

SINGLET-TRIPLET INTERVAL RATIOS FOR sp , sd , sf ,
 p^2s AND d^2s CONFIGURATIONS*

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ABSTRACT

A systematic comparison of the known data on the singlets and triplets arising from sp , sd , sf , p^2s and d^2s configurations with the theory of Houston shows that the theory gives a good account of the deviations from the Landé interval rule which accompany departure from Russell-Saunders coupling. There are numerous significant discrepancies, however. Writing 1L_l and ${}^3L_{l+1}$, 3L_l , ${}^3L_{l-1}$ with $L = P, D, F$, when $l = 1, 2, 3$ for the term values, we plot as abscissa $({}^3L_{l-1} - {}^3L_{l+1}) / |{}^3L_l - {}^1L_l|$ and $({}^3L_{l-1} - {}^3L_l) / ({}^3L_l - {}^3L_{l+1})$ as ordinate if $({}^3L_l - {}^1L_l)$ is positive, otherwise the reciprocal of this quantity. Houston's equations (12) give functional relations between these interval ratios which are compared with the experimental values.

HOUSTON¹ has worked out an approximate quantum mechanical theory of the relation of the triplet interval ratio to the singlet-triplet interval for two electron configurations in which one of the electrons is in an s state, but the comparison he makes with experimental data gives one very little idea as to just how accurate the theory is. The purpose of the present paper is therefore to make a systematic comparison of the available data with Houston's theory.

If we write the terms of a singlet-triplet system as 1L_l , ${}^3L_{l+1}$, 3L_l , ${}^3L_{l-1}$, with $L = P, D, F$, when $l = 1, 2, 3$, Houston's equations (12) give the following relative term values in terms of the parameter X , which is the ratio of the exchange perturbation integral to the perturbation theory integral which measures the spin energy. (The first classification applies to $X > 0$, the second to $X < 0$.):

$$\begin{aligned} {}^1L_l, {}^3L_l &= -\frac{1}{2}(X-1) - \frac{1}{2}\{(X+1)^2 + 4l(l+1)\}^{1/2} \\ {}^3L_{l+1} &= -l \\ {}^3L_l, {}^1L_l &= -\frac{1}{2}(X-1) + \frac{1}{2}\{(X+1)^2 + 4l(l+1)\}^{1/2} \\ {}^3L_{l-1} &= l+1. \end{aligned}$$

Fig. 1. illustrates the behavior of these intervals. With X large and positive the triplet has the Landé interval and the singlet is high above the triplet as at a . As X approaches zero the levels approach coincidence in pairs as at b ; and with X large and negative we again have the Landé interval with the singlet far below the triplet as at c .

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¹ W. V. Houston, Phys. Rev. **33**, 297 (1929).

TABLE I (continued)

Ne I	2p ³ 3s	0.544	1.161	Int. Crit. Tab.
	5s	1.047	0.122	
	6s	1.027	0.0578	
	7s	1.017	0.0322	
	8s	1.010	0.0211	
	9s	1.006	0.0125	
	10s	1.003	0.0084	
	11s	1.001	0.0061	
Xe I	5p ⁶ 6s	0.999	0.120	B. S. JI. Res. 3, 756

Elem.	Config.	Abs.	sd (X > 0) Ord.	Sources
Ba I	6s5d	0.258	0.476	Int. Crit. Tab.
Ca I	4s3d	0.0237	0.640	Int. Crit. Tab.
	5d†	0.0258	0.607	
	7d†	0.712	0.750	
Cr V	4s3d	0.1678	0.525	Phys. Rev. 33, 542
Lu II	6s5d	0.490	0.363	Bul. Am. Ph. Soc. Apr. 10, 1930, 11
²⁰⁸ Pb III	6s5d	0.220	0.466	Phys. Rev. 34, 397
Sb IV	5s5d	0.258	0.628	Phys. Rev. 34, 402
	6d	0.562	0.683	
Sc II	4s3d	0.0718	0.616	Sci. Papers, Bur. Stand. 22, 329
	5s3d	0.301	0.483	
Sn III	5s3d	0.250	0.659	Int. Crit. Tab.
	6d	0.129	0.610	Phys. Rev. 34, 402
Sr I	5s4d	0.083	0.600	Int. Crit. Tab.
	6d	0.409	0.398	
Te V	5s5d	0.257	0.616	Phys. Rev. 34, 402
Tl III	4s3d	0.1031	0.592	Astro. JI. 66, 13
V IV	4s3d	0.1335	0.558	Phys. Rev. 33, 542
Yt II	5s4d	0.272	0.506	B. S. JI. Res. 2, 738
	6s4d	0.996	0.124	

