

## Exchange-Degenerate $\Lambda$ -Trajectory Interpretation of $K^+p$ Backward Elastic Scattering Data\*

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A baryon Regge exchange model for  $K^+p$  backward elastic scattering is constructed in accordance with the requirements of duality and the absence of  $Y=+2$  baryons. The dominant exchange-degenerate amplitude is associated with the  $\Lambda_\alpha$  ( $1115, \frac{1}{2}^+$ ) and  $\Lambda_\gamma$  ( $1520, \frac{3}{2}^-$ ) particles. Recent data on  $d\sigma/du$  ( $K^+p$ ) from 2 to 7 GeV/c are successfully reproduced.

NEW data on  $K^+p$  elastic scattering over the 2–7-GeV/c momentum range exhibit a backward peak with a rapid energy dependence.<sup>1–6</sup> Since the  $K^+p$  elastic amplitude has no prominent direct-channel resonances, we expect Reggeized baryon exchange to account for the backward peak down to rather low momenta.<sup>7</sup> To describe the present data on  $d\sigma/du$  ( $K^+p$ ), we propose a simple model based on the following assumptions: (i)  $\Sigma$ -exchange contributions can be neglected relative to  $\Lambda$ -exchange contributions; (ii) only exchanges of the leading  $\Lambda$  trajectories are important; (iii) the Regge amplitudes are exchange-degenerate, in accord with the duality principle<sup>8</sup> and the absence of prominent  $K^+p$  direct channel resonances. By the above arguments the dominant amplitudes will be associated with the  $\Lambda_\alpha$  ( $1115, \frac{1}{2}^+$ ) and  $\Lambda_\gamma$  ( $1520, \frac{3}{2}^-$ ) particles, which lie on a degenerate trajectory and have the same residue function. Quantitative fits are presented which reproduce the  $K^+p$  data with few parameters. Further tests of the model are proposed.

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<sup>1</sup> D. Cline, C. Moore, and D. Reeder, Phys. Rev. Letters **19**, 675 (1967).

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<sup>4</sup> W. F. Baker, K. Berkelman, P. J. Carlson, G. P. Fisher, P. Fluery, D. Hartill, R. Kalback, A. Lundby, S. Mukkin, R. Nierhaus, K. P. Pretzl, and J. Worlds, in *Proceedings of the Fourteenth International Conference on High Energy Physics, Vienna, 1968* (CERN, Geneva, 1968).

<sup>5</sup> G. S. Abrams, L. Eisenstein, T. A. O'Halloran, W. Shufeldt, and J. Whitmore, Phys. Rev. Letters **21**, 1407 (1968).

<sup>6</sup> G. Bellettini, in *Proceedings of the Fourteenth International Conference on High Energy Physics, Vienna, 1968* (CERN, Geneva, 1968), p. 329.

<sup>7</sup> Meson Regge exchange models for  $K^+p$  elastic and charge-exchange scattering appear to fit the forward data rather well at low momenta ( $\sim 1$  BeV/c); A. A. Hirata *et al.*, University of California Radiation Laboratory Report No. UCRL-18322, 1968 (unpublished); R. W. Bland *et al.*, University of California Radiation Laboratory Report No. UCRL-18323, 1968 (unpublished).

<sup>8</sup> R. Dolen, D. Horn, and C. Schmid, Phys. Rev. **166**, 1768 (1968); G. F. Chew and A. Pignotti, Phys. Rev. Letters **20**, 1078 (1968).

Hyperon exchange models for the  $K^+p$  amplitude are complicated by the numerous possibilities for  $\Lambda$  and  $\Sigma$  exchanges. However, the following considerations suggest that the  $\Lambda$ -exchange contributions will dominate in  $K^+p$  backward scattering.

(i) Coupling constants from dispersion relations give<sup>9</sup>

$$g_{p\Lambda K^2} \gg g_{p\Sigma^0 K^2}.$$

(ii) Direct-channel resonance contributions to  $K^-p$  elastic scattering from  $I=0$  resonances are appreciably larger than from  $I=1$  resonances.

(iii) From the known residues of  $N_\alpha$  and  $\Delta_\delta$  Regge exchanges in  $\pi^\pm p$  backward scattering,<sup>10</sup>  $SU(3)$  estimates for the residues of  $\Lambda_\alpha$ ,  $\Sigma_\alpha$ , and  $\Sigma_\delta$  exchanges in  $K^+p$  backward scattering give  $\Lambda_\alpha$  amplitude  $\gg \Sigma_\alpha$  or  $\Sigma_\delta$  amplitudes.

