

SINGLET-TRIPLET INTERVAL RATIOS FOR sp , sd , sf ,
 p^5s AND d^9s CONFIGURATIONS*

By E. U. CONDON AND G. H. SHORTLEY

UNIVERSITY OF MINNESOTA

(Received April 29, 1930)

ABSTRACT

A systematic comparison of the known data on the singlets and triplets arising from sp , sd , sf , p^5s and d^9s configurations with the theory of Houston shows that the theory gives a good account of the deviations from the Landé interval rule which accompany departure from Russell-Saunders coupling. There are numerous significant discrepancies, however. Writing 1L_l and ${}^3L_{l+1}$, 3L_l , ${}^3L_{l-1}$ with $L=P, D, F$, when $l=1, 2, 3$ for the term values, we plot as abscissa $({}^3L_{l-1} - {}^3L_{l+1}) / |{}^3L_l - {}^1L_l|$ and $({}^3L_{l-1} - {}^3L_l) / ({}^3L_l - {}^3L_{l+1})$ as ordinate if $({}^3L_l - {}^1L_l)$ is positive, otherwise the reciprocal of this quantity. Houston's equations (12) give functional relations between these interval ratios which are compared with the experimental values.

HOUSTON¹ has worked out an approximate quantum mechanical theory of the relation of the triplet interval ratio to the singlet-triplet interval for two electron configurations in which one of the electrons is in an s state, but the comparison he makes with experimental data gives one very little idea as to just how accurate the theory is. The purpose of the present paper is therefore to make a systematic comparison of the available data with Houston's theory.

If we write the terms of a singlet-triplet system as 1L_l , ${}^3L_{l+1}$, 3L_l , ${}^3L_{l-1}$, with $L=P, D, F$, when $l=1, 2, 3$, Houston's equations (12) give the following relative term values in terms of the parameter X , which is the ratio of the exchange perturbation integral to the perturbation theory integral which measures the spin energy. (The first classification applies to $X > 0$, the second to $X < 0$):

$$\begin{aligned} {}^1L_l, {}^3L_l &= -\frac{1}{2}(X-1) - \frac{1}{2}\{(X+1)^2 + 4l(l+1)\}^{1/2} \\ {}^3L_{l+1} &= -l \\ {}^3L_l, {}^1L_l &= -\frac{1}{2}(X-1) + \frac{1}{2}\{(X+1)^2 + 4l(l+1)\}^{1/2} \\ {}^3L_{l-1} &= l+1. \end{aligned}$$

Fig. 1. illustrates the behavior of these intervals. With X large and positive the triplet has the Landé interval and the singlet is high above the triplet as at a . As X approaches zero the levels approach coincidence in pairs as at b ; and with X large and negative we again have the Landé interval with the singlet far below the triplet as at c .

* This paper was presented at the Washington meeting of the American Physical Society, April 24, 1930.

¹ W. V. Houston, Phys. Rev. 33, 297 (1929).

The theoretical curves of Figs. 2 and 3 are plotted as follows: For abscissa is used $(^3L_{l-1} - ^3L_{l+1}) / |^3L_l - ^1L_l|$. For the lower curve, $X > 0$, the ordinate is $(^3L_{l-1} - ^3L_l) / (^3L_l - ^3L_{l+1})$, which goes from $l/(l+1)$ with abscissa zero, to zero with abscissa one. For the upper curve, $X < 0$, the ordinate

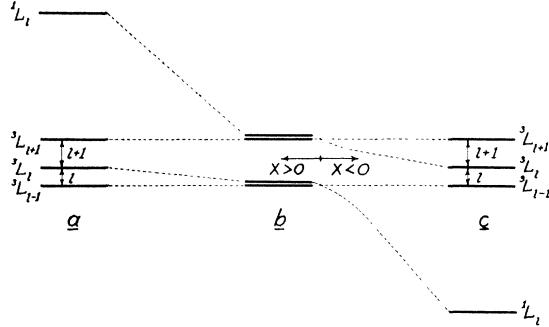


Fig. 1. Diagram illustrating the behavior of the singlet-triplet intervals. (a) X large and positive. (b) X very small and positive. (c) X large and negative.

is the reciprocal of this, $(^3L_l - ^3L_{l+1}) / (^3L_{l-1} - ^3L_l)$, which goes from zero with abscissa one to $(l+1)/l$ with abscissa zero. The fact that the abscissa starts increasing (to $(l+\frac{1}{2})/(l(l+1))^{1/2}$ at $X = -1$) shows that $^1L_l - ^3L_l$ becomes less than the whole spread of the triplet until the ordinate equals $1/2l$ at $X = -2$.

