

## Mass and Beta Decay of $\text{Na}^{20}\dagger$

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(Received 22 May 1967; revised manuscript received 27 July 1967)

The  $\beta$  and  $\gamma$  spectra of  $\text{Na}^{20}$  have been measured with a wedge-gap magnetic spectrometer and a 3-cm<sup>3</sup> lithium-drifted germanium detector, respectively.  $\text{Na}^{20}$  was produced by a  $(p,\alpha n)$  reaction on a natural magnesium target in the UCLA cyclotron. The mass excess of  $\text{Na}^{20}$  is measured to be  $6863 \pm 40$  keV. The half-life of  $\text{Na}^{20}$  is  $0.408 \pm 0.006$  sec. Two  $\beta$  branches of  $11.25 \pm 0.04$  and  $5.55 \pm 0.15$  MeV were observed, and their relative intensities measured to be  $100 \pm 2$  and  $9 \pm 1$ , respectively. The more energetic  $\beta$  branch feeds the first excited state of  $\text{Ne}^{20}$ ; the excitation energy is determined to be  $1633.2 \pm 1.0$  keV from a study of its  $\gamma$  transition. The 5.55-MeV  $\beta$  branch establishes a link with the known delayed- $\alpha$  spectra from the highly excited states of  $\text{Ne}^{20}$ , and allows us to propose a decay scheme for  $\text{Na}^{20}$ . Limits on the intensities of possible unobserved  $\beta$  transitions are discussed.

### 1. INTRODUCTION

INVESTIGATION of weak  $\beta$  branches from short-lived isotopes is extremely difficult from a measurement of a  $\beta$  spectrum alone. The study of delayed-proton and delayed- $\alpha$  spectra, whenever such processes are energetically possible, can significantly reduce this difficulty. Conversely, these particular studies require a link to the  $\beta$  spectrum in order to enable one to extract  $\beta$  branching ratios from them.

This work, concerned with the measurement of the  $\text{Na}^{20}$   $\beta$  spectrum, establishes such a link to the delayed- $\alpha$  measurements<sup>1,2</sup> from excited states of  $\text{Ne}^{20}$ , allowing us to propose a decay scheme of  $\text{Na}^{20}$  with accurate  $ft$  values. From these  $ft$  values, the  $T=1$ ,  $T_z=0$  analog state of the mass-20 triplet is unambiguously determined. In addition, the measurement of the highest endpoint energy of the  $\text{Na}^{20}$   $\beta$  spectrum can determine the mass excess of  $\text{Na}^{20}$  and test the validity of the mass formula in the isobaric triplet. The first experimental value of the  $\text{Na}^{20}$  mass excess<sup>3</sup> has been in disagreement with more recent determinations<sup>4,5</sup> by more than 1 MeV. This experiment agrees with these later values and improves their accuracy.

A decay scheme of  $\text{Na}^{20}$  is constructed from the data of this measurement and those of delayed- $\alpha$ -particle studies.<sup>1</sup> On the basis of the deduced  $ft$  values, spin and parity assignments are made, confirming previous determinations.<sup>6</sup> From the  $\gamma$ -ray measurements and the systematics of  $\log ft$  values of  $\beta$  transitions, upper limits on some unobserved  $\beta$  branches are discussed.

### 2. EXPERIMENTAL TECHNIQUES

Since the measurement of  $\beta$  spectra from short-lived isotopes requires the accumulation of data from many short bombardments of a given target, a number of experimental difficulties arise. Beam-intensity fluctuations during the target irradiation and the slow build up of long-lived reaction products on the target introduce normalization problems. Also, other short-lived isotopes may be produced and mask the desired activity.

Some of these difficulties are overcome by exploiting the focusing properties of a single wedge-gap magnetic spectrometer. Several detectors covering a sizable momentum range of a  $\beta$  spectrum (15%) can be placed along the *focal line* of the instrument. In this manner one has a way to normalize the data by partially overlapping the momentum interval analyzed in two successive runs.

The problems of extraneous half-lives can be handled by independently multiscaling the output of each of the detectors of the spectrometer. This is, of course, conveniently done in an on-line computer. After sufficient statistical accuracy is achieved in one run, the relevant data are readily accessible for the decomposition of the time spectrum into intensities and half-lives by a least-squares fitting program. Using such a system allows one to extract a normalized  $\beta$  spectrum from a strong "background" of competing activities.

