused by small changes in the crystal dimensions and symmetry. The hexagonal bromate shows less extension and complexity of upper multiplets and the excited lower levels are closer to the basic one than in the less symmetrical monoclinic chloride previously studied.

THIS is the third of a series of papers describing the effects of chemical composition and crystal symmetry on the energy levels of Sm IV (Sm^{+++} ion) in the solid state. In the previous papers¹ the absorption spectra of monoclinic crystals of $\text{SmCl}_3 \cdot 6\text{H}_2\text{O}$ have been described. It was demonstrated at that time that the absorption lines result from transitions between several lower levels and higher ones, the latter being grouped in multiplets. The behavior of these multiplets with temperature variation strongly indicates that they are composed of components which have resulted from the splitting of one or more excited levels by the electric fields of the neighboring ions in the crystal (a crystal Stark effect). This conclusion is in excellent agreement with the discoveries of Spedding and Nutting² concerning the nature of the excited levels of the Gd^{+++} ion in various salts. It will also be substantiated in this and subsequent papers.

The lines arising from the excited lower levels are readily distinguished, not only by the frequently occurring constant energy differences of the spectrum, but also by the behavior of the

lines with temperature variation. As is well known, the intensity of a line is proportional to the number of atoms in the energy state from which it originates, and that number in turn depends on the temperature according to the Boltzmann relation. Because of this fact, as the temperature is lowered the intensities of lines arising from excited lower levels become less, and below a certain critical temperature, which depends on $\Delta\nu$, such lines become too faint to be observed.

This paper deals with the absorption spectra of hexagonal crystals of $\text{Sm}(\text{BrO}_3)_3 \cdot 9\text{H}_2\text{O}$ as shown by single crystals and by conglomerates of small crystals. In general the spectra resemble those described for the chloride and behave in a similar manner with regard to temperature shifts of multiplets and lines within them, changes in distinctness and intensity, etc. We shall describe only the points wherein the two spectra differ and shall refer the reader to the previous paper for general details.

EXPERIMENTAL PART

The bromate was prepared by the addition of a $\text{Ba}(\text{BrO}_3)_2$ solution to one of $\text{Sm}_2(\text{SO}_4)_3$ in a slightly less than equivalent amount. The $\text{Sm}(\text{BrO}_3)_3$ solution was then evaporated and the hydrated crystals extracted and purified by re-

¹Spedding and Bear, *Phys. Rev.* **42**, 58 (1932); **42**, 76 (1932).

²Spedding and Nutting, *Phys. Rev.* **38**, 2294 (1931); *J. Am. Chem. Soc.* **55**, 496 (1933).

crystallization. The crystals were measured by Dr. A. Pabst of this university. His preliminary results were as follows: The crystals are hexagonal prisms with very ragged multiple terminations by hexagonal pyramid faces. From the measurement of four crystals the angle $(10\bar{1}0) : (10\bar{1}1)$ was found to be $57^{\circ}1'$. This gives an axial ratio of 0.5620.

The samarium was from the same lot that had been used for the chloride preparation and was of such purity that no lines of other rare earths were observed. In addition to photographing single crystals with 3 to 6 mm absorbing paths, conglomerates of 5 to 10 mm in thickness were used. The differences between the spectra for the two cases were less than with the chloride, since the low-lying levels are much closer to the basic one. Even at hydrogen temperatures (20°) the population in these states is appreciable, making extreme absorbing layers unnecessary.

The apparatus and methods used in photographing the spectra at low temperatures have been described previously.¹

RESULTS

In Fig. 1 are reproduced the photographs of the spectra observed at 20° , 78° , 169° and 298°K , and in Table I are given measurements of the positions of lines and bands at these temperatures. Between 78° and 20° occurs a remarkable phenomenon which we have observed so far only for the bromate. Practically all the lines split into two components, the red one of which is almost certainly complex (see Fig. 2). Above 78° the spectra behave similarly to the other samarium salts that we have studied.³ The shift in position with temperature variation and the widening of the multiplets at lower temperatures may perhaps be slightly greater in the bromate case than in the chloride.

