Errata

Rescattering Model Applied to Y^* Production, C. P. SINGH AND B. K. AGARWAL [Phys. Rev. **173**, 1611 (1968)]. In Eq. (7), the normalization factor $\frac{1}{2}(2\pi i)^2/(2\pi)^4$ should be replaced by $|\mathbf{q}'|m_b/8\pi W$. In Eq. (8), the corresponding change would be to replace $1/4(2\pi)^4$ by a factor $|\mathbf{q}'|^2 m_b^2/64\pi^2 W^2$, where $|\mathbf{q}'|$ is the c.m. momentum of the intermediate s-channel particles. This change would decrease our values of differential cross section by a factor of about $\frac{2}{5}$ for the case of $Y_1^*(1385)$ production at a K^- momentum of 1.46 GeV/c. This may be compensated by adjusting the value of $g_{Kp\Lambda}^2/4\pi$, which can be anything from 4.8 to 16 [see R. H. Graham, S. Pakvasa, and K. Raman, Phys. Rev. **163**, 1774 (1967)]. Thus our results remain essentially the same.

Soft-Photon Theorem for Bremsstrahlung in a Potential Model, LEON HELLER [Phys. Rev. **174**, 1580 (1968)]. P. Signell has pointed out an error in Sec. III, where the formulas are written for particles with spin. When the functions $T_n(\nu,t,\Delta_i,\Delta_f)$ and the spin operators $A_n(\mathbf{L}_1,\mathbf{L}_2,\mathbf{L}_3)$ in Eq. (29) are expanded in powers of the photon's energy, one must make certain that the values of the arguments of T_n about which one expands are consistent with the values of the arguments of A_n, as prescribed in Eqs. (23)–(25) and Eq. (12). The choice of arguments which was made in Eq. (30), described after that equation, is not consistent. This can be seen directly from the fact that t is required by definition to equal L_1^2, but $(\mathbf{p}_2'-\mathbf{p}_2)^2$ differs from Q^2 by $2\mu\boldsymbol{\gamma}\cdot\mathbf{Q}/m_1$. To correct this it is sufficient to expand the T_n in Eq. (29) about $t=Q^2$, with all other variables as before. The only change in the appearance of Eq. (30) is the addition of $(2\mu\boldsymbol{\gamma}\cdot\mathbf{Q}/m_1)\partial T_n/\partial t$ to the brackets which at present contain derivatives of T_n with respect to ν and Δ. The functions T_n in Eqs. (32) and (33) should also be evaluated at $t=Q^2$, and Eq. (33) acquires an extra term

$$\frac{e_1}{m_1}\left(\frac{\mathbf{p}_1'}{\Delta}+\frac{\mathbf{p}_1}{\Delta'}\right)\frac{2\mu\boldsymbol{\gamma}\cdot\mathbf{Q}}{m_1}\sum_n \frac{\partial T_n}{\partial t}\left(\frac{\bar{q}^2}{2\mu},Q^2,0,0\right)A_n(\mathbf{Q},\mathbf{N},\mathbf{P}).$$

Another way to avoid this problem is to write the general T-matrix element, Eq. (23), in terms of *unit vectors*:

$$\langle \mathbf{k}_1', \mathbf{k}_2'; s'm' | T(e) | \mathbf{k}_1, \mathbf{k}_2; sm\rangle = \sum_n \bar{T}_n(\nu,t,\Delta_i,\Delta_f)\langle s'm'|A_n(\hat{L}_1,\hat{L}_2,\hat{L}_3)|sm\rangle,$$

as is customary in elastic scattering. [The relation between the $\bar{T}_n$ and the T_n can be gotten from Eq. (34).] These variables are independent, so $\bar{T}_n$ and A_n can be expanded independently.