

The \mathbf{J}_μ and $\mathbf{J}_{5\mu}$ in Appendix D should be changed to

$$\begin{aligned} \mathbf{J}_\mu &= -\xi\gamma_\mu\frac{1}{2}\tau\xi + \rho^2\pi \times \partial_\mu\pi + f\bar{\xi}\gamma_\mu\gamma_5\tau \times \pi\xi\rho \\ &\quad + 2f^2\bar{\xi}\gamma_\mu\left(\frac{\pi^2\tau - (\tau \cdot \pi)\pi}{1+\sigma}\right)\xi\rho^2, \\ \mathbf{J}_{5\mu} &= -\bar{\xi}\gamma_\mu\gamma_5\frac{1}{2}\tau\xi + \frac{\rho\sigma}{2f}\partial_\mu\pi \\ &\quad + f\{\pi(\partial_\mu\pi^2)\frac{1}{2}(\sigma\rho' - \rho\sigma') + \bar{\xi}\gamma_\mu\tau \times \pi\xi\rho\} \\ &\quad + 2f^2\bar{\xi}\gamma_\mu\gamma_5\left(\frac{\pi^2\tau - (\tau \cdot \pi)\pi}{1+\sigma}\right)\rho^2\xi. \end{aligned}$$

The axial currents $\mathbf{J}_{5\mu}$ connected with the three specific models that appear in Eqs. (4.1a), (4.3a), and (4.5a) should also be corrected accordingly.

The last sentence of Sec. 4 should read: "This accounts for the absence of the higher-order part of the π - N cross term in the expression of the axial-vector current that appears in Schwinger's model."⁶

We are indebted to Dr. V. Ogievetsky for pointing out an error that led to incomplete expressions for the vector and axial-vector currents.

Operator Formalism for Daughter Trajectories in the Bethe-Salpeter Equation, VICTOR CHUNG AND JON WRIGHT [Phys. Rev. **162**, 1716 (1967)]. The sentence after Eq. (4.8) is in error. For unequal masses, there is a term

$$\lim_{\sqrt{s} \rightarrow 0} \frac{\langle \chi(\sqrt{s}) | \cos\beta | \chi(\sqrt{s}) \rangle}{\sqrt{s}}$$

which can be shown to lead to the matrix elements

$$\langle D_\kappa^{\alpha-\kappa+1} | \cos\beta | D_{\kappa-1}^{\alpha-\kappa+1} \rangle$$

and

$$\langle D_\kappa^{\alpha-\kappa+1} | \cos\beta | D_{\kappa+1}^{\alpha-\kappa+1} \rangle.$$

The slope formula Eq. (4.12) remains unchanged.

We wish to thank Professor N. Nakanishi for bringing this error to our attention.

Perturbational-Variational Approach to the Calculation of Variational Wave Functions. I. Theory, JEREMIAH N. SILVERMAN AND JON C. VAN LEUVEN [Phys. Rev. **162**, 1175 (1967)]. p. 1180, right column, line 3: delete "quantized." p. 1180, right column, line 30 and last line: read " $\bar{W}_{m,n}$ " instead of " $W_{m,n}$." p. 1181, left column, line 1, and right column, line 14: read " $\bar{W}_{m,m}$ " instead of " $W_{m,m}$." p. 1182, Table I, cycle $n+1$: read " $\langle A \rangle_n$ " instead of " $\langle A \rangle_{n-1}$." p. 1189, Eq. (B4) should read

$$\prod_{j=1}^{2n} \frac{d_j^{k_j}}{k_j!} = d_s \prod_{j=1}^{2n-s} \frac{d_j^{k_j}}{k_j!}, \quad n+1 \leq s \leq 2n.$$

Pole Approximations in the N/D Equations, MITCHEL J. SWEIG [Phys. Rev. **165**, 1893 (1968)]. It has been called to the author's attention that the use of Padé approximants in the N/D equations was first suggested by M. Bander [J. Math. Phys. **5**, 1427 (1964)]. There is also a recent paper by A. K. Common [J. Math. Phys. **8**, 1669 (1967)] which contains many of the ideas of the present work, with a somewhat more detailed discussion of the mathematics involved.

High-Spin Baryons. I. The Pion-Baryon-Baryon Vertex, P. CARRUTHERS [Phys. Rev. **152**, 1345 (1966)]. The following sign errors in Sec. IV should be corrected. Equation (4.4) should be replaced by

$$T^{(\rho)}_{\nu\mu} = (-1)^\rho (2T'+1)^{1/2} C(T'1T; \nu-\rho\mu).$$

The left-hand side of Eq. (4.5) should be

$$g \sum_\rho (-1)^\rho (\pi^\rho)^* \psi_{T'}^\dagger T^{(-\rho)} \psi_T + \text{H.c.},$$

and the left-hand side of Eq. (4.10) is simply $T^{(\rho)}_{\mu\nu}$. Finally, the over-all sign of T_ν in Eq. (4.11) should be changed.