Elem.	Config.	Abs.	sd (X < 0) Ord.	Sources
Ba I	6s6d	0.0485	1.228	Int. Crit. Tab.
	7d	0.0456	0.438	
	8d	0.0284	1.190	
	9d	0.0746	1.595	
	10d†	0.0324	4.61	
Ca I	4s4d	0.0207	1.474	Int. Crit. Tab.
	6d†	0.0544	1.200	
Cd I	5s5d	0.1074	1.555	Int. Crit. Tab.
	6d	0.0621	1.414	
	7d	0.0553	1.388	
	8d†	0.0454	6.000	
Ga II	4s4d	0.0096	1.360	Phys. Rev. 34, 714
	5s4d	0.0029	1.416	
	6s4d	0.0022	1.125	
Ge III	4s4d	0.0099	1.505	Phys. Rev. 34, 697
Hg I	6s6d*	1.510	0.585	Int. Crit. Tab.
	7d	1.018	0.923	
	8d	0.808	1.051	
	9d	0.717	1.438	
	10d	0.653	0.556	
²⁰⁸ Pb III	6s5d*	0.250	2.145	Phys. Rev. 34, 397
Sr I	5s5d	0.183	1.520	Int. Crit. Tab.
	7d	0.273	1.195	
	8d	0.163	1.250	
	9d	0.298	1.271	

Elem.	Config.	Abs.	d ^{2s} (X < 0) Ord.	Sources
Tl II	6s6d	0.532	1.384	Phys. Rev. 35, 236
	7d	0.237	1.173	
	8d*	0.118	1.861	
Zn I	4s4d	0.0283	1.618	Int. Crit. Tab.
	5d†	0.0128	1.818	
Yt II	5s5d	0.0533	1.509	B. S. JI. Res. 2, 738

Elem.	Config.	Abs.	sf (X > 0) (not plotted) Ord.	Sources
Ag II	4d ⁹ 5s	0.862	0.526	Phys. Rev. 31, 317
	6s	1.022	0.0896	
Au II	5d ⁹ 6s	1.062	0.2569	Phys. Rev. 34, 19
Cd III	4d ⁹ 5s	0.886	0.4916	Phys. Rev. 31, 778
Cu II	3d ⁹ 4s	0.605	0.798	Phys. Rev. 29, 386
	5s	1.018	0.1835	
	6s	1.018	0.0692	Phys. Rev. 34, 1128
	7s	1.010	0.0348	
Ga IV	3d ⁹ 4s	0.708	0.686	Phys. Rev. 31, 750
Ge V	3d ⁹ 4s	0.759	0.622	Phys. Rev. 31, 750
Hg III	5d ⁹ 6s	1.033	0.2568	Phys. Rev. 34, 19
In IV	4d ⁹ 5s	0.913	0.447	Phys. Rev. 31, 778
Ni I	3d ⁹ 4s	0.596	0.810	Phys. Rev. 29, 386
	5s	1.024	0.1393	
	6s	1.016	0.0519	Phys. Rev. 34, 828
Pd I	4d ⁹ 5s	0.890	0.510	Phys. Rev. 29, 386
	6s	1.018	0.065	
Pt I	5d ⁹ 6s	0.796	0.0829	Phys. Rev. 34, 19
	7s	1.001	0.0374	Phys. Rev. 34, 190
Tl IV	5d ⁹ 6s	1.042	0.235	Phys. Rev. 34, 19
Zn III	3d ⁹ 4s	0.652	0.748	Phys. Rev. 30, 381

Elem.	Config.	Abs.	sf (X > 0) (not plotted) Ord.	Sources
Al II	3s4f†	0.0978	0.750	Int. Crit. Tab.
	5f	0.0506	0.783	
	6f	0.0573	0.777	
	7f	0.0389	0.761	
	8f	0.0118	0.657	
	9f†	0.0049	0.750	
	10f†	0.0030	0.786	
	11f†	0.0016	0.667	
	12f†	0.0019	0.800	
⁶⁸ Ge III	4s4f	1.384	0.275	Phys. Rev. 34, 697
Sn III	5s4f	0.0556	0.378	Int. Crit. Tab.
	5f	0.0312	0.357	Phys. Rev. 30, 574
Sr I	5s4f†	0.0056	0.630	Int. Crit. Tab.
	5f†	0.0113	0.889	
	6f†	0.1765	0.049	
	7f†	0.314	0.165	

Elem.	Config.	Abs.	sf (X < 0) (not plotted) Ord.	Sources
Ba I	6s4f	0.0046	1.014	Int. Crit. Tab.
	5f	0.0847	4.34	
	6f	0.0092	1.594	

* S. Smith, Phys. Rev. 34, 397 (1929) has a question as to which of two singlets belong to this configuration. These are both listed, one for $X > 0$, plotted with a question mark and one for $X < 0$, which fell off the graph.

† The singlet is within the triplet.