TABLE I. Coordinates.

\dagger indicates a point which is not plotted on account of large probable error (see text).

$*$ indicates a point which does not fall within the limit of Figs. 2 and 3.

Elem.	Config.	$s^p(X > 0)$		Sources
		Abs.	Ord.	
Al II	$3s3p$	0.0084	0.493	Int. Crit. Tab.
	$4p$	0.0293	0.481	
	$5p$	0.1145	0.445	
Ba I	$6s6p$	0.2301	0.422	Int. Crit. Tab.
	$7p$	0.1410	0.420	
	$8p$	0.1181	0.619	
C I	$3s2p$	0.0527	0.499	Proc. Roy. Soc. A118, 43
Ca I	$4s4p$	0.0179	0.494	Int. Crit. Tab.
	$5p^\dagger$	0.1253	0.110	
Cd I	$5s5p$	0.1315	0.462	Int. Crit. Tab.
	$6p$	0.1695	0.406	
	$7p$	0.202	0.355	
	$8p$	0.211	0.378	
	$9p$	0.219	0.367	
Ga II	$4s4p$	0.0604	0.478	Phys. Rev. 34, 714
	$5s4p$	0.1476	0.424	
Ge I	$5s4p$	0.719	0.177	Phys. Rev. 31, 786
	$6s4p$	0.914	0.095	
Ge III	$4s4p$	0.0818	0.465	Phys. Rev. 34, 698
	$5p$	0.276	0.366	
Hg I	$6s6p$	0.436	0.382	Int. Crit. Tab.
	$7p$	1.035	0.094	
	$8p$	0.952	0.056	
	$9p$	0.0046	0.487	Int. Crit. Tab.
	$10p$	0.680	0.235	Proc. Roy. Soc. A114, 662
Mg I	$3s3p$	0.135	0.437	
	$4s2p$	0.305	0.373	
	$5s2p$	0.0658	0.458	Proc. Roy. Soc. A118, 43
Pb I	$7s6p$	0.972	0.032	Phys. Rev. 33, 301
Pb III	$6s6p$	0.600	0.274	Phys. Rev. 34, 397
Sb IV	$5s5p$	0.278	0.388	Phys. Rev. 34, 402
Si I	$4s3p$	0.221	0.395	Zeits. f. Phys. 40, 530
Sn I	$6s5p$	0.918	0.074	Phys. Rev. 30, 574
Sn III	$5s5p$	0.230	0.407	Int. Crit. Tab.
	$6p$	0.598	0.226	
Sr I	$5s5p$	0.0808	0.474	Int. Crit. Tab.
	$6p$	0.607	0.395	

Elem.	Config.	$s^p(X < 0)$		Sources
		Abs.	Ord.	
Al II	$3s6p^*$	0.1045	2.193	Int. Crit. Tab.
	$4s6p$	0.0140	1.950	Int. Crit. Tab.
	$5s6p$	0.3935	5.356	Int. Crit. Tab.
Ca I	$4s6p$	0.1273	7.135	
	$5s6p$	0.1131	8.73	
Hg I	$10p^*$	0.0902	2.100	Int. Crit. Tab.
Sr I	$5s7p^*$	0.814	0.823	Phys. Rev. 31, 501
	$5s^*$	1.109	0.269	
	$5s$	1.063	0.241	

