

## Comments and Addenda

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### Consistency of Low-Energy Parameters and Soft-Pion Scattering Theory\*

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The parameters of low-energy  $\pi\pi$ ,  $\pi N$ , and  $\pi N \rightarrow \pi\pi N$  scattering and their consistency with current-algebra predictions are discussed. It is found that the isospin-zero  $\pi\pi$  scattering length is  $(0.24 \pm 0.03)m_\pi^{-1}$  quite independent of the exact value of  $f_\pi$ . An effective value of  $f_\pi$  near 85 MeV seems to be preferred. Experiments are suggested which will clarify the situation further.

THE application of soft-pion techniques to scattering processes has always been subject to the uncertainty in the pion decay constant  $f_\pi$ . When evaluated directly from the pion decay lifetime<sup>1</sup>  $f_\pi = 94$  MeV is obtained, but when  $f_\pi$  is related to  $\pi N$  scattering through the Goldberger-Treiman (GT) relation we get  $f_\pi = (M/G)g_A/g_V = 82$  MeV. The purpose of this paper is to reexamine this problem and some related ones in the light of recently published work.

The unsubtracted, antisymmetric  $\pi-N$  dispersion relation evaluated<sup>2</sup> using the most recent experimental data gives the following relation between the isospin scattering and the Born term:

$$a_{1/2} - a_{3/2} = 5.25 f^2 - (0.122 \pm 0.002). \quad (1)$$

In addition to the usually accepted value of the  $\pi N N$  coupling constant<sup>3</sup>  $f^2 = 0.081 \pm 0.003$ , independent determinations of  $f^2 = 0.079 \pm 0.002$  and  $f^2 = 0.076 \pm 0.004$  are obtained from one-pion exchange (OPE) fits to  $N-N$  data.<sup>4</sup> Using the sum rule (1), a value for the antisymmetric scattering length  $a_{1/2} - a_{3/2} = 0.29 \pm 0.02$  is obtained. Direct experimental evidence from low-energy

$\pi^\pm p$  elastic scattering gives<sup>5</sup>  $a_{1/2} - a_{3/2} = 0.30 \pm 0.02$ . The soft-pion calculation for  $\pi N$  scattering gives

$$\begin{aligned} a_{1/2} - a_{3/2} &= (45/f_\pi)^2 = 0.23, \quad \text{for } f_\pi = 94 \text{ MeV} \\ &= 0.30, \quad \text{for } f_\pi = 82 \text{ MeV}. \end{aligned}$$

A value  $f_\pi$  near 85 MeV is indicated.

For  $\pi\pi$  scattering a sum for the  $p$ -wave scattering length has been proposed<sup>6</sup> which can be quite accurately evaluated. This rule has a number of attractive features, including the small contribution of threshold  $s$  waves and asymptotic part relative to the intermediate region (primarily the  $\rho$  meson). Using the latest values of the  $\rho$  resonance parameters<sup>7</sup> ( $m_\rho = 770 \pm 4$  MeV,  $\Gamma_\rho = 111 \pm 6$  MeV) from  $e^+e^-$  colliding beams, we find that (in units of  $m_\pi^{-3}$ )

$$a_1 = 0.035 \pm 0.005.$$

The soft-pion calculation of  $a_1$  is free of ambiguity and is given by<sup>8,9</sup>

$$a_1 = (16/f_\pi)^2, \quad (2)$$

resulting in the value  $f_\pi = 86 \pm 6$  MeV, again closer to the GT value of 82 MeV than the  $\pi$ -decay value 94 MeV.

The soft-pion calculation of the mass-shell  $\pi-\pi$  amplitude is not unique. A new parameter<sup>9</sup>  $\xi$  measuring the deviation from the  $\sigma$  model must be given before the ratio

$$a_0/a_2 = (\frac{5}{2}\xi - 7)/(\xi + 2) \quad (3)$$

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<sup>1</sup> See, for example, Stephen L. Adler and Roger F. Dashen, *Current Algebras and Applications to Particle Physics* (W. A. Benjamin, Inc., New York, 1968).