(iv) The strong backward peak observed in data<sup>11</sup> on the reaction  $K^-p \rightarrow \pi^0 \Lambda_\gamma(1520, \frac{3}{2}^-)$  leads us to expect a significant contribution from  $\Lambda_\gamma$  exchange in  $K^+p \rightarrow K^+p$ .

Since the trajectory for  $\Lambda_\beta(1405, \frac{1}{2}^-)$  is expected to lie about a unit below the  $\Lambda_\alpha$  or  $\Lambda_\gamma$  trajectories,<sup>12</sup> the  $\Lambda_\beta$ -exchange contribution to  $K^+p$  scattering should die out extremely rapidly with energy. We therefore shall retain only  $\Lambda_\alpha$ - and  $\Lambda_\gamma$ -exchange contributions in subsequent calculations. These trajectories are illustrated in Fig. 1.

The Regge plot of the  $\Lambda_\alpha$  and  $\Lambda_\gamma$  particles<sup>13</sup> in Fig. 1 displays approximate exchange degeneracy for the trajectories. This degeneracy has been speculatively correlated with the nonexistence of  $K^+p$  direct-channel resonances through the duality principle.<sup>14–16</sup> In order that counterclockwise loops are not present in the Argand diagrams of the partial-wave decomposition

<sup>9</sup> J. Kim, Phys. Rev. Letters **19**, 1079 (1968).

<sup>10</sup> V. Barger and D. Cline, Phys. Rev. Letters **21**, 392 (1968).

<sup>11</sup> J. Badier *et al.*, Saclay Report, 1966 (unpublished).

<sup>12</sup> V. Barger and D. Cline, Phys. Rev. **155**, 1792 (1967).

<sup>13</sup> N. Barash-Schmidt *et al.*, University of California Radiation Laboratory Report No. UCRL-8030, 1968 (unpublished).

<sup>14</sup> C. Schmid, Phys. Rev. Letters **20**, 689 (1968); Proceedings of the 1968 Institute for Theoretical Physics at the University of Colorado (unpublished).

<sup>15</sup> G. Veneziano, CERN Report No. Th. 924, 1968 (unpublished).

<sup>16</sup> V. Barger and D. Cline, University of Wisconsin Report (unpublished).

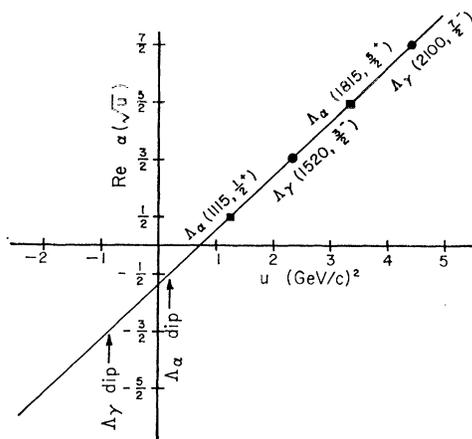


FIG. 1. Exchange degenerate ( $\Lambda_\alpha, \Lambda_\gamma$ ) Regge trajectory. The  $u$  positions of amplitude zeros at wrong-signature nonsense points are indicated by vertical arrows.

of the Regge exchange amplitude, the  $\Lambda_\alpha$  and  $\Lambda_\gamma$  contributions must be exchange degenerate (by duality, counterclockwise loops would be associated with the existence of  $K^+p$  direct-channel resonances<sup>14</sup>). With this exchange-degenerate requirement, the  $u$ -channel amplitude for  $K^+p$  backward scattering has the form<sup>17,18</sup>

$$\tilde{f}_2(\sqrt{u}, s) = \frac{1}{\Gamma(\alpha + \frac{1}{2})} \gamma \left( \frac{1 + e^{-i\pi(\alpha-1/2)}}{\cos\pi\alpha} + \frac{1 - e^{-i\pi(\alpha-1/2)}}{\cos\pi\alpha} \right) s^{\alpha-1/2},$$

where  $\alpha(\sqrt{u})$  and  $\gamma(\sqrt{u})$  denote the trajectory and residue, respectively. The first term above represents the  $\Lambda_\alpha$  exchange, the second  $\Lambda_\gamma$ . The  $\Lambda_\alpha$  amplitude vanishes at  $\alpha = -\frac{1}{2}$  ( $u \simeq +0.2$ ) and the  $\Lambda_\gamma$  amplitude vanishes at  $\alpha = -\frac{3}{2}$  ( $u \simeq -0.8$ ), as noted in Fig. 1. However, the sum of the two contributions is a smoothly varying function of  $u$ .