Elem.	Config.	$p^6s(X < 0)$		Sources
		Abs.	Ord.	
A I	$3p^6s$	0.854	0.756	Zeits. f. Phys. 40, 839
	$6s$	1.008	0.0578	
	$7s$	1.021	0.0282	
Ca III	$3p^6s$	0.95	0.0078	
	$5s$	1.002	0.0078	
	$9s$	1.109	0.269	
K II	$3p^6s$	0.819	0.383	Phys. Rev. 31, 501
	$5s$	1.063	0.241	
Kr I	$4p^6s$	1.058	0.221	B. S. Jl. Res. 3, 154
Na II	$2p^6s$	0.442	1.293	Phys. Rev. 31, 967

TABLE I (continued)

Ne I	$2p^63s$	0.544	1.161	Int. Crit. Tab.	Tl II	$6s6d$	0.532	1.384	Phys. Rev. 35, 236
	5s	1.047	0.122			7d	0.237	1.173	
	6s	1.027	0.0578			8d*	0.118	1.861	
	7s	1.017	0.0322		Zn I	$4s4d$	0.0283	1.618	Int. Crit. Tab.
	8s	1.010	0.0211			5d†	0.0128	1.818	
	9s	1.006	0.0125		Yt II	$5s5d$	0.0533	1.509	B. S. Jl. Res. 2, 738
	10s	1.003	0.0084						
	11s	1.001	0.0061						
Xe I	$5p^66s$	0.999	0.120	B. S. Jl. Res. 3, 756					
<hr/>									
$sd(X>0)$									
Elem.	Config.	Abs.	sd (X>0)	Ord.					
Ba I	$6s5d$	0.258	0.476	Int. Crit. Tab.	Ag II	$4d^95s$	0.862	0.526	Phys. Rev. 31, 317
Ca I	$4s3d$	0.0237	0.640	Int. Crit. Tab.		6s	1.022	0.0896	
	5d†	0.0258	0.607		Au II	$5d^66s$	1.062	0.2569	Phys. Rev. 34, 19
	7d†	0.712	0.750		Cd III	$4d^95s$	0.886	0.4916	Phys. Rev. 31, 778
Cr V	$4s3d$	0.1678	0.525	Phys. Rev. 33, 542		3d ⁹ s	0.605	0.798	Phys. Rev. 29, 386
Lu II	$6s5d$	0.490	0.363	Bul. Am. Ph. Soc. Apr. 10, 1930, 11		5s	1.018	0.1835	
³ Pb III	$6s5d$	0.220	0.466	Phys. Rev. 34, 397		6s	1.018	0.0692	Phys. Rev. 34, 1128
Sb IV	$5s5d$	0.258	0.628	Phys. Rev. 34, 402	Ga IV	$3d^94s$	0.708	0.686	Phys. Rev. 31, 750
	6d	0.562	0.683		Ge V	$3d^94s$	0.759	0.622	Phys. Rev. 31, 750
Sc II	$4s3d$	0.0718	0.616	Sci. Papers, Bur. Stand. 22, 329	Hg III	$5d^66s$	1.033	0.2568	Phys. Rev. 34, 19
	5s3d	0.301	0.483		In IV	$4d^95s$	0.913	0.447	Phys. Rev. 31, 778
Sn III	$5s5d$	0.250	0.659	Int. Crit. Tab.	Ni I	$3d^94s$	0.596	0.810	Phys. Rev. 29, 386