³ This is also true below 78° if only the red components of the doublets at 20° are considered. On the other hand, the violet components shift quite markedly to the shorter wave-lengths, which is contrary to the usual behavior. As a result, at 20° the "center of gravity" of a doublet frequently appears to have reversed the direction of temperature shift. This phenomenon may account for an apparent reversal between 78° and 20° in the direction of displacement of several diffuse lines that are probably unresolved doublets of this type.

The multiplets of excited energy levels, and consequently the low-temperature multiplets of the spectrum, on the whole obey rules that are similar to those found by Spedding and Nutting² for the Gd^{+++} ion. In crystals of higher symmetry the over-all spread of the spectral multiplets is less and frequently the number of components fewer than in crystals of less symmetry. The samarium bromate multiplets, however, are shifted to the red with respect to those of the chloride, while with the gadolinium salts the opposite is true. In some cases the samarium bromate multiplets seem to have more components than the corresponding ones in the chloride. This is almost certainly an illusion, since the extra lines may arise from the excited lower levels. In the case of the chloride these lines are absent at 20°K , while in the bromate they can be quite intense at that temperature as a result of the smaller separations between these low levels. Another factor which tends to complicate matters is an apparent difference in the selection rules that are effective in the two salts. Thick conglomerates (long optical path) of the chloride tend to bring out many faint lines, causing the appearance of the spectrum to be markedly changed from that of a single crystal. In the bromate only a few new lines appear.

The lower levels of the ion in the bromate are also "pulled-in" from their chloride position.⁴ Fig. 3 shows an energy-level diagram of several of the multiplets whose individual lines could be most satisfactorily studied. Lines arising from the low excited levels were identified, as before with the chloride, not only by their constant energy differences but also by their temperature

⁴ The close proximity of the low excited levels to the basic one makes the bromate spectrum an excellent example of a type whose temperature variation has frequently been erroneously reported in the literature. When spectroscopic instruments of low dispersion are employed related high and low-temperature lines appear to form an unresolved band. On account of the fading of the former and the increasing intensity of the latter type of line as the temperature is lowered, such a band may appear to shift to shorter wave-lengths, whereas the individual lines themselves may be actually displaced in the opposite direction. That the movement of multiplets and most lines is to longer wave-lengths at low temperatures has now been confirmed in this laboratory for quite a few gadolinium and samarium salts.