<sup>2</sup> M. Olsson, Phys. Rev. **171**, 1681 (1968); all scattering lengths are expressed in pion mass units. All momenta are expressed in GeV/c.

<sup>3</sup> J. Hamilton and W. Woolcock, Rev. Mod. Phys. **35**, 737 (1963); V. Samaranayake and W. Woolcock, Phys. Rev. Letters **15**, 936 (1965).

<sup>4</sup> R. Seamon *et al.*, Phys. Rev. **165**, 1579 (1968); M. MacGregor, R. Arndt, and R. Wright, *ibid.* **169**, 1128 (1968).

<sup>5</sup> R. A. Donald, W. H. Evans, W. Hart, P. Mason, D. E. Plane, and E. J. C. Read, Proc. Phys. Soc. (London) **87**, 445 (1966).

<sup>6</sup> M. G. Olsson, Phys. Rev. **162**, 1338 (1967).

<sup>7</sup> J. E. Augustin *et al.*, Phys. Letters **288**, 508 (1969).

<sup>8</sup> S. Weinberg, Phys. Rev. Letters **17**, 616 (1966).

<sup>9</sup> M. G. Olsson and Leaf Turner, Phys. Rev. Letters **20**, 1127 (1968).

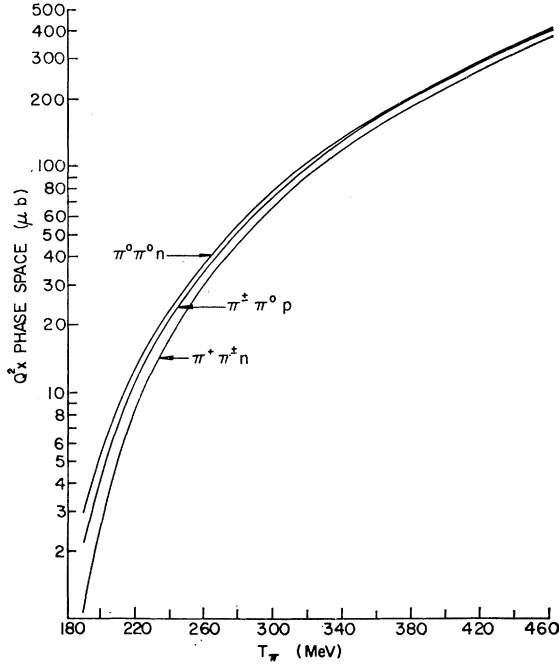


FIG. 1. The pion lab kinetic energy ( $T_\pi$ ) dependence of  $Q^2 \times \text{phase space}$  for  $\pi N \rightarrow \pi\pi N$ .

of  $s$ -wave scattering lengths can be calculated. It can be further shown<sup>9</sup> that the same parameter  $\xi$  enters into the threshold amplitude for the process  $\pi N \rightarrow \pi\pi N$ . One such reaction  $\pi^- p \rightarrow \pi^- \pi^+ n$  has been measured close to production threshold. The ratio of measured cross section to three-body phase space has the following dependence<sup>10</sup> on  $f_\pi$  and  $\xi$

$$\left[ \frac{\sigma(\pi^- \pi^+ n)}{Q^2 \times \text{phase space}} \right]^{1/2} = \left( \frac{82}{f_\pi} \right)^2 (1.36 - 0.6\xi) = 1.7 \pm 0.16. \quad (4)$$

The measurements used are at the two lowest energies where data are available.<sup>11</sup> The  $p$ -wave  $\pi$ - $\pi$  scattering length is related to the  $s$ -wave lengths in the soft-pion scheme<sup>8,9</sup> by

$$2a_0 - 5a_2 = 18a_1 = (68/f_\pi)^2. \quad (5)$$

Using (3)–(5), we find that

$$a_0 = (0.20 \pm 0.02) + (16/f_\pi)^2, \quad (6)$$

<sup>10</sup> Leaf Turner, Ph. D. thesis, University of Wisconsin, 1969 (unpublished). The notation used is described in detail later on in this article.

