Liquid-Solid Phase Transition in  $He<sup>3</sup>-He<sup>4</sup>$  Mixtures, P. M. Tedrow and D. M. Lee [Phys. Rev. 181, 399 (1969)]. Paragraphs on pp. 410 and 411 were inadvertently transposed. For convenience, both corrected pages are herewith reproduced in their entirety, and should be substituted for pp. 410- 411 in the issue of note.

....three phases in accordance with the discussion of the phase rule. Figure 13 resembles an ordinary eutectic diagram with two liquids and one solid rather than two solids and one liquid coexisting.

Figure 14 is drawn for a pressure of 25.6 atm. The question of what happens as the He'-rich hexagonal close-packed (hcp) solid begins to form at about 25.0 atm is related to the melting curve minimum in pure  $He<sup>4</sup>$ . This is treated in detail minimum in pure He<sup>4</sup>. This is treated in detail<br>by LePair *et al*.<sup>11</sup> and will not be discussed here. At 25.6 atm, the hcp solid stretches along the left side of the diagram and has isolated a pocket of liquid, L,. Two more univariants have appeared, hcp-L<sub>1</sub>-L<sub>2</sub> and hcp-bcc<sub>1</sub>-L<sub>1</sub>. The univariant hcp $bcc<sub>1</sub>$ -L<sub>1</sub> appears twice because it is thought to have a minimum in the  $P-T$  plane, and the plane  $P=25.6$ atm cuts it at two different temperatures. The projections of the hcp- $L_1$ - $L_2$  and bcc<sub>1</sub>- $L_1$ - $L_2$  univariants in Fig. 14 onto the P-T plane are shown in Fig. 5. These figures show that  $\text{hcp-bcc}_{1}$ -L, and bcc<sub>1</sub>-L<sub>1</sub>-L<sub>2</sub> move to lower temperatures as P



FIG. 15. At 26 atm and  $0.37\degree$ K, the four phases, hcp solid, bcc<sub>1</sub> solid, liquid L<sub>1</sub>, and liquid L<sub>2</sub> are in equilibrium, i. e. , a quadrupole point exists at this temperature and pressure.  $\bigcirc$  This work;  $\nabla$  Vignos and Fairbank<sup>6</sup>;  $\Box$  Le Pair et al.<sup>11</sup>;  $\Delta$  Zinovieva<sup>22</sup>;  $\blacktriangle$  Lipschultz et al. $30$ 

increases while  $\text{hcp-L}_1 - \text{L}_2$  moves to higher temperatures.

In Fig. 15 at 26 atm the three univariants in Fig. 14 have met at about  $0.37\textdegree$ K. At higher pressures, no He<sup>4</sup>-rich liquid  $L_1$  can exist. This point is the quadruple point labeled  $Q<sub>2</sub>$  deduced from the experimental results earlier in the paper [Sec. IV(B)] corresponding to the coexistence of four phases, hcp, bcc<sub>1</sub>,  $L_1$ , and  $L_2$ . The four intersecting univariants associated with this quadruple point are  $bcc_1-L_1-L_2$ ,  $hc p-L_1-L_2$ ,  $bcc_1-hcp-L_1$ , and  $bcc_1-hcp-L_2.$ 

The interpretation of the experimental results obtained in this work in terms of a quadruple point is supported by observations by Zinovieva<sup>22</sup> who visually detected the stratification of an 86.3% mixture into two liquid phases at temperatures of 0.30 and 0.35'K and a pressure of 26 atm with no solid being present. These observations agree qualitatively with the data in Fig. 5 and with the proposed phase diagram presented here if her observations were made in the two-phase region between hcp- $L_1$ - $L_2$  and bcc<sub>1</sub>- $L_1$ - $L_2$  (Fig. 14). In the same manner that a triple point is useful in establishing a fixed temperature point in a single-phase system, a quadruple point in the  $He<sup>3</sup>-He<sup>4</sup>$  system provides a fixed temperature point below  $1^{\circ}$ K, which is not a function of pressure or impurity content.

The location of the line  $bcc_1-hcp-L_2$  is shown in Fig. 16 at 26.2 atm. This line also appears in Fig. 5. In the high-temperature region of the diagram, the univariant hcp-bcc<sub>1</sub>-L<sub>1</sub> has disappeared and the boundaries of the two solid phases have met at the  $X = 0$  axis. This point corresponds to the lower triple point<sup>6</sup> in pure  $He<sup>4</sup>$ .

In Fig. 17 at 26.5 atm, the solid-phase separation' first appears and is thought to be a phase separation into two bcc phases labeled bcc, and  $bcc<sub>2</sub>$ . This gives rise to another univariant, bcc, $bcc<sub>2</sub>-L<sub>2</sub>$ . CP marks the critical point of the solidphase separation. Along the pure  $He<sup>4</sup>$  axis, the thin region of bcc solid<sup>6</sup> in pure He<sup>4</sup> is in evidence at about  $1.5^{\circ}$ K. The data points in this figure belong to the  $bcc_1 - hcp - L_2$  line and not to the  $bcc_1$  $bcc_2-L_2$  univariant. The projection of the two univariants onto the  $P-T$  plane is shown in Fig. 20.

At 26.<sup>75</sup> atm, in Fig. 18, the univariants bcc,  $bcc_2-L_2$  and  $bcc_1-hcp-L_2$  have met, giving rise to a second quadruple point, with the four phases  $bcc_1$ ,  $bcc_2$ , hcp, and  $L_2$  coexisting.