	6d	0.129	0.610	Phys. Rev. 34, 402		5s	1.024	0.1393	
Sr I	$5s4d$	0.083	0.600	Int. Crit. Tab.		6s	1.016	0.0519	Phys. Rev. 34, 828
	6d	0.409	0.398		Pd I	$4d^95s$	0.890	0.510	Phys. Rev. 29, 386
Te V	$5s5d$	0.257	0.616	Phys. Rev. 34, 402		6s	1.018	0.065	
Tl III	$4s3d$	0.1031	0.592	Astro. Jl. 66, 13	Pt I	$5d^66s$	0.796	0.0829	Phys. Rev. 34, 19
V IV	$4s3d$	0.1335	0.558	Phys. Rev. 33, 542	Tl IV	$5d^66s$	1.042	0.235	Phys. Rev. 34, 19
Yt II	$5s4d$	0.272	0.506	B. S. Jl. Res. 2, 738	Zn III	$3d^94s$	0.652	0.748	Phys. Rev. 30, 381
	6s4d	0.996	0.124						
<hr/>									
$sd(X<0)$									
Elem.	Config.	Abs.	sd (X<0)	Ord.					
Ba I	$6s6d$	0.0485	1.228	Int. Crit. Tab.	Al II	$3s4f$	0.0978	0.750	Int. Crit. Tab.
	7d	0.0456	0.438			5f	0.0506	0.783	
	8d	0.0284	1.190			6f	0.0573	0.777	
	9d	0.0746	1.595			7f	0.0389	0.761	
	10d†	0.0324	4.61			8f	0.0118	0.657	
Ca I	$4s4d$	0.0207	1.474	Int. Crit. Tab.		9f†	0.0049	0.750	
	6d†	0.0544	1.200			10f†	0.0030	0.786	
Cd I	$5s5d$	0.1074	1.555	Int. Crit. Tab.		11f†	0.0016	0.667	
	6d	0.0621	1.414			12f†	0.0019	0.800	
	7d	0.0553	1.388		⁴ Ge III	$4s4f$	1.384	0.275	Phys. Rev. 34, 697
	8d†	0.0454	6.000		Sn III	$5s4f$	0.0556	0.378	Int. Crit. Tab.
Ga II	$4s4d$	0.0096	1.360	Phys. Rev. 34, 714		5f	0.0312	0.357	Phys. Rev. 30, 574
	5s4d	0.0029	1.416		Sr I	$5s4f$	0.0056	0.630	Int. Crit. Tab.
	6s4d	0.0022	1.125			5f†	0.0113	0.889	
Ge III	$4s4d$	0.0099	1.505	Phys. Rev. 34, 697		6f†	0.1765	0.049	
Hg I	$6s6d$ *	1.510	0.585	Int. Crit. Tab.		7f†	0.314	0.165	
	7d	1.018	0.923						
	8d	0.808	1.051						
	9d	0.717	1.438						
	10d	0.653	0.556						
³ Pb III	$6s5d$ *	0.250	2.145	Phys. Rev. 34, 397					
Sr I	$5s5d$	0.183	1.520	Int. Crit. Tab.					
	7d	0.273	1.195						
	8d	0.163	1.250						
	9d	0.298	1.271						