For data fitting we assume an approximately linear trajectory, as suggested by the Regge plot in Fig. 1. The residue is parametrized with a simple form consistent with the absence of  $J^P = \frac{1}{2}^-$  and  $J^P = \frac{3}{2}^-$  MacDowell reflected particles,<sup>19</sup>

$$\gamma(\sqrt{u}) = \left( 1 + \frac{\sqrt{u}}{M_\Lambda} \right) \left( 1 + \frac{\sqrt{u}}{M_{Y_0^*}} \right) \beta(u),$$

<sup>17</sup> The cross-section formulas for fermion Regge exchange are given in Ref. 12. The amplitude  $f_2$  of Ref. 12 is related to  $\tilde{f}_2$  by  $f_2(\sqrt{u}, s) = \tilde{f}_2(\sqrt{u}, s)(E_u + M)/\sqrt{u}$ . The units of  $s$ ,  $t$ , and  $u$  are  $(\text{GeV}/c)^2$ .

<sup>18</sup> Since we retain only the leading power in the Regge expansion, some question exists as to whether  $s^\alpha$  or  $(s-t)^\alpha$  represents the better approximation at low energies. We have tried both forms, with similar results for  $K^+p$ . Fits to the 3-7-GeV/c  $K^+p$  data were somewhat better with the  $(s-t)^\alpha$  form, but fits to the more accurate 2-3-GeV/c data were better with the  $s^\alpha$  form. Baryon-exchange calculations of  $\bar{p}p \rightarrow K^-K^+$  at low energy are extremely sensitive to the choice of an energy-dependent factor.

<sup>19</sup> V. Barger and D. Cline, Phys. Letters **26B**, 83 (1967). The observed  $\Lambda_B(1830, \frac{5}{2}^-)$  particle is presumably the MacDowell reflected state associated with  $\Lambda_\alpha(1815, \frac{5}{2}^+)$ .

where  $M_\Lambda \simeq 1.115$  GeV and  $M_{Y_0^*} \simeq 1.52$  GeV. For this residue structure, the  $u$ -channel amplitude is

$$\tilde{f}_2(\sqrt{u}, s) = \frac{2\beta(u)}{\Gamma(\alpha + \frac{1}{2}) \cos\pi\alpha} \left( 1 + \frac{\sqrt{u}}{M_\Lambda} \right) \left( 1 + \frac{\sqrt{u}}{M_{Y_0^*}} \right) s^{\alpha-1/2}.$$

The fits to the 2-7-GeV/c data<sup>1-6</sup> are shown in Figs. 2 and 3. The  $(\Lambda_\alpha, \Lambda_\gamma)$  trajectory was determined to be

$$\alpha = -0.72 + 1.01u$$

as compared with

$$\alpha = -0.7 + 0.96u$$

from the Regge spin-mass plot of particles in Fig. 1. The curves in Figs. 2 and 3 correspond to the residue parametrization (in  $\text{GeV}^{-1}$ )<sup>20</sup>

$$\beta = 9.1e^{2.1u} + 3.1e^{-1.5u}.$$

The second term in  $\beta$  was needed to account for the tendency of the  $d\sigma/du$  data to flatten at larger  $|u|$ . The model seems to provide a reasonable over-all description of the present  $K^+p$  data.

The measurement of  $K^+n \rightarrow K^0p$ ,  $K_2p \rightarrow K_1p$ , or  $K^+n \rightarrow K^+n$  backward scattering will provide a straightforward experimental test of the  $\Lambda$ -exchange dominance assumption of this model. We expect the

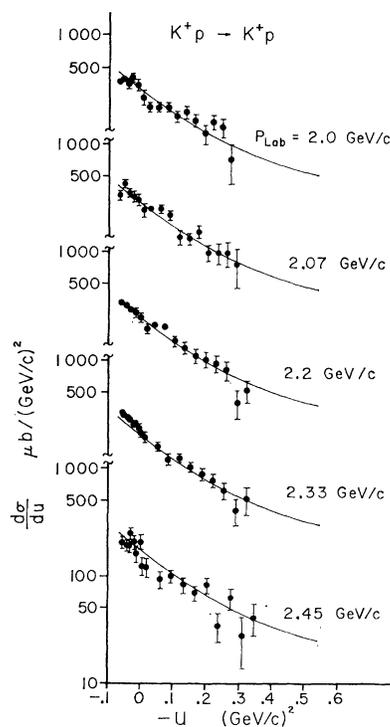


FIG. 2. Fitted curves for  $d\sigma/du$  ( $K^+p$ ) based on  $(\Lambda_\alpha, \Lambda_\gamma)$  exchange-degenerate Regge model. Data from Ref. 3.