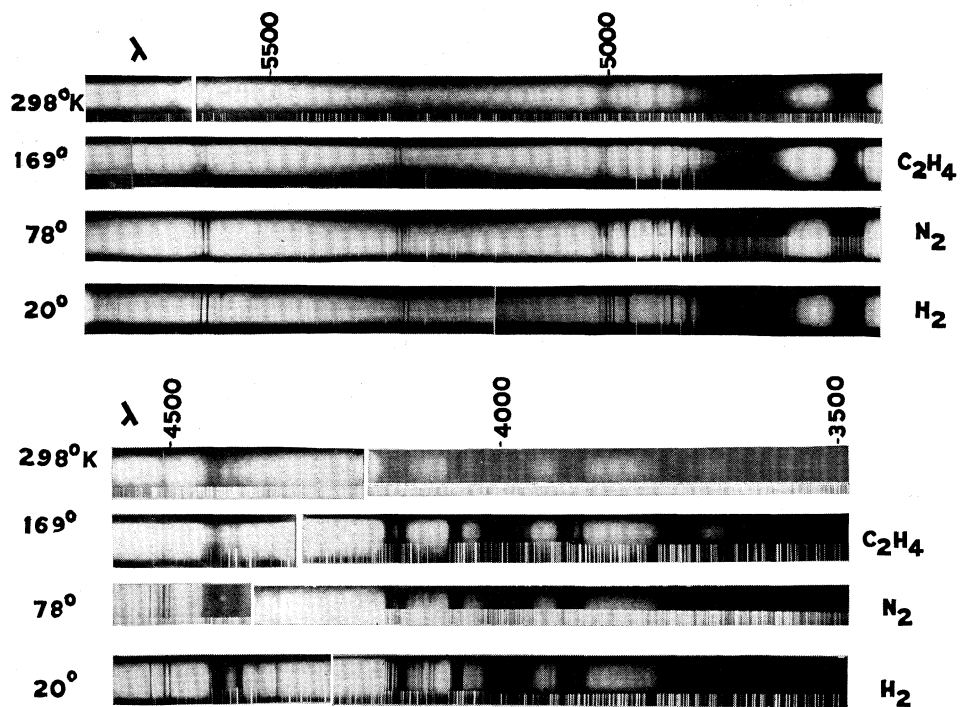


FIG. 1. Conglomerate absorption of $\text{Sm}(\text{BrO}_3)_3 \cdot 9\text{H}_2\text{O}$. The substances at whose boiling points the spectra were photographed are indicated on the right, the corresponding temperatures on the left. The three conspicuously simple multiplets on the long wave-length end are considered in more detail in Figs. 2 and 3.

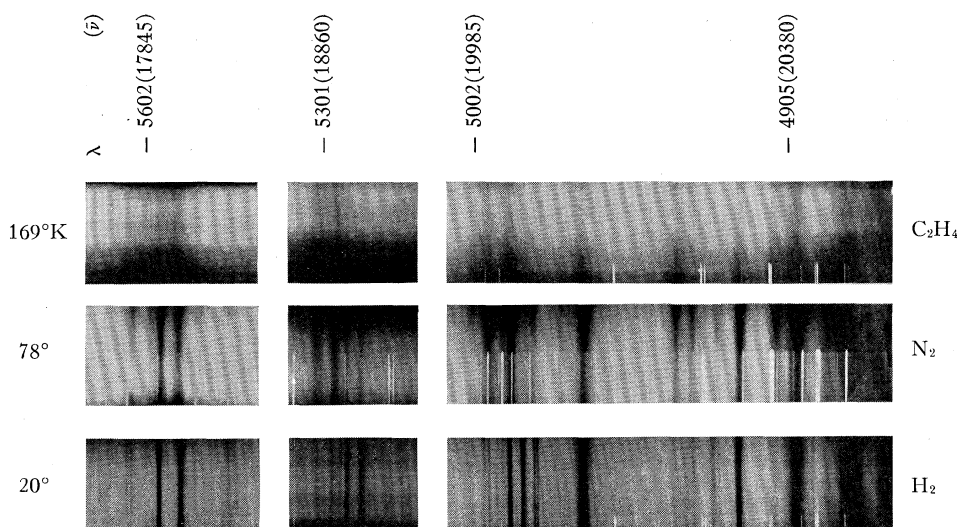


FIG. 2. Selected multiplets of $\text{Sm}(\text{BrO}_3)_3 \cdot 9\text{H}_2\text{O}$ absorption. Several of the multiplets of simplest structure are shown here magnified. The temperatures are given on the left, these being the boiling points of the substances on the right. Above are quoted wave-lengths and wave numbers for purposes of orientation with respect to the other figures and tables.

