

S-Wave $\pi^- + p$ Scattering and a Possible η - n Bound State

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The $\pi^- + p$ $T = \frac{1}{2}$, $J = \frac{1}{2}^-$ scattering length and phase shift are calculated in a coupled-channel model of π - N and η - N scattering. The η - n system is treated as a Regge pole associated with quasibound state. The Born terms associated with ρ and nucleon exchange are included in the calculation. The results are compared with experimental data.

AN extension of Regge¹ representation for studying the low-energy behavior of scattering amplitudes has been applied to $T = \frac{1}{2}$, $J = \frac{1}{2}^-$ state of π - N scattering. The Khuri representation² was modified by Bohm and Rashid³ for $T = \frac{1}{2}$, $J = \frac{1}{2}^-$ π - N scattering, and by Pal⁴ for \bar{K} - N scattering, by introducing Born terms where possible. Bulos *et al.*⁵ and Peterson *et al.*⁶ found a rapid rise in the inelasticity in $\pi^- p$ scattering associated with the threshold for η production in the reaction $\pi^- + p \rightarrow \eta + n$. The elegant phase-shift analysis of Auvil *et al.*⁷ suggests that a discrete baryon state may exist in the neighborhood of $\eta + n$ threshold. Further evidence for such a state is obtained from the recent experimental data,⁸ which show that the total cross section for the reaction $\pi^- + p \rightarrow \eta + n$ rises steeply from threshold to a value of 1 mb at an incident pion kinetic energy (T_{π^-}) between 655 and 704 MeV, and then falls gradually to 0.25 mb at 1300 MeV. The quantum numbers of this baryon state are of course $T = \frac{1}{2}$, $J = \frac{1}{2}^-$.

In the present paper, we use the modified Khuri representation for π - N scattering and the foregoing data on the η - n system to calculate the $T = \frac{1}{2}$, $J = \frac{1}{2}^-$ scattering length and phase shift. The η - n system will be treated as a Regge pole associated with a quasibound state in all the channels $\pi^- + p \rightarrow \pi^- + p$, $\eta + n \rightarrow \eta + n$ and $\pi^- + p \rightarrow \eta + n$ in S -wave $T = \frac{1}{2}$ scattering.

The $S_{1/2}$ partial wave for these reactions can be written as a matrix in the following way:

$$T = \begin{pmatrix} T_{11} & T_{12} \\ T_{21} & T_{22} \end{pmatrix}. \quad (1)$$

Using the Khuri representation, the contribution of the η - n Regge pole to the elements of the T matrix can be calculated in a straightforward way. The result for T , including the Born terms associated with nucleon and ρ -meson exchange in $\pi^- + p \rightarrow \pi^- + p$ elastic scattering, is

$$T_{11} = \frac{\beta}{\alpha(W) - \frac{1}{2}} [e^{(\alpha-1/2)\xi_1} + e^{(\alpha-1/2)\xi_2}] - \frac{g_{NN\pi}^2}{4\pi W^2} \left\{ [(W+m)^2 - m_\pi^2] \left(\frac{W-m}{2k^2} \right) Q_0 \left(\frac{2E^2 + 2m_\pi^2 - W^2 - m^2}{2k^2} \right) \right. \\ \left. + [(W-m)^2 - m_\pi^2] \left(\frac{W+m}{2k^2} \right) Q_1 \left(\frac{2E^2 + 2m_\pi^2 - W^2 - m^2}{2k^2} \right) \right\} - \frac{2g_{\rho\pi\pi}g_{\rho NN}}{4\pi W^2} \left\{ [(W+m)^2 - m_\pi^2] \right. \\ \left. \times \left(\frac{W-m}{k^2} \right) Q_0 \left(\frac{2E^2 - 2m^2 + m_\rho^2}{2k^2} \right) + [(W-m)^2 - m_\pi^2] \left(\frac{W+m}{k^2} \right) Q_1 \left(\frac{2E^2 - 2m^2 + m_\rho^2}{2k^2} \right) \right\}, \quad (2)$$

$$T_{12} = \frac{\beta_{12}}{\alpha(W) - \frac{1}{2}} [e^{(\alpha-1/2)\xi_{1'}} + e^{(\alpha-1/2)\xi_{2'}}], \quad (3)$$

$$T_{22} = \frac{\beta_2}{\alpha(W) - \frac{1}{2}} [e^{(\alpha-1/2)\xi_{1'}} + e^{(\alpha-1/2)\xi_{2'}}]. \quad (4)$$