<sup>20</sup> For an  $(s-t)^\alpha$  energy dependence (see Ref. 18), the fitted residue was  $\beta = 7.8e^{1.9u} + 1.2e^{-1.7u}$ . The parametric form for the residue is not well suited for extrapolation outside the  $u$  range considered in the data analysis.

following relations to hold:

$$\frac{d\sigma}{du}(K^+n \rightarrow K^0p) \approx \frac{d\sigma}{du}(K^+p \rightarrow K^+p),$$

$$\frac{d\sigma}{du}(K_2p \rightarrow K_1p) \ll \frac{d\sigma}{du}(K^+p \rightarrow K^+p),$$

$$\frac{d\sigma}{du}(K^+n \rightarrow K^+n) \ll \frac{d\sigma}{du}(K^+p \rightarrow K^+p).$$

The assumption of exchange degeneracy is somewhat more difficult to evaluate experimentally. The polarization is predicted to be zero in the exchange-degenerate limit, but polarization phenomena are typically sensitive to small background contributions. Finite-energy sum rules (FESR) for the  $u$  channel provide an alternative test of the exchange degeneracy, since the  $\Lambda_\alpha$  and  $\Lambda_\gamma$  amplitudes can be separated through their opposite signature.<sup>21</sup> However, lack of information on the  $\bar{p}p \rightarrow \bar{K}K$  channel will cause uncertainties in such FESR evaluations.

The absence of prominent resonances in the  $K^+p$  elastic channel makes it a particularly interesting reaction for phenomenological study. Our analysis indicates that  $\Lambda$  exchanges can account for the backward peak at intermediate energies. It will be of interest to see to what extent this model can account for  $K^+p$  backward scattering data below 2 GeV/c. At the same time, experiments above 7 GeV/c are highly desirable to determine if the measured energy dependence for

<sup>21</sup> D. Beder and J. Finkelstein, Phys. Rev. **160**, 1363 (1967); C. Chiu and M. Der Sarkissian, Nuovo Cimento **55A**, 396 (1968).

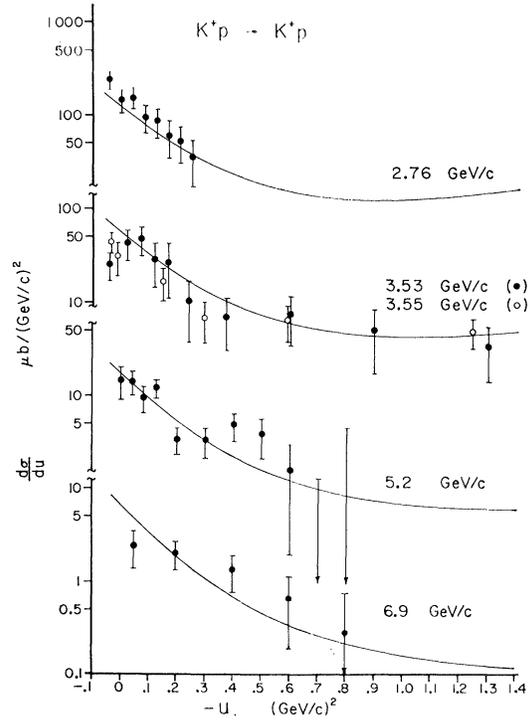


FIG. 3. Fitted curves for  $d\sigma/du(K^+p)$  based on  $(\Lambda_\alpha, \Lambda_\gamma)$  exchange-degenerate Regge model. Data from Refs. 1 (3.53 GeV/c), 2 (3.55 GeV/c), 4 (5.2 and 7.0 GeV/c), and 5 (2.76 GeV/c).

$d\sigma/du(K^+p)$  remains compatible with the predictions from  $(\Lambda_\alpha, \Lambda_\gamma)$  exchanges.

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