The details of the spectrometer operation have been published previously.<sup>7</sup> Four detectors (Si) with a momentum resolution of 1% were used and covered a relative momentum range of 12%. The spectrometer was calibrated to 0.1% in momentum. The targets exposed to the internal beam of the UCLA sector focused cyclotron were extracted to the spectrometer source position in 250 msec with the use of a fast, air operated, pneumatic rabbit system. The multiscaling, half-life separation, and decomposition of the Kurie plot of the  $\beta$  spectrum was done on-line with an SDS 925 computer.

† Supported in part by the U. S. Office of Naval Research under Contract Nonr 233-44.

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<sup>6</sup> T. Lauritsen and F. Ajzenberg-Selove, in *Nuclear Data Sheets*, compiled by K. Way *et al.* (U. S. Government Printing Office, National Academy of Sciences—National Research Council, Washington 25, D. C.), NRC 61-5,6-307.

<sup>7</sup> A. J. Armini, R. M. Polichar, and J. W. Sunier, Nucl. Instr. Methods **48**, 309 (1967).

### 3. EXPERIMENTAL RESULTS

$\text{Na}^{20}$  was produced by the  $(p, \alpha n)$  reaction at 45 MeV on a natural, self-supporting magnesium target. The bombardment energy, considerably above the reaction threshold of 25 MeV, was chosen to minimize the  $(p, n)$  reactions on the three magnesium isotopes  $\text{Mg}^{24}$ ,  $\text{Mg}^{25}$ , and  $\text{Mg}^{26}$ .

The half-life was measured to be  $408 \pm 6$  msec, which is an average of over 100 independent measurements. The Kurie plot of the  $\beta$  spectrum associated with this half-life is shown in Fig. 1, for positron energies between 3.5 and 11.5 MeV. Above 3.5 MeV the  $\text{Al}^{24}$  (2.08-sec) and  $\text{Al}^{24m}$  (0.13-sec) contaminants were adequately separated. Below this limit, strong branches from  $\text{Al}^{24}$ ,  $\text{Al}^{25}$ ,  $\text{Mg}^{23}$ , and  $\text{Al}^{26m}$  could not be adequately separated from the  $\text{Na}^{20}$  activity, mainly because of their overwhelming intensity. This prevented the measurement of the superallowed transition from  $\text{Na}^{20}$  to its analog. The analysis of the Kurie plot gave three branches of allowed shape. Their endpoint energies are  $11.25 \pm 0.04$ ,  $7.15 \pm 0.10$ , and  $5.55 \pm 0.15$  MeV. Their relative intensities are  $100 \pm 2$ ,  $17 \pm 2$ , and  $9 \pm 1$ , respectively.

The 7.15-MeV branch is inconsistent with the known positive-parity levels in  $\text{Ne}^{20}$ , and too strong to be a first forbidden transition. This fact prompted an investigation of the  $\text{Na}^{20}$   $\beta$  spectrum through the use of another reaction. The production of  $\text{Na}^{20}$  was attempted by the  $(p, tn)$  reaction on natural sodium, with a threshold of 32 MeV. The  $\beta$  spectrum was severely contaminated by the 0.8-sec  $\beta$  activity of  $\text{B}^8$ , produced by the  $(p, \alpha n)$  reaction on the residual mineral oil in the sodium. However carefully the target was prepared and cleaned, the relative yield of these two reactions was such that the  $\text{B}^8$  could not be eliminated. Because of the relatively low counting rate obtained in this experiment, the half-life separation was judged too uncertain and the total  $\beta$  spectrum was recorded for 400 msec and analyzed. The

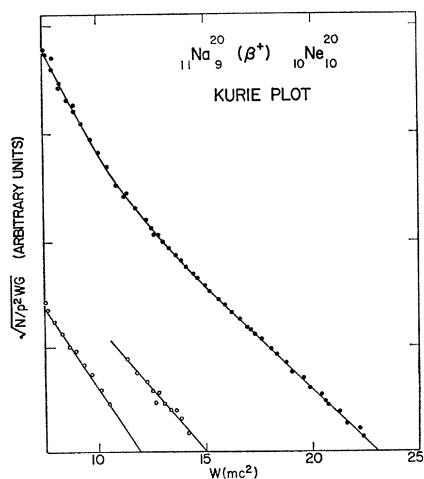


FIG. 1. The  $\text{Na}^{20}$  Kurie plot from the  $\text{Mg}^{24}(p, \alpha n)$  reaction on natural magnesium. The branch with the  $W=15$  endpoint is an unidentified contaminant.