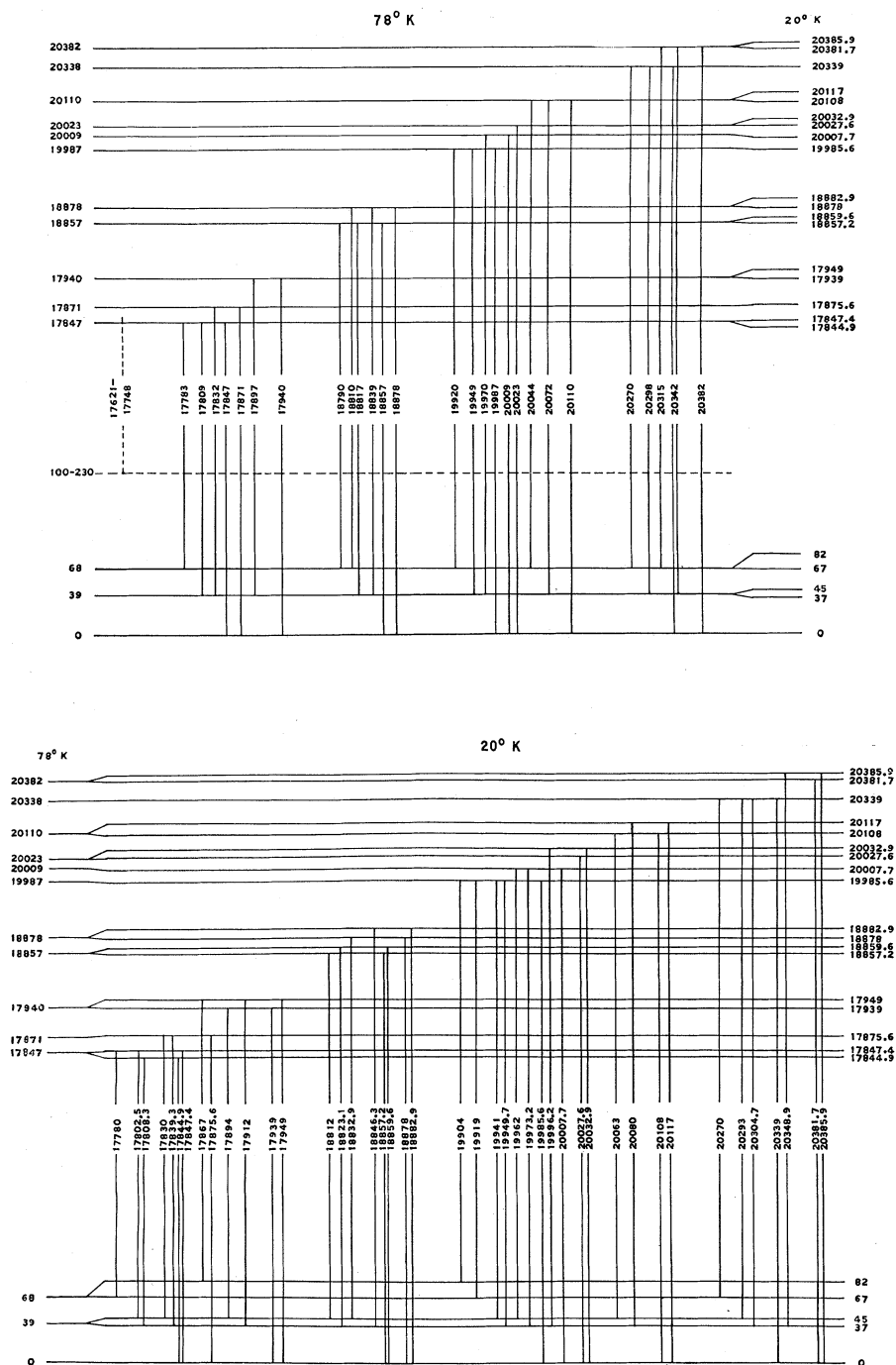


FIG. 3. Energy levels of $\text{Sm}(\text{BrO}_3)_3 \cdot 9\text{H}_2\text{O}$. Fig. 3a shows the levels needed to account for the multiplets of Fig. 2 at 78°K . The dotted lines indicate indistinct levels and transitions which are most evident at ethylene temperature (for other examples see text). In Fig. 3b are given the same levels after the changes which occur between 78° and 20° . The figures are wave numbers (cm^{-1}).

TABLE I. (Continued).

