Alphas Branching Ratio in the Decay of ¹⁴⁹Tb[†]

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4.1-h ¹⁴⁹Tb has been found to decay by α emission in (22.6±2.3)% of its disintegrations. This was determined by measurements of the α activity in a chemically and isotopically separated ¹⁴⁹Tb sample and by determination of the amount of ¹⁴⁹Gd formed in this sample through electron-capture decay. Branching ratios of 0.475 ± 0.055 and 0.104 ± 0.013 were deduced for the 894- and 654-keV γ rays of ¹⁴⁶Eu formed in the ¹⁴⁹Tb α decay.

LPHA-PARTICLE emission in neutron-deficient A rare-earth nuclides was discovered by Thompson et al. in 1949,¹ and among the first of these α emitters reported was a nuclide of 4-h half-life emitting 4.0-MeV α particles, later identified² as ¹⁴⁹Tb. Among the rareearth α emitters with appreciable α branches (>1%) this turns out to be the only one with a half-life in the range between 18 min (151Dy) and 84 years (148Gd). For this reason, the α particles emitted by ¹⁴⁹Tb can be readily detected even in the presence of complex mixtures of other rare-earth nuclides and from thick samples. The formation of ¹⁴⁹Tb in foils of high-Z materials, particularly gold, has therefore often served for the monitoring of high-energy proton beams.^{3,4} To put such measurements on an absolute basis, cross sections for production of the α branch of ¹⁴⁹Tb from gold by protons of various energies have been carefully determined.³⁻⁶ However, the fraction of ¹⁴⁹Tb decays proceeding by α emission is still not accurately known; the literature values⁷⁻¹⁰ for this fraction range from 0.1 to 0.2, with large uncertainties for the individual values. We have therefore redetermined the α -branching ratio of 4.1-h 149Tb.

The ¹⁴⁹Tb sample was produced from high-energy spallation of tantalum. The tantalum target was irradiated in the circulating beam of the alternating gradient synchrotron at Brookhaven with 28-GeV protons for 30 min. After irradiation the target was dissolved and standard chemical separations were performed to separate the rare-earth group from the tantalum and from other radioactivities produced by the high-energy protons. The gross rare-earth fraction was finally pre-

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cipitated as oxalate. The sample was dried and ignited to oxides in a furnace at 900°C. The oxides were introduced into the ion source of the Brookhaven electromagnetic isotope separator, where they were converted into chlorides by passing carbon tetrachloride vapor over them. The volatile chlorides were then ionized. The mass-separated samples were collected on a long, thin aluminum foil and a strip containing mass 149 with $\ll 1\%$ contamination by neighboring masses was used for the subsequent measurement. The α -emission rate of this sample was measured repeatedly with a calibrated surface-barrier silicon detector connected to a multichannel pulse-height analyzer. The α -particle energy (3.95 MeV) and half-life (4.1 h) found were those characteristic of ¹⁴⁹Tb. After these α -particle measurements, the mass-149 strip was dissolved in acid solution with Tb, Gd, and Eu carriers. The sample was loaded on a Dowex-50 cation exchange column. The eluent used was $0.4M \alpha$ -hydroxyisobutyric acid solution adjusted to pH 4.1 with ammonium hydroxide. The column was operated at 80°C. The terbium fraction came off the column about 50 min after the beginning of the elution. For the further measurements, a narrow cut of the terbium elution peak was used so as to minimize the amount of 149Gd present in the 149Tb sample. The terbium was precipitated as oxalate. Later yield determinations showed the cross contamination of gadolinium in this terbium sample to be less than 0.1%, whereas the terbium recovery was about 60%.