<sup>11</sup> Yu. A. Batusov, S. A. Bunyatov, V. M. Sidorov, and V. A. Yarba, *Yadern. Fiz.* **1**, 526 (1965) [English transl.: *Soviet J. Nucl. Phys.* **1**, 374 (1965)].

so that  $a_0$  is only slightly dependent on  $f_\pi$ . For the usual values of  $f_\pi$ ,  $a_0$  changes by only 0.01, giving

$$a_0 = 0.24 \pm 0.03. \quad (7)$$

The parameter  $\xi$  falls between zero and minus one for  $f_\pi$  ranging between 80 and 95 MeV. The isospin-two scattering length  $a_2$  is, however, more sensitive to the value of  $f_\pi$ ;  $a_2 = (0.08 \pm 0.01) - (29/f_\pi)^2$ , giving an  $a_2$  varying between  $-0.01$  and  $-0.04$ . Thus the amplitude for  $\pi^- p \rightarrow \pi^- \pi^+ n$  is nearly proportional to  $a_0$ . Hence,  $a_0$  is determined from the production cross section almost independently of  $f_\pi$  and  $\xi$ .

On the other hand, a measurement of  $\pi^+ p \rightarrow \pi^+ \pi^+ n$  near threshold would determine  $a_2$  directly, since both of these quantities depend on  $\xi$  and  $f_\pi$  in similar ways. A third production reaction  $\pi^- p \rightarrow \pi^0 \pi^0 n$  is useful, since it is sensitive to the relative sign of  $a_0$  and  $a_2$ .

From Ref. 9 we can explicitly calculate the above cross sections:

$$\begin{aligned} \sigma(\pi^- p \rightarrow \pi^- \pi^+ n) &= |a(-+n)|^2 Q^2 \times \text{phase space}, \\ \sigma(\pi^- p \rightarrow \pi^0 \pi^0 n) &= |a(00n)|^2 \frac{1}{2} Q^2 \times \text{phase space}, \\ \sigma(\pi^+ p \rightarrow \pi^+ \pi^+ n) &= |a(++n)|^2 \frac{1}{2} Q^2 \times \text{phase space}, \end{aligned}$$

where (with  $f_\pi = 82$  MeV)

$$\begin{aligned} a(-+n) &= -1.36 + 0.6\xi, \\ a(00n) &= 2.11 - 0.3\xi, \\ a(++n) &= 1.51 + 0.6\xi. \end{aligned}$$

The factor of  $\frac{1}{2}$  appearing in the later two cross sections takes into account the identity of the final pions. The quantity  $Q^2 \times \text{phase space}$  (where  $Q$  is the incident c.m. pion momentum) is shown in Fig. 1. Use of the above relations gives cross sections in  $\mu\text{b}$ . Other production cross sections near threshold are much smaller;

$$\sigma(\pi^+ p \rightarrow \pi^+ \pi^0 p) = \sigma(\pi^- p \rightarrow \pi^- \pi^0 p) = \frac{1}{4} \sigma(\pi^+ p \rightarrow \pi^+ \pi^+ n).$$

The above predictions for production cross sections are expected to be accurate in an energy range where all the final particles are in relative  $s$  waves.

We have shown that by use of pion production data the  $s$ -wave  $\pi\pi$  scattering lengths can be determined directly independent of the parameters  $\xi$  and  $f_\pi$ . A value of  $f_\pi$  consistent with all pion scattering processes is  $f_\pi \approx 85$  MeV, which is close to the GT value. This  $f_\pi$  is an "effective" value containing all of the uncertainties of the zero-mass extrapolation, and seems to be the one appropriate for use in soft-pion calculations.