A fairly complete search of the literature was made and the points plotted on Figs. 2 and 3. The coordinates are tabulated in the accompanying tables together with a brief reference to the source of the data. No points were plotted for which any of the intervals were less than 2.5 cm^{-1} , since the accuracy of these points did not seem to be sufficient for a fair comparison with the theory. Such points are listed in the table with a dagger (†). Those points

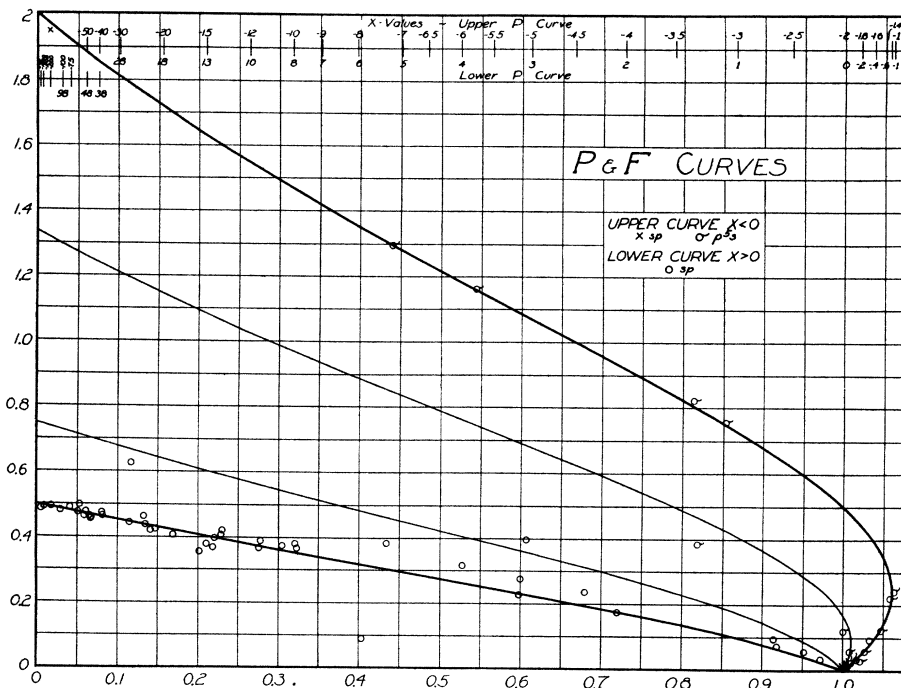


Fig. 2. *P* and *F* curves showing *sp* and *p^s* points. (*F* curve is drawn lightly and no points plotted.)

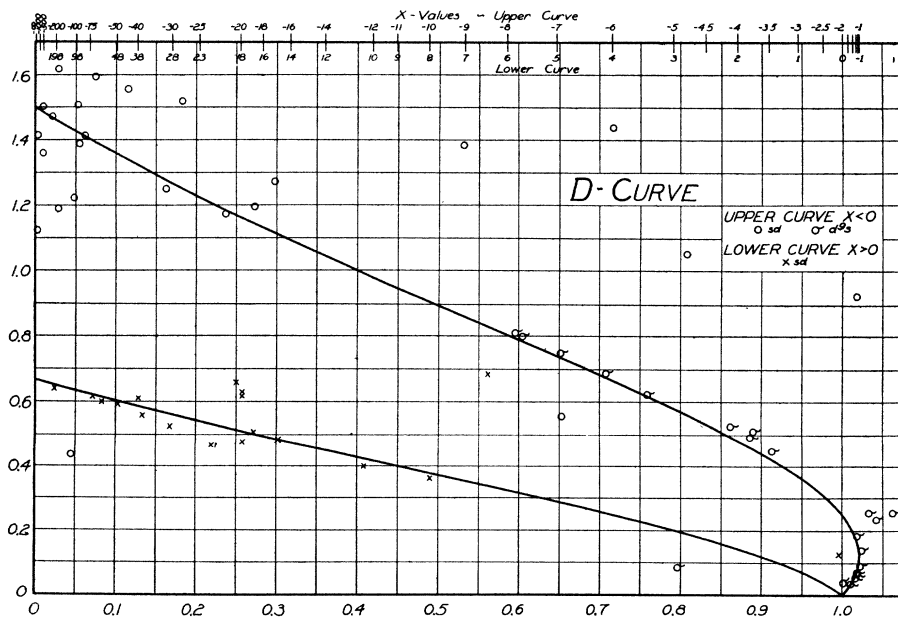


Fig. 3. *D* curve showing *sd* and *d^s* points.

whose coordinates fell off the graph are listed with an asterisk (*). No sf points were plotted because there were just a few clustered near Russell-Saunders coupling.

It is found that the theory works well for p^5s and d^9s configurations, in which we are one p or d electron short of a closed shell, except that the whole system of lines is inverted. These points all fall on the upper curves.

Pb III² and Tl II³ have sf configurations in which 3F_3 and 1F_3 are on opposite sides and outside of the $^3F_2 - ^3F_4$ interval. A similar thing happens in the case of the $3p^5 8s$ configuration of A I⁴. These partially inverted triplets might have been plotted with negative ordinates on our graphs.

As to the accuracy with which the points fit the theory, it can be said that in general the d^9s and p^5s fit best, sp and sd for $X > 0$ next, and sp and sd for $X < 0$ poorest. Also in general, elements of high atomic number show especially pronounced disagreements. The disagreements are to be regarded provisionally as cases in which the second-order perturbations are not negligible rather than as essential defects in the basic theory.

² S. Smith, Phys. Rev. **34**, 397 (1929).

³ S. Smith, Phys. Rev. **35**, 236 (1930).

⁴ A. Meissner, Zeits. f. Physik **40**, 839 (1927).