3. Page 1648, line 10 should read:

$$\frac{f_A(Y_1^*)^2}{M_{Y_0^*}+M_N}m_K\left(1-\frac{m_K}{M_N}\right)$$
$$-\frac{2}{3}\frac{g_A(Y_1^*)^2}{M_{Y_1^*}-M_N}\left(1+\frac{m_K}{2M_N}+\frac{m_K M_N}{M_{Y_1^*}{}^2}+\frac{m_K}{6M_{Y_1^*}}\right)\Big].$$

$K\pi$ Scattering, κ Meson, and Current Algebra, Probir Roy [Phys. Rev. **168**, 1708 (1968)]. There are some incorrect, although inconsequential, remarks in Sec. V of this paper. Although current algebra predicts that $(\sigma_{K^+\pi^-}{}^{\text{el}})_{\text{thr}}=(\sigma_{K^+\pi^+}{}^{\text{el}})_{\text{thr}}$, the process $K^+\pi^-\to K^0\pi^0$ contributes to $\sigma_{K^+\pi^-}{}^{\text{tot}}$. Thus it follows that $\Delta_{K\pi}$ does not vanish at threshold but is equal to

$$\frac{1}{2\pi}\left(\frac{m_K m_\pi}{m_K+m_\pi}\right)^2\frac{1}{F_\pi^4},$$

neglecting σ terms which are of higher order in the pion mass. We can now estimate the nonresonant S-wave contribution to the right-hand side of Eq. (1) if we assume that $\Delta_{K\pi}(s)$ equals the constant value at threshold and cut the integral off at $s=m_{K^*}{}^2$. This turns out to be $\simeq 0.03/m_\pi^2$, which is so much smaller than our discrepancy of $(0.220\pm0.036)/m_\pi^2$ as to be negligible.

The remarks following Eq. (6) in Sec. VI should now be modified somewhat. Unitarity implies that

$$\lim_{k\to 0}\operatorname{Im}M_{K^+\pi^\pm}{}^f=\lim_{k\to 0}\sum\frac{32\pi k}{\sqrt{s}}\left|\frac{M_{K^+\pi^\pm}}{16\pi}\right|^2, \qquad (6)$$

where the summation is over all outgoing channels. From Eq. (5), we see that near the physical threshold $\operatorname{Im}M_f$ is linear in k to the lowest order. Using the current-algebra[18] predictions for

$$m(K^+\pi^-\to K^+\pi^-)$$

and

$$M(K^+\pi^-\to K^0\pi^0),$$

we can now write

$$\operatorname{Im}M_{K^+\pi^-}{}^f=\frac{3k}{2\pi}\frac{m_\pi^2 m_K^2}{F_\pi^4}\frac{1}{m_\pi+m_K}+O(k^2).$$

Our estimate of the contribution of the interference term is now

$$\sigma_{K^+\pi^-}{}^{\text{intf}}(m_\kappa^2)=\frac{3}{2\pi}\frac{1}{m_\kappa(m_K+m_\pi)}\frac{m_K^2 m_\pi^2}{F_\pi^4},$$

so that the contribution to the right-hand side of Eq. (1) is still only $\sim 2\%$ of the contribution from the κ.

The following misprints should also be corrected. (1) Equation (2), page 1709 should read:

$$\Delta_{K\pi}{}^\rho\sim\frac{2\gamma_{\rho K}\gamma_{\rho\pi}}{q_{K\pi}s^{1/2}}(\sqrt{\pi})\frac{\Gamma(2)}{\Gamma(3/2)}\left(\frac{s-m_\pi^2-m_K^2}{s_1}\right)^{1/2}.$$

(2) The last two sentences on page 1709 should read: "If we neglect $O((m_\pi^2+m_K^2)/s)$ above s_0, Eq. (2) gives $\Delta_{K\pi}\sim 8.28/\sqrt{(ss_1)}$. Then the Regge contribution to the second integral is $\simeq 0.001/m_\pi^2$, which $\cdots$."

Symmetric Quark Model of Baryon Resonances, O. W. Greenberg and M. Resnikoff [Phys. Rev. **163**, 1844 (1967)]. The numerical value of the matrix element $\langle\Lambda(\mathbf{8},2)|T_{L\cdot S}{}^8|\Lambda(\mathbf{8},2)\rangle=-\frac{2}{3}$ was omitted from Table X.

New Formulation of the Axially Symmetric Gravitational Field Problem. II, Frederick J. Ernst [Phys. Rev. **168**, 1415 (1968)]. It was intended that the inequality $0\leqslant y\leqslant 1$ above Eq. (20) should read $-1\leqslant y\leqslant 1$. However, all that one can actually assert is that a scale of lengths may be chosen so that $-1+b\leqslant y\leqslant 1+b$, where b is some constant. Identifying y with $\cos\theta+b$, we are then led to the generalization of the Kerr and Newman metrics discovered by B. Carter [Bull. Am. Phys. Soc. **13**, 571 (1968)].