R. T. $\bar{\nu}$ (cm ⁻¹)	E. T. $\bar{\nu}$ (cm ⁻¹)	N. T. $\bar{\nu}$ (cm ⁻¹)	Int.	Type	H. T. λ (Å)	H. T. $\bar{\nu}$ (cm ⁻¹)	Int.	R. T. $\bar{\nu}$ (cm ⁻¹)	E. T. $\bar{\nu}$ (cm ⁻¹)	N. T. $\bar{\nu}$ (cm ⁻¹)	Int.	Type	λ (Å)	H. T. $\bar{\nu}$ (cm ⁻¹)	Int.
†24593'	†24579' 24588	24588.9	8nd	L	4065.85c	24588.2c	9nd	26509'		26549	8d	L	3771.9	26505	2d
					4065.05c	24593.1c	9nd						3766.4	26543	6d
24715'					4056.8'	24643'	2d						3764.6	26556	6d
					4054.2'	24659'	8d						3762.7'	26569'	8nd "2"
					4051.8'	24674'	2d						3761.30c	26579.2c	9db
					4048.8'	24692'	L						3758.8	26597	9db "2"
					4044.5'	24718'	4vd						3755.4	26621	10db
					4041.5'	24740'	L						3753.0'	26638'	9d "2"
					4038.1'	24757'	3vd						3747.7'	26675'	
					4034.8'	24777'	L						3746.4'	26685'	
					4032.7'	24790'	L						3744.1'	26701'	
					4029.9	24807'	4d						3742.8	26710	3d
					4027.2c	24824c	10d						3741.6'	26719'	9db
					4006.3'	24953'	L						3737.1'	26751'	
					4004.3	24966	L						3725.5	26834	8d
					3997.8	25007'	L						3722.8	26854	8d
					3989.6'	25058'	L						3718.5	26885	4d
3987.4'	25072'	L	3716.6	26899	6d										
3985.5'	25084'	L	3713.8	26919	8d										
3984.5'	25090'	L	3712.0	26932	9d										
3982.7'	25102'	L	3710.4	26944	7d										
3981.2'	25111'	L	3706.9	26969	7d										
3978.9'	25126'	L	3702.7	27000	6d										
3978.2'	25130'	L	3698.4'	27031'	6d										
3976.1'	25143'	L	3695.0'	27052'											
3974.6	25153	L	27042												
3971.6	25171	L	27273												
3968.50	25191.2	L	27288'												
3965.1'	25213'	L	27297'												
3961.1'	25238'	L	27302												
3958.0	25258	L	27317'												
3957.1'	25264'	L	27344	27340											
3956.25	25260.4	L	27359	27350											
3954.3	25282	L	27363												
	Faintly evident	H	27457'	27453'											
3939.75	25375.1	L	27488'	27481'	9d										
3935.3	25404	L	27499'	27492'	8db										
3932.7	25421	L	27524'	27512'	8db										
3930.2	25437	L	27541'	27534'	8db										
3922.2	25489	L	27571'	27617'	8db										
		L	27591'	27637'											
3919.80	25504.4	L	27621'	27662'											
3917.8	25517	L	27634'	27809'											
3914.3'	25540'	L	27663'												
3908.8'	25576'	L	27810'	27862'											
3906.7	25590	L	27866'	27886'											
3904.4	25605	L	27879'												
3902.7'	25616'	L	27919'	27926'											
3900.7'	25629'	L		28055											
3900.0'	25634'	L		28087											
3897.9'	25648'	L		28067											
3896.55	25656.6	L		28094											
3893.50	25676.5	L		28125											
3891.15	25692.0	L		28130											
3889.3'	25704'	L		28159											
3888.5	25710	L		28178											
3886.3	25724	L		28184											
3883.9	25740	L		28207											
3882.8'	25747'	L		28244											
3880.80	25760.5	L		28250											
3878.65	25774.8	L		28283											
3877.45	25783.0	L		28289											
3873.0	25812	L		28640											
3867.9'	25846'	L		28643											
3865.1'	25863'	L		28673											
3860.4'	25897'	L		28673											
3857.2'	25919'	L		28782											
3840.2'	26033'	L		28788											
3834.1'	26074'	L		28812											
26324'	26327	26321.1	2s	H		28845	28845	5d							
26388'	26362	26352.1	1s	H		3787.65	26394.1	3nd							

behavior. Tables II and III give examples of the constant intervals found. The lower levels of the diagram and tables account for all of the important lines of the multiplets considered and for most of the high-temperature lines in other parts of the spectrum. The complexity of the other multiplets seems to be the chief factor hindering attainment of complete agreement.

