TABLE II. Table for a, β, and ϵ for $PN \to P'N$.

λ_1	λ_2	a	b	β	ϵ
$+\frac{1}{2}$	$+\frac{1}{2}$	0	0	α	$-$
$+\frac{1}{2}$	$-\frac{1}{2}$	1	1	$\alpha-1$	$+$

fortunately, except for $\pi^+p \to \eta N^*$, the A_2 trajectory is masked by the π exchange in $\pi N \to \rho N$; $\pi N \to \rho N^*$ and $KN \to K^*N^*$, or it is accompanied by ρ exchange such as in $KN \to KN^*$. It would, therefore, be very helpful to have more detailed data for the reaction $\pi^+p \to \eta N^*$.

ACKNOWLEDGMENTS

It is a pleasure to express my gratitude to Professor Reinhard Oehme, who suggested the present subject, for stimulating discussions, and for continued guidance and encouragement. I also would like to thank the Institute of Technology, Bandung, and the Center for Developmental Change, University of Kentucky, for their arrangements.

TABLE III. Table for a, b, β, and ϵ for $PN \to VN$.

λ_V	λ_1	λ_2	a	b	β	ϵ
$+1$	$+\frac{1}{2}$	$+\frac{1}{2}$	1	1	$\alpha-1$	$-$
	$+\frac{1}{2}$	$-\frac{1}{2}$	0	2	$\alpha-1$	$+$
0	$+\frac{1}{2}$	$+\frac{1}{2}$	0	0	α	$+$
	$+\frac{1}{2}$	$-\frac{1}{2}$	1	1	$\alpha-1$	$+$
-1	$+\frac{1}{2}$	$+\frac{1}{2}$	1	1	$\alpha-1$	$+$
	$+\frac{1}{2}$	$-\frac{1}{2}$	2	0	$\alpha-1$	$+$

APPENDIX

In Tables II and III we list values of a, b, β, and ϵ for $PN \to P'N$ and $PN \to VN$, respectively.

Asymptotic forms for $N_{\lambda_i,\lambda_f}{}^{\alpha}P_{\beta}{}^{(a,b)}(z)$ are the following:

$$N_{0,0}{}^{\alpha}P_{\alpha}{}^{(0,0)}(z) \sim n_{\alpha}(2z)^{\alpha},$$
$$N_{1,0}{}^{\alpha}P_{\alpha-1}{}^{(1,1)}(z) \sim 2[\alpha/(\alpha+1)]^{1/2} n_{\alpha}(2z)^{\alpha-1},$$
$$N_{1,1}{}^{\alpha}P_{\alpha-1}{}^{(0,2)}(z) \sim 2[\alpha/(\alpha+1)] n_{\alpha}(2z)^{\alpha-1},$$
$$N_{1,-1}{}^{\alpha}P_{\alpha-1}{}^{(2,0)}(z) \sim 2[\alpha/(\alpha+1)] n_{\alpha}(2z)^{\alpha-1},$$

where

$$n_{\alpha} = (1/\sqrt{\pi})[\Gamma(\alpha+\tfrac{1}{2})/\Gamma(\alpha+1)].$$

Errata

Radiative Corrections. I. High-Energy Bremsstrahlung and Pair Production, KJELL MORK AND HAAKON OLSEN [Phys. Rev. **140**, B1661 (1965)]. The numerical values for F_1 given in Table I are incorrect. The correct values are as follows:

ω_1/ϵ_1	0.1	0.2	0.3	0.4	0.5	0.6	0.7	0.8	0.9
$F_1 \times 10^2$	0.095	0.19	0.30	0.43	0.60	0.80	1.08	1.50	2.30

We are indebted to Dr. H. D. Schulz for pointing out these errors to us.

$K_{\mu 3}$ and K_{e3} Form Factors at Finite Momentum Transfer, M. FITELSON AND E. KAZES [Phys. Rev. **159**, 1236 (1967)]. Equation (23a) should read

$$(a_1{}^2C_1 + a_2{}^2C_2)/\sqrt{2} = \langle \pi^0 | J''_3{}^1(0) | K^+ \rangle_{p=\infty}.$$

Equation (26) should read

$$-\frac{\cos Ma_1}{\cos Ma_2} = \frac{B_2 - b_2C_2M \tan Ma_2}{B_1 - b_1C_1M \tan Ma_1}.$$

Castillejo-Dalitz-Dyson Poles and Asymptotic Fields, STANLEY JERNOW AND EMIL KAZES [Phys. Rev. **160**, 1428 (1967)]. A typographical error appeared in Eq. (2.2) which gave the form of the interaction Hamiltonian. The equation should read

$$G^{\dagger} \equiv \int d\mathbf{p}\, f(\omega_p)\theta^{\dagger}(p), \quad \omega_p = (\mu^2 + p^2)^{1/2}. \qquad (2.2)$$

Also, the left-hand side of Eq. (2.21) should be the time derivative of the field and should read $-\dot{\beta}_j{}^{\dagger}(t)$.