Again there must be four univariants meeting at this quadruple point. The other two,  $bcc_1-bcc_2$ hcp and  $bcc_2$ -hcp-L<sub>2</sub>, are shown in Fig. 19 at 30 atm. These surfaces, too, are shown in Fig. 20, projected onto the  $P-T$  plane. At 30 atm, pure He' has begun to solidify, isolating a small pocket of  $L_2$  in the lower right-hand corner. This is re-

lated to the minimum in the He<sup>3</sup> freezing curve.

The location of all of the univariant surfaces is summarized in Fig. 20 where their projections on the  $P-T$  plane are shown. The parts of the surfaces which have been measured in this laboratory have data points shown.

A few words must be said concerning the extrapolation of some of the univariants beyond the region where they have been measured. The ways of terminating univariants were mentioned earlier in this section. They can terminate at critical point lines or pure substance lines, and at quadruple points where four must meet.

The  $bcc_1-hcp-L_2$  univariant has been drawn so that it meets the pure He<sup>3</sup> melting curve at high pressure and low temperature. This must be the case if the hcp solid is essentially inert and does not affect the situation at all. Then nearly pure solid He<sup>3</sup> (bcc<sub>1</sub>) and nearly pure liquid He<sup>3</sup> ( $L_2$ ) are what determine the path of  $bcc_1-hcp-L_2$ . Since pure liquid He<sup>3</sup> and pure solid He<sup>3</sup> are in equilibrium along the He<sup>3</sup> melting curve, the behavior shown in Fig. 20 is reasonable if the hcp solid is pure He4 at absolute zero.

The hcp- $L_1$ - $L_2$  curve is shown merging into the pure He4 melting curve as the temperature approaches zero. If Fig. 14 is correct, then as the pressure decreases from 25.6 atm and the hcp- $L_1$ - $L_2$  line moves toward lower temperatures,  $L_1$  and  $L_2$  become purer until at absolute zero the pure hcp solid will be in equilibrium with the pure He4 and pure He' liquids assuming complete phase separation at absolute zero. In this condition the solid will be at the bottom of a container with the  $He<sup>4</sup>$  liquid above it and the  $He<sup>3</sup>$  liquid on top. If phase separation is complete at absolute zero, He' atoms cannot spontaneously pass into either liquid  $He<sup>4</sup>$  or solid  $He<sup>4</sup>$  whereas the  $He<sup>4</sup>$  atoms can move freely from liquid to solid. Thus it seems reasonable that the three-phase equilibrium curve hcp- $L_1$ -L, should extrapolate to the He<sup>4</sup> melting curve. If the liquid phase separation is incomplete at absolute zero, as is the case for saturated vapor pressures<sup>3,4</sup> that is, if  $L_1$  can have finite concentration of He' at absolute zero, then the situation can be quite different. If the hcp solid meets the liquid phase separation surface in such a way that it "traps" a region of  $L_1$  on the low-temperature side, then the univariant hcp- $L_1$ - $L_2$  will have a minimum. Figure 21 shows a series of  $T-X$  diagrams for incomplete phase separation which leads to a minimum in the hcp- $L_1$ - $L_2$  line on the P-T diagram.

The other univariants end more routinely. Figures 11 and 12 show why bcc<sub>1</sub>-L<sub>1</sub>-L<sub>2</sub> must end at the peak of the liquid phase separation. The termination of  $bcc_1-bcc_2-L_2$  at the critical point

of the solid phase separation can be understood by noticing how Fig. 17 evolves from Fig. 16. Less is known about the termination of  $bcc_1-bcc_2$ hcp. It is thought that the critical curve (labeled CP) for the solid phase separation is nearly independent of pressure.<sup>5</sup> As the pressure is increased the hcp-bcc phase transition occurs at higher He<sup>3</sup> concentrations. As the hcp-bcc transition moves across toward higher He' concentrations, it carries  $bcc_1 - bcc_2 - hcp$  with it until it hits the peak of the curve, CP in Fig. 19. At this point on the  $P-T$  diagram bcc<sub>1</sub>-bcc<sub>2</sub>-hcp should terminate. Very little is known about the solid phase separation, however, so this statement lacks verification. Finally, the  $bcc_1-hcp-L$ , line must end at the pure He4 melting curve at the lower triple point. In Fig. 20, points shown on the bcc,-hcp-L, univariant were obtained by Lipschultz  $et$  al. in ultrasonic experiments.<sup>30</sup> From the shape of the  $T-X$  diagrams in Figs. 10-19, the location of the azeotropic line on the  $P-T$  diagram is also clear.

Large portions of the phase diagram and the existence of many of the three-phase surfaces have not been verified experimentally suggesting the need for further experiments.

Alternative diagrams have been suggested by Alternative diagrams have been suggested by several other groups.  $11,22,31$  The diagram proposed by LePair  $et$   $al.^{11}$  mainly agrees with the one 11<br>11 presented here. In addition, they discuss the. ...



FIG. 16. At 26.2 atm only the univariant  $bcc_1-hcp-L_2$ remains.  $L_1$  cannot exist at this pressure. The lower triple point of pure He<sup>4</sup> is at about 1.45°K.  $6 \text{ }$  O.This work;  $\nabla$  Vignos and Fairbank<sup>6</sup>;  $\square$  Le pair et al.<sup>11</sup>;  $\triangle$  Zinovieva<sup>22</sup>; A Lipschultz et al.<sup>30</sup>

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