TABLE II. Prominent "low-temperature" lines and their "high-temperature" satellites (78°K).

L.-t. line	Satellite A	$\Delta\bar{\nu}_A$	Satellite B	$\Delta\bar{\nu}_B$
17847	17809	38	17783	64
17871	17832	39		
18857	18817	40	18790	67
18878	18839	39	18810	68
19987	19949	38	19920	67
20009	19970	39		
20110	20072	38	20044	66
20382	20342	40	20315	67
21509	21472	37	21442	67
21641	21600	41		
22043.0	22003	40	21975	68
22144	22104	40	22076	68
22185			22117	68
22708	22669	39	22642	66
23898.8	23859	40	23830	69
24588.9	24549	40		
26926	26888	38		
27340	27302	38	27273	67
28125	28087	38	28055	70
28244	28207	37	28178	66

The lines at nitrogen temperatures are somewhat more diffuse for the bromate than for the chloride. This probably results from the fact that the lines already have started to split apart. At hydrogen temperature the splitting is so pronounced that at first sight the spectrum seems to present an entirely different appearance. All of the levels evidently have split into two components. This splitting amounts to from 5 to 15 cm^{-1} for the lower levels and is somewhat less for the extreme upper ones. Of the two components resulting from each of the lower levels the one with higher energy appears not to be single (uncertain for the 67–82 pair). It is probably double and may be still more complex. Of particular interest is the fact that after the levels have split all the lines resulting therefrom cannot be referred to common upper levels, except in cases where it is obvious that the low temperature lines in question are complex and are not completely resolved. As a result, the levels 0, 45, and 67 seem to be associated with one set of excited

levels and 0, 37 and 82 with another.⁵ Only a few lines arising from the excited lower level at 82 cm^{-1} are observed, since the lines are very faint because of the low population in this state at 20°K. Because of this low intensity the above rule has not as much confirmation for this level as might be desired, though the behavior described is certainly true of the 37 and 45 levels.

It is not at all certain just what causes this splitting of the levels, though in any case the effect is probably brought about by a slight change in the crystal structure. The symmetry of the field about the Sm^{+++} ion may become less at the lower temperatures, causing the levels to be less degenerate. On the other hand, the selection rules mentioned above suggest that in the crystalline change some of the samarium ions in the unit cell have become different from the others (with respect to field), and that in a mole of ions in the crystal one has a fraction of a mole of one type of Sm^{+++} and the remainder of another kind.⁶

Any crystalline change that might bring about this effect is small, however, as appears from the fact that the bromate crystals show no pronounced tendency to crack when they are suddenly cooled, nor do they undergo crystal transitions during prolonged exposures at low temperatures.

There is some evidence that other lower levels exist. New lines make their appearance at liquid ethylene temperature, though they are quite diffuse and faint. The probabilities of transitions from such levels seem to be small and only a few of them appear. Thicker absorbing layers tend to bring these lines out more strongly, but in any case they are hard to measure accurately (see dotted level and band of Fig. 3). The most

⁵ It will be noticed that the method employed here cannot determine whether the basic level, 0, is or is not split at low temperatures. The diagram, Fig. 3, is based upon the assumption that it is not. If the low-temperature doublets were all of the same separation, one might conclude that the high upper levels are not split and that the basic one has divided this much. However, according to our measurements such cannot be the case, even if there is some splitting in the normal level. This ground state probably does split similarly to the next higher one, though it does so in smaller amount.

⁶ The ions in a crystal of $\text{SmCl}_3 \cdot 6\text{H}_2\text{O}$ must all have the same energy states, since no such selection rules have been found for that salt.

TABLE III. Prominent "low-temperature" lines and their "high-temperature" satellites (20°K).