The coupling constants $g_{NN\pi}$, $g_{\rho\pi\pi}$, and $g_{\rho NN}$ are taken from experiment. In (2)-(4), $\alpha(W)$ is the trajectory function, and the β 's are the corresponding residue

¹ N. N. Khuri, Phys. Rev. **130**, 429 (1963).

² N. N. Khuri and B. M. Udgankar, Phys. Rev. Letters **10**, 72 (1963).

³ A. Bohm and R. A. Rashid, Phys. Letters **10**, 451 (1964).

⁴ B. Pal, Nuovo Cimento **36**, 1392 (1965).

⁵ F. Bulos *et al.*, Phys. Rev. Letters **13**, 486 (1964).

⁶ V. Z. Peterson *et al.*, University of California Radiation Laboratory Report No. UCRL-11575 (unpublished).

⁷ P. Auvil *et al.*, Phys. Letters **12**, 76 (1964); C. Lovelace, Nuovo Cimento **33**, 473 (1964).

⁸ W. Bruce Richards *et al.*, Phys. Rev. Letters **16**, 1221 (1966).

functions. The ξ are given by

$$\cosh \xi_1 = 1 + 8m_\pi^2/k^2, \quad (5)$$

$$\cosh \xi_2 = [W^2 - (m - m_\pi)^2]/2k^2 - 1, \quad (6)$$

$$\cosh \xi_1' = 1 + 2m_\pi^2/q^2, \quad (7)$$

$$\cosh \xi_2' = (W^2 - m^2 - 2m_\pi^2)/2q^2 - 1, \quad (8)$$

$$\cosh \xi_1'' = [(m_\pi + m_\eta)^2 - 2m^2 \\ + 2(k^2 + m^2)^{1/2}(q^2 + m^2)^{1/2}]/2kq, \quad (9)$$

$$\cosh \xi_2'' = [W^2 + m^2 - m_\pi^2 - m_\eta^2 \\ - 2(k^2 + m^2)^{1/2}(q^2 + m^2)^{1/2}]/2kq. \quad (10)$$

