Optical Properties of Platinum*

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The optical functions $\epsilon_1, \epsilon_2, \alpha$, and $-\text{Im}(1/\epsilon)$ have been deduced from reflectance data on electron-gunevaporated Pt samples. They were obtained by two different methods: (1) the Kramers-Kronig analysis of the normal-incidence reflectance data, and (2) the analysis of reflectance at two different incident angles. The absorption coefficient α shows small peaks at $h\nu = 0.4$, 7.7, and 10.5 eV, and a broad peak centered about 4 eV. The imaginary part of the dielectric function, ϵ_2 , has structure at $h\nu = 6.5$ and about 11 eV.

I. INTRODUCTION

O complete the experimental studies of the electronic structure of transition metals at the end of the 3d (Ni), 4d (Pd), and 5d (Pt) series,^{1,2} we have undertaken the optical and photoemission studies of platinum. These metals are particularly interesting because Ni is ferromagnetic while Pd and Pt are not, although they all have the same number of outer electrons. In this paper, optical functions deduced from reflectance data will be presented and discussed. These data will be used along with photoemission results in a subsequent paper³ to deduce the electronic structure of Pt over an energy range about 10 eV above and below the Fermi level.

II. DETERMINATION OF THE OPTICAL FUNCTIONS OF PLATINUM

Recently, normal incidence reflectivity data of electron-gun evaporated Pt films have been reported⁴ over the spectral range 0.06 eV $\leq h\nu \leq 5.63$ eV. In Table I, data taken at several different angles for $6.88 \le h\nu$ \leq 21.23 eV are presented. The films were prepared in the same manner as reported previously.⁴ In Fig. 1 the normal-incidence reflectance curve is shown. This curve was obtained using the data of Ref. 4 and the 10° data from Table I and extrapolating from 5.63 to 6.88 eV. The reflectance has slope changes at 0.6 and 6.7 eV and a small peak is present at 10.1 eV.

Two approaches were used in calculating the optical functions. They were first calculated from the reflectance data at two different incident angles (10°, 60°, 60°, 80°, and 10°, 80°) at the available 14 data points in the region 6.88 eV $\leq h\nu \leq$ 21.23 eV by solving the two

exactly.⁵ Second, the optical functions over the complete range were calculated by the Kramers-Kronig (KK) analysis⁶ using the "normal-incidence" (10°) data. The high photon energy reflectance extrapolation was adjusted so that the optical functions obtained from the KK analysis agree with those from reflectance at two angles of incidence. In Fig. 2 ϵ_1 and ϵ_2 from KK analysis are shown along with results from the two-angle incidence data. This is the best match obtained for the two results. The agreement of the two ϵ_2 's is within 15% but it is worse (within about 50% at $h\nu = 16 \text{ eV}$) for the ϵ_1 's. The extrapolation has been adjusted so that the ϵ_2 values by the two methods agree most closely over the range 7 eV $\leq h\nu \leq 11$ eV, which is the region of the most interest. The curves by the two methods have different over-all slopes which could not be made equal by adjusting the high photon energy extrapolation. This discrepancy is probably due to the following reasons: (1) The light used for reflectivity measurements probably was not completely unpolarized; (2) the data points for $h\nu > 12$ eV were far apart (~3 eV). The reflectance be-

equations relating the optical functions with reflectivity



⁶ R. Tousey, J. Opt. Soc. Am. 29, 235 (1939); A. Y-C. Yu, Ph.D. thesis, Stanford University, 1967 (unpublished). ⁶ J. L. Shay (private communication).

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¹A. J. Blougett, Jr., and W. E. Spicer, Phys. Rev. (to be published).
² A. Y-C. Yu and W. E. Spicer, Phys. Rev. (to be published).
³ A. Y-C. Yu and W. E. Spicer, Phys. Rev. (to be published).
⁴ L. F. Durmmeter and G. Hass, *Physics of Thin Films* (Academic Press Inc., New York, 1964), Vol. 2, p. 337.

tween data points was extrapolated by exponentials. The error introduced by this will be appreciable since data points are relatively far apart. However, since for 7 eV $\leq h\nu \leq 11$ eV the data points were quite close together (~0.5 eV) and since the agreement between results of the two different methods is good, the optical functions in this region should be very reliable. Note that location of the major pieces of structure in ϵ_2 is the same for the two ϵ_2 curves.

TABLE I. Reflectance data of Pt (G. Hass).

<i>hv</i> (eV)	Angle of Incidence (deg)	Reflectance (%)
21.23	80 60 10	61.17655 31.13701 21.72616
16.84	80 60 10	55.77633 24.39885 14.69829
14.86	80 60 10	57.02338 26.25589 15.66564
13.68	80 60	54.39103 24.30080 14.04017
12.08	10 80 60	55.85290 25.81948
10.92	10 80 60	16.48173 56.58180 27.93711
10.19	10 80 60	19.20076 55.60826 28.26232
9.71	10 80 60	20.66767 54.63517 27.69484
9.23	10 80 60	19.90159 54.70237 27.37519
8.83	10 80 60	19.87917 54.32747 27.63155
8.33	10 80 60	20.25131 54.79938 28.18447
7.7	10 80 60	20.81851 55.84696 29.76346
7.29	10 80 60	22.51346 55.66986 29.88077
6.88	10 80 60	23.33367 56.53064 30.76324
	10	23.85045

In Fig. 3 the absorption coefficient α and the loss function $-\text{Im}(1/\epsilon)$ obtained from KK analysis are shown. The agreement between the absorption coefficients obtained from the two different methods is very good. α shows a small piece of structure at 0.4 eV, a very broad maximum centered around 4 eV, and a shoulder at 10.5 eV. The origin of this structure will be discussed in a subsequent paper.³ The loss function has



a small peak at 6.5 eV. Note that the electron energy loss spectra⁷ show a small peak at 6.2 eV, in reasonable agreement with the optical results. However, this peak is probably not due to an "ideal" plasma resonance because ϵ_2 is quite large (~3) there and damping will be severe.

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⁷ C. J. Powell, Proc. Phys. Soc. (London) 76, 583 (1960).