L.-t. line	A_I	$\Delta\tilde{\nu}_{A_I}$	A_{II}	$\Delta\tilde{\nu}_{A_{II}}$	B_{II}	$\Delta\tilde{\nu}_{B_{II}}$	B_I	$\Delta\tilde{\nu}_{B_I}$
17844.9 (I)	17808.3	36.6						
17847.4 (II)			17802.5	44.9	17780	67		
17875.6 (I, II)	17839.3	36.3	17830	46				
17939 (II)			17894	45				
17949 (I)	17912	37					17867	82
18857.2 (II)			18812	45				
18859.6 (I)	18823.1	36.5						
18878 (II)			18832.9	45				
18882.9 (I)	18846.3	36.5						
19985.6 (I, II)	19949.7	35.9	19941	45	19919	67	19904	82
20007.7 (I, II)	19973.2	34.5	19962	46				
20032.9 (I)	19996.2	36.7						
20108 (II)			20063	45				
20117 (I)	20080	37						
20339 (I, II)	20304.7	34	20293	46	20270	69		
20385.9 (I)	20348.9	37.0						
20485 (I, II)	20448	37	20440	45				
20679 (I, II)	20644	35	20636	43				
21511 (?)			21465	46	21445	66		
22040.7 (IIa)			21998	43	21977	64		
22044.0 (IIb)			21998	46	21977	67		
22048.5 (I)	22011	38					21965	84-
22145 (I, II)	22112	33	22100	45				
22187 (I, II)					22120	67		
22489.3 (I, II)	22453	36	22445	44	22425	64		
22503.5 (I, II)	22467	37	22459	45				
23013 (I, II)			22968	45	22945	68	22931	82
23901.4 (II)			23857.2	44.2				
23904.5 (I)	23867.6	36.9						
24533 (I, II)	24497.9	35	24487	46	24466.0	67		
26579.2 (I, II)	26543	36						
28125.2 (II)			28080	45				
28130.6 (I)	28094	37						
28249 (I, II)	28218	31	28204	45				

The Roman numerals refer to the two components formed at 20° from single lines at 78°.

definite cases of such lines are those at 23,660 to 23,783 (see Table I). Judging from these bands the levels would be separated from the basic one by about 100 to 230 cm^{-1} , probably existing as two doublets with centers at 140 and 220 cm^{-1} .

CONCLUSIONS

The absorption spectra of $\text{Sm}(\text{BrO}_3)_3 \cdot 9\text{H}_2\text{O}$ crystals are composed of lines that originate in transitions from a basic state and other low ones to excited levels which are grouped in multiplets. These multiplets probably arise from excited levels that are split apart in the electric field of the crystal. The over-all splitting of these high

levels is somewhat less than in the monoclinic chloride.

The separations of the lower levels are also less than in the chloride, being chiefly from 39 to 68 cm^{-1} as compared with 145 to 217 cm^{-1} in the chloride, though greater intervals probably also occur in both. The existence of more than three such levels in the samarium salts makes the assumption of more than one low-lying electronic state, as advocated by one of us,⁷ highly probable.⁸

⁷ Spedding, J. Am. Chem. Soc. **54**, 2593 (1932).

⁸ The basic level of Sm^{+++} as predicted by Hund is a ${}^6H_{5/2}$ and should yield only $J+\frac{1}{2}$ or 3 levels in an electric field.

