some terms in the sum (6.5) may vanish for some $\omega_l \neq \omega_k$. We will not discuss here the general case and will restrict ourselves to the case when all ω's are equal to each other, $\omega_i = \omega_j = \omega$. Taking this assumption into account, we get from (6.3) ,with the help of (6.4),

$$E_{j_1 \cdots j_{2n}}{}^{(n,n)}(\mathbf{r}_1 \cdots \mathbf{r}_{2n};\omega) = \sum_\pi E_{j_1 j_p}{}^{(1,1)}(\mathbf{r}_1 - \mathbf{r}_p, \omega) \cdots \times E_{j_n j_s}{}^{(1,1)}(\mathbf{r}_n - \mathbf{r}_s, \omega), \quad (6.6)$$

where $E_{j_k j_l}{}^{(1,1)}(\mathbf{r}_k - \mathbf{r}_l, \omega)$ are the second-order correlation functions introduced in Sec. 3. If all coordinates $\mathbf{r}_k$ and polarizations j_k coincide,

$$E_{j \cdots j}{}^{(n,n)}(\mathbf{r},\omega) = n![E_{j,j}{}^{(1,1)}(0,\omega)]^n. \quad (6.7)$$

Our considerations of higher-order correlation spectral tensors can also be performed in wave-vector space. The higher-order correlation tensors introduced in this way are expressible by the second-order correlation functions in a way similar to (6.4) and (6.7). Of course all these relations are a direct consequence of the fact that the probability distributions of the fluctuating fields of blackbody radiation are Gaussian in any Lorentz frame.

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Errata

Equations of Motion in Nonequilibrium Statistical Mechanics. II. Energy Transport, BALDWIN ROBERTSON [Phys. Rev. **160**, 175 (1967)].

1. Footnote 30 should read ". . . only the first term will survive. . . ".

2. For complete generality, Eqs. (A7) and (B1) should read

$$P(t)\bar{A}\sigma(t) = \bar{F}\sigma(t)\cdot\left[\frac{\delta\langle A\rangle_t}{\delta\langle F\rangle_t} - \left\langle\frac{\delta A}{\delta\langle F\rangle_t}\right\rangle_t\right], \quad (A7)$$

$$\mathcal{P}(t)A = F\cdot\left[\frac{\delta\langle A\rangle_t}{\delta\langle F\rangle_t} - \left\langle\frac{\delta A}{\delta\langle F\rangle_t}\right\rangle_t\right]. \quad (B1)$$

However, in almost all applications, the second term will be zero as assumed before.

3. Appendix C, third paragraph, the second sentence should be deleted and replaced by the sentence: Furthermore, even if we were given a $\lambda(t)$, the matrix $\langle F\bar{F}\rangle_t^{-1}$ would not exist if we use the definition (6).[43a]

4. The following footnote should be added:

[43a] Proof: Eqs. (5) and (6) give $\langle\bar{A}\rangle_t = 0$, and so Eqs. (5) and (C1) gives $\zeta\cdot\langle F\bar{F}\rangle_t = 0$. Since ζ is nonzero, the determinant of $\langle F\bar{F}\rangle_t$ must be zero.

5. The sentence including Eq. (C8) should read "Combining Eqs. (9), (1), and (C7), ..."

6. In Eq. (27), there should be a bar over the rightmost **J**.

Pressure Shifts of Highly Excited States of Atoms, M. H. MITTLEMAN [Phys. Rev. **162**, 81 (1967)]. There is an error in the derivation Eq. (6) for the local momentum of the electron. The term $-2\Delta E$ arises from the polarization of the M.P. by the Cs core and by the electron. Polarization is a second-order effect, so these effects are not additive. Instead of $2\Delta E$ in Eq. (6) we should have

$$-\sum_\beta\left[\frac{2\alpha}{Z_\beta{}^4} - \frac{2\alpha Z_\beta\cdot(Z_\beta - r)}{Z_\beta{}^3|Z_\beta - r|^3}\right].$$

The second term arises from the cross term between electron interactions and core interactions. Upon conversion of the sum to an integral this becomes

$$2\Delta E + 4\pi\alpha\rho U(r).$$

Replacing $2\Delta E$ in Eq. (6) by this expression modified Eq. (6) to the extent that $U(r)$ is multiplied by the factor $(1+4\pi\alpha\rho)$. The $U(r)$ arising in this way is multiplied by another factor of ρ, so that there will be no change in our results to the order of ρ given in the paper.

Similar considerations apply to Eq. (25).