⁴ S. Smith, Phys. Rev. 34, 397 (1929) has a question as to which of two singlets belong to this configuration. These are both listed, one for $X > 0$, plotted with a question mark and one for $X < 0$, which fell off the graph.

⁶ The singlet is within the triplet.

A fairly complete search of the literature was made and the points plotted on Figs. 2 and 3. The coordinates are tabulated in the accompanying tables together with a brief reference to the source of the data. No points were plotted for which any of the intervals were less than 2.5 cm^{-1} , since the accuracy of these points did not seem to be sufficient for a fair comparison with the theory. Such points are listed in the table with a dagger (\dagger). Those points

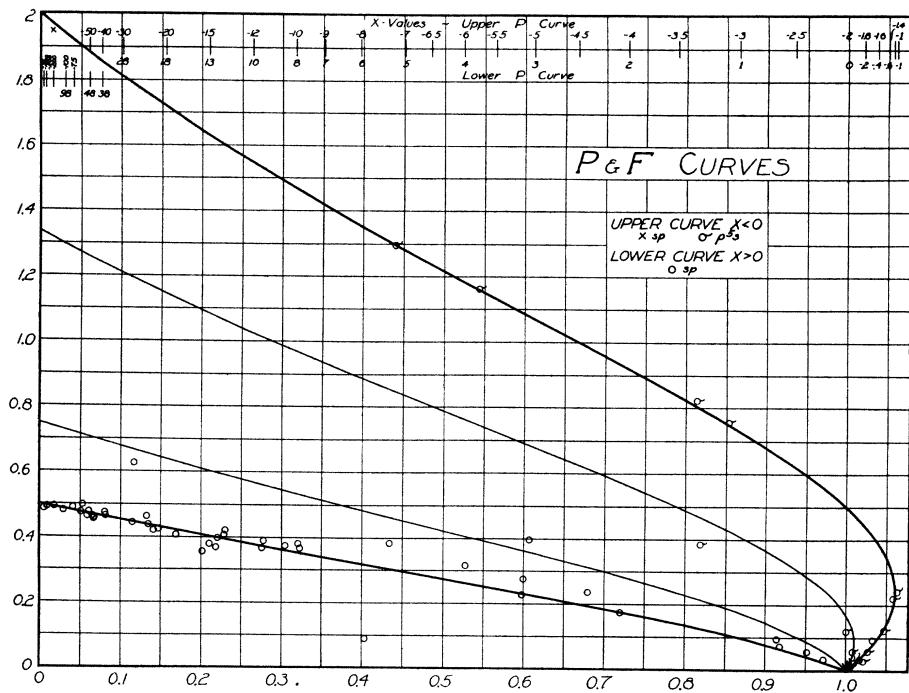


Fig. 2. P and F curves showing sp and ps points. (F curve is drawn lightly and no points plotted.)

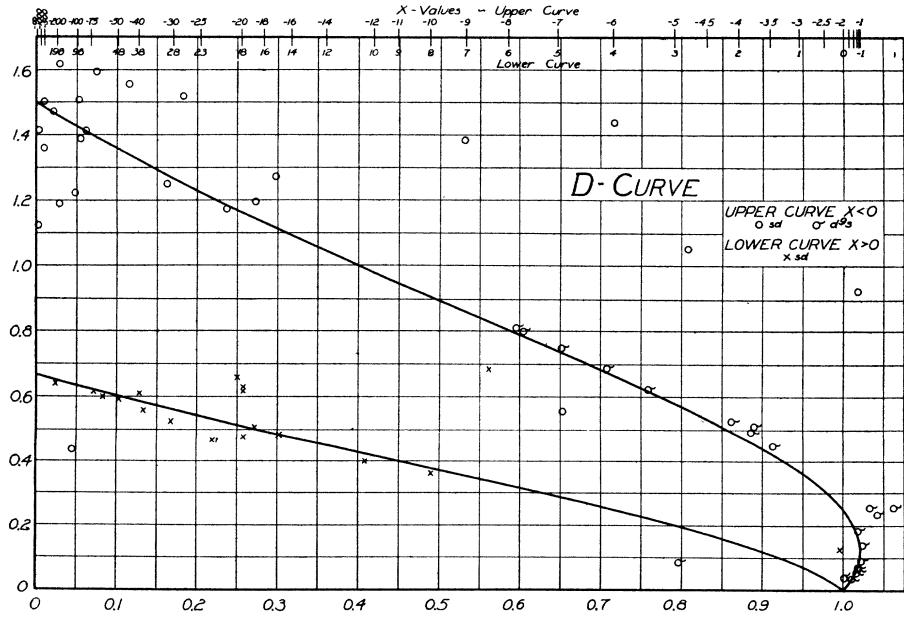


Fig. 3. D curve showing sd and d^9s points.

whose coordinates fell off the graph are listed with an asterisk (*). No *sf* points were plotted because there were just a few clustered near Russell-Saunders coupling.

It is found that the theory works well for p^5s and d^9s configurations, in which we are one p or d electron short of a closed shell, except that the whole system of lines is inverted. These points all fall on the upper curves.

Pb III² and Tl II³ have *sf* configurations in which 3F_3 and 1F_3 are on opposite sides and outside of the ${}^3F_2 - {}^3F_4$ interval. A similar thing happens in the case of the $3p^5 8s$ configuration of A I⁴. These partially inverted triplets might have been plotted with negative ordinates on our graphs.

As to the accuracy with which the points fit the theory, it can be said that in general the d^9s and p^5s fit best, sp and sd for $X > 0$ next, and sp and sd for $X < 0$ poorest. Also in general, elements of high atomic number show especially pronounced disagreements. The disagreements are to be regarded provisionally as cases in which the second-order perturbations are not negligible rather than as essential defects in the basic theory.

² S. Smith, Phys. Rev. **34**, 397 (1929).

³ S. Smith, Phys. Rev. **35**, 236 (1930).

⁴ A. Meissner, Zeits. f. Physik **40**, 839 (1927).