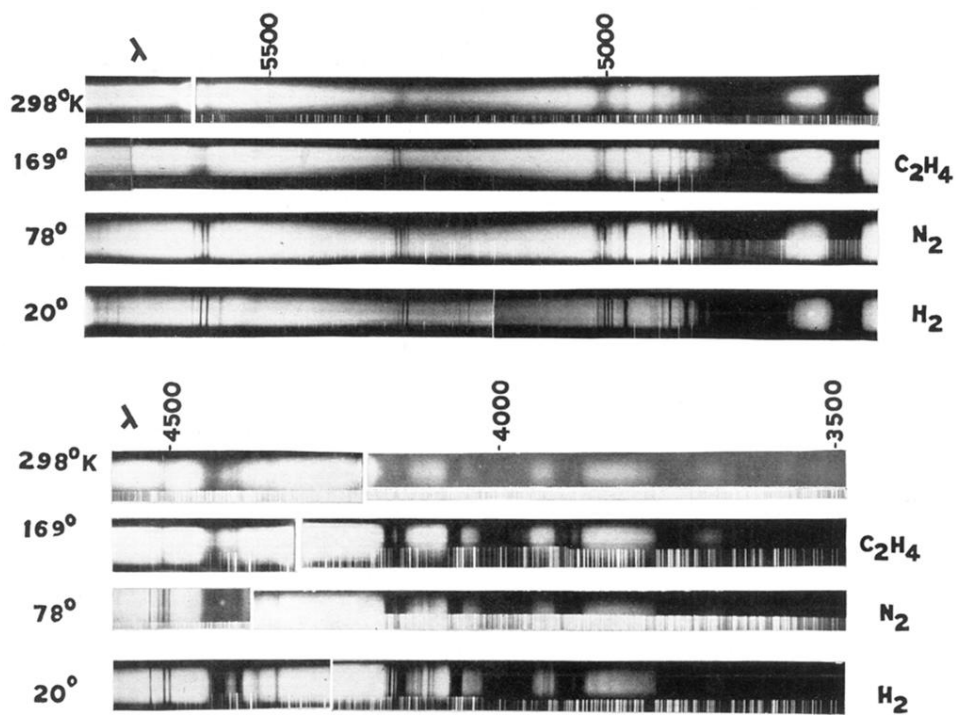


FIG. 1. Conglomerate absorption of $\text{Sm}(\text{BrO}_3)_3 \cdot 9\text{H}_2\text{O}$. The substances at whose boiling points the spectra were photographed are indicated on the right, the corresponding temperatures on the left. The three conspicuously simple multiplets on the long wave-length end are considered in more detail in Figs. 2 and 3.

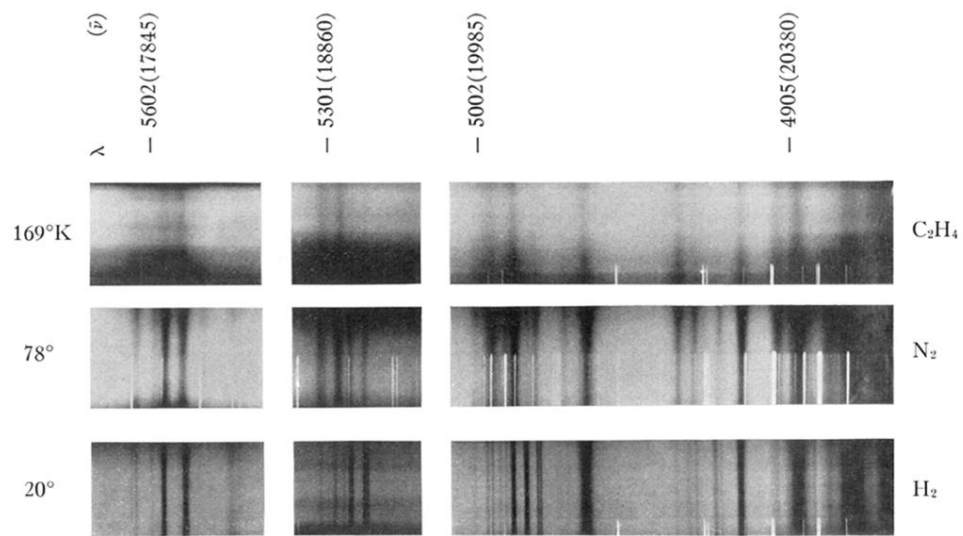


FIG. 2. Selected multiplets of $\text{Sm}(\text{BrO}_3)_3 \cdot 9\text{H}_2\text{O}$ absorption. Several of the multiplets of simplest structure are shown here magnified. The temperatures are given on the left, these being the boiling points of the substances on the right. Above are quoted wave-lengths and wave numbers for purposes of orientation with respect to the other figures and tables.