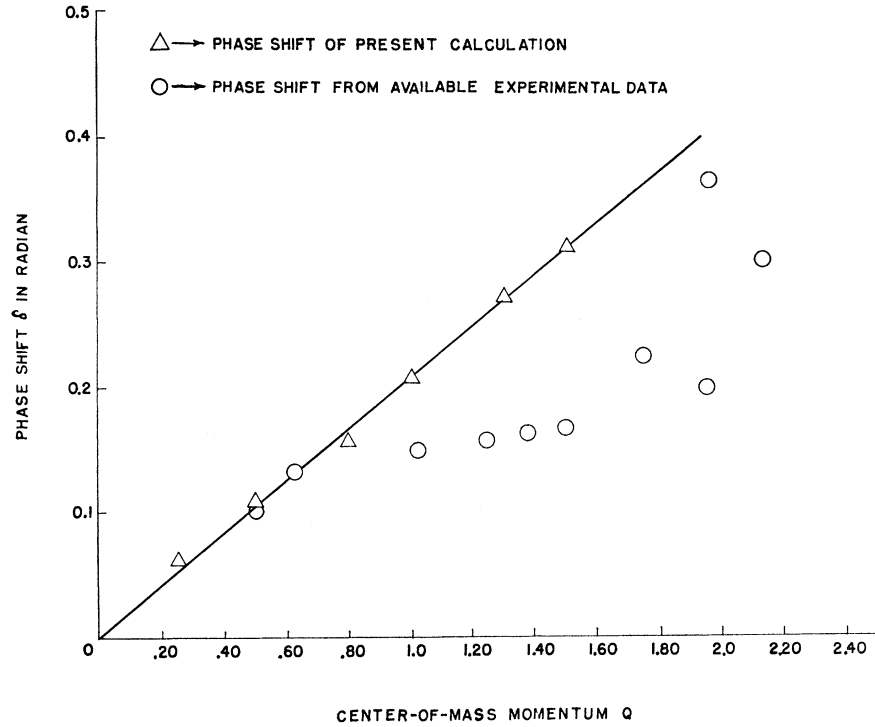


FIG. 1. Comparison of the $T=\frac{1}{2}$, $J=\frac{1}{2}^-$ phase shift calculated with the present model, and the experimental results of Miyake *et al.* (Ref. 11). Q is the pion momentum in the center-of-mass system in units of $m_{\pi}c$.

k and q are the center-of-mass momenta in the $\pi^- + p \rightarrow \pi^- + p$ and $\eta + n \rightarrow \eta + n$ channels, respectively. In (5)–(10), W is the total c.m. energy and m is the nucleon mass. m_{π} , m_{η} , and m_p denote the masses of the corresponding particles. The β 's are given by the relations based on the threshold behavior and the factorization theorem.⁹ The η - n trajectory may be written in the form

$$\text{Re}\alpha(W) = \frac{1}{2} + \epsilon(W - m_{\eta n}), \quad (11)$$

$$\text{Im}\alpha(W) = \frac{1}{2} \epsilon \Gamma \left(\frac{W - m - m_{\pi}}{m_{\eta n} - m - m_{\pi}} \right)^{\alpha_0}, \quad (12)$$

TABLE I. Comparison of the $T=\frac{1}{2}$, $J=\frac{1}{2}^-$ phase shift calculated with the present model, and the experimental results of Miyake *et al.* (Ref. 11).

Center-of-mass momentum of pion in units of $m_{\pi}c$	Phase shift δ in radians	
	Experimental value	Present calculation
0.50	0.100	0.108
0.63	0.143	0.128
1.02	0.150	0.210
1.25	0.158	0.258
1.38	0.163	0.263
1.50	0.168	0.310

⁹ S. K. Bose and S. N. Biswas, Phys. Rev. **135**, B1045 (1964).

where $\alpha_0 = \alpha(W)|_{k^2 \rightarrow 0}$. In the above, $m_{\eta n}$ and Γ are the mass and width of the η - n state. For the slope of the η - n trajectory, we take the same value as that found for the nucleon, $\epsilon = 0.4$. The complex scattering length for the reaction $\pi^- + p \rightarrow \eta + n$ is $a + ib$, where $a = 0.8304$ F and $b = 0.05$ F according to Peterson *et al.* The width Γ of the η - n state can be calculated very crudely in terms of a and b ; we obtain $\Gamma = b/\mu_{\eta} a^3 = 83.4$ MeV, where μ_{η} is the reduced mass of η and n . Now the expression for T_{11} becomes

$$T_{11} = \frac{C_1}{\alpha(W) - \frac{1}{2}} [1 + e^{(\alpha-1/2)(\xi_2 - \xi_1)}] + \text{Born terms.} \quad (13)$$

The constant C_1 may be determined, using unitarity, in terms of ϵ and Γ . This procedure yields $C_1 \simeq -\epsilon\Gamma/4k_{\eta n} \simeq -0.02$. The calculated value of the S -wave $T=\frac{1}{2}$ scattering length is $A_{11} = T_{11}(W)|_{k \rightarrow 0} \simeq 0.15$, which can be compared with the experimental data by Woolcock,¹⁰ $A_1 = 0.17 \pm 0.005$. The $S_{1/2}$ phase shifts calculated for the reaction $\pi^- + p \rightarrow \pi^- + p$ are compared with the experimental data by Miyake *et al.*¹¹ in Table I and Fig. 1. The quantity Q is the momentum of meson in units of $m_{\pi}c$ in the center-of-mass system.

¹⁰ W. S. Woolcock, in *Proceedings of the Aix-en-Provence International Conference on Elementary Particles, 1961* (Centre d'Etudes Nucléaires de Saclay, Seine et Oise, 1961), p. 459.

¹¹ K. Miyake *et al.*, Phys. Rev. **126**, 2188 (1962).