In the chemically separated Tb sample, the 4.1-h ¹⁴⁹Tb was allowed to decay and then the γ radiations of 9.25-day ¹⁴⁹Gd and 5.9-day ¹⁴⁵Eu, grown in as a result of the electron capture and α decays of ¹⁴⁹Tb, respectively, were determined with a calibrated Ge(Li) detector. From the α -disintegration rate of ¹⁴⁹Tb measured in the gross A = 149 sample, from the chemical yield of Tb in the chemically separated sample, and from the amount of ¹⁴⁹Gd formed in this sample by electron-capture decay, the α -branching ratio of ¹⁴⁹Tb was calculated. Since the abundances of ¹⁴⁵Eu γ rays determined by different authors¹¹⁻¹⁴ vary widely-from

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0.41 to 0.68 per disintegration for the most abundant γ ray at 894 keV-the present measurements were used to deduce, from the ¹⁴⁹Tb α -disintegration rate and from the intensities of the 894- and 654-keV γ rays of ¹⁴⁵Eu in the ¹⁴⁹Tb sample, the branching ratios of these two γ rays.

The counting geometry of the surface-barrier Si detector used for the α measurements was calibrated with an ²⁴¹Am standard. The linewidth of the α peaks was 35-40 keV full width at half-maximum (FWHM). The resolution obtained with the 149Tb sample was comparable to that of the ²⁴¹Am standard. The γ radiations of both ¹⁴⁹Gd and ¹⁴⁵Eu were measured with a 9-cm² \times 1-cm Ge(Li) detector with a resolution of 3.5 keV FWHM at 661 keV.

The amount of ¹⁴⁹Gd formed in the ¹⁴⁹Tb sample was determined through measurements of the most abundant γ ray emitted in ¹⁴⁹Gd decay, at 150 keV. This γ ray was found to decay with a half-life of 9.25 ± 0.10 days as compared with literature values^{15,16} ranging from 9.0 to 10.3 days for the half-life of ¹⁴⁹Gd. The abundance of the 150-keV γ ray was taken from Ref. 15 to be 0.48 per disintegration. The same value results from the data in two recent papers^{17,18} under the plausible assumption that no electron-capture decays proceed directly to the ground state of ¹⁴⁹Eu. An independent check of this branching ratio was made as follows: The intensities of the 150-keV γ ray were

measured in two samples of ¹⁴⁹Gd which were then allowed to decay into ¹⁴⁹Eu. The disintegration rates of the resulting ¹⁴⁹Eu were determined through measurement of its K x-ray emission with a calibrated, 2-mmthick NaI(Tl) detector and single-channel analyzer. These duplicate determinations both yielded values of 0.47 for the branching ratio of the 150-keV γ ray of 149Gd.

Using the above-mentioned half-lives of ¹⁴⁹Tb and ¹⁴⁹Gd, the branching ratio of the 150-keV γ ray, and the measured disintegration rates, the α -branching ratio of 149 Tb was found to be 0.226.

In estimating the error of this branching ratio, there are several sources of uncertainties. Since the amount of decay of ¹⁴⁹Tb into ¹⁴⁹Gd during the chemical separation on the cation column was not negligible, the time at which the terbium fraction was, in fact, freed of gadolinium is somewhat uncertain. This factor was estimated to contribute an uncertainty of $\pm 4\%$ to the value of the branching ratio. Other uncertainties were estimated as follows: abundance of the 150-keV γ ray of ¹⁴⁹Gd—5%; efficiency calibration of γ detector—7%; efficiency calibration of α detector—2%; chemical yield of ¹⁴⁹Tb—2%; counting statistics—1%. The root-meansquare combination of these uncertainties is about 10%. Thus we report a value of 0.226 ± 0.023 for the α branching of 4.1-h ¹⁴⁹Tb.

From this α -branching ratio of ¹⁴⁹Tb, from the α decay rate of the ¹⁴⁹Tb sample, and from the amount of 894- and 654-keV γ radiation of ¹⁴⁵Eu grown in, the abundances of these two γ rays were found to be 0.475 ± 0.055 and 0.104 ± 0.013 per disintegration, respectively.

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