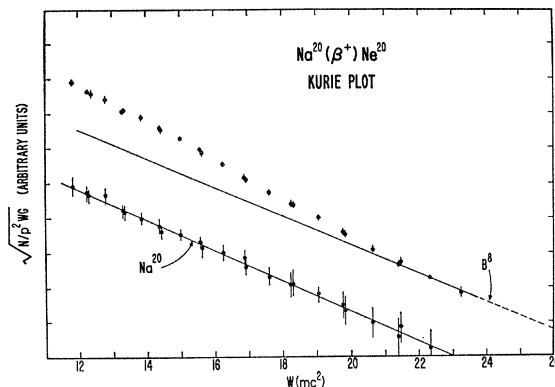


FIG. 2. The  $\text{Na}^{20}$  Kurie plot from the  $\text{Na}^{23}(p, tn)$  reaction. The extension of the  $\text{B}^8$  contaminant spectrum at the higher energies has been omitted.

resulting Kurie plot and its separation are shown in Fig. 2. After subtraction of the  $\text{B}^8$  component, the remaining  $\text{Na}^{20}$  spectrum is still sufficiently accurate to definitely rule out any 7-MeV transition of intensity comparable to what had been observed from the Mg target. The reason for this 7-MeV branch and its nature must be due to an unknown contaminant of half-life very close to 0.4 sec. This contaminant is presently under further investigation. A spectroscopic analysis of the target material showed that it consisted of 99.9% Mg.

The  $\gamma$  rays from  $\text{Na}^{20}$  were observed using a magnesium target. A 3.2-cm<sup>3</sup> lithium-drifted germanium detector with an energy resolution of 3.6 keV for the 1.33-MeV  $\gamma$  transition of  $\text{Co}^{60}$  was coupled to the rabbit irradiation system. A 1024-channel time-sequenced spectrum was used for the pulse-height analysis. Only one  $\gamma$  ray was found which had a half-life consistent with the 0.4 sec of  $\text{Na}^{20}$ . Its energy was measured to be  $1633.2 \pm 1.0$  keV from an external calibration with  $\text{Co}^{60}$ ,  $\text{Ga}^{66}$ , and  $\text{Na}^{24}$  sources, and an internal calibration with the two standard  $\gamma$  transitions of  $\text{Mg}^{24}$ , present in the target from the competing  $\text{Mg}^{24}(p, n)\text{Al}^{24}$  reaction. This result is in excellent agreement with the value of  $1632.6 \pm 0.8$  keV reported by Alburger and Jones.<sup>8</sup>

### 4. THE $\text{Na}^{20}$ DECAY SCHEME AND ITS DISCUSSION

The 11.25-MeV end point and the subsequent 1633.2-keV  $\gamma$  transition allow us to determine the mass excess of  $\text{Na}^{20}$ . The result obtained is  $6863 \pm 40$  keV and is in excellent agreement with the measurements of Donovan and Parker<sup>4</sup> and Pehl and Cerny.<sup>5</sup>

The necessary link to the delayed- $\alpha$  measurements of Polichar *et al.*<sup>1</sup> is established through the intensity ratio of the 11.25- and 5.55-MeV  $\beta$  transitions. The 5.55-MeV  $\beta$  branch feeds the 7.43-MeV level of  $\text{Ne}^{20}$ , which breaks up into  $\text{O}^{16}$  and the strongest delayed- $\alpha$  line observed

<sup>8</sup> D. E. Alburger and K. W. Jones, Phys. Rev. **149**, 743 (1966).

TABLE I. Combined data of this measurement and the delayed- $\alpha$  spectra of Polichar *et al.*<sup>a</sup>

Level in Ne <sup>20</sup> Energy (MeV) $I^\pi$	$\beta$ in (MeV) <sup>b</sup>	Branching ratio %	Log $ft$
1.6332±0.0010 2 <sup>+</sup>	$\beta_0$ 11.25±0.04	90.0	4.89±0.05
7.43 ±0.01 2 <sup>+</sup>	$\beta_1$ 5.45±0.04	8.1	4.49±0.10
7.84 ±0.03 2 <sup>+</sup>	$\beta_2$ 5.04±0.05	0.38	5.51±0.13
8.74 ±0.03 1 <sup>-</sup>	$\beta_3$ 4.14±0.05	0.024	6.35±0.16
9.48 ±0.02 2 <sup>+</sup>	$\beta_4$ 3.40±0.05	0.11	5.27±0.11
10.28 ±0.01 2 <sup>+</sup>	$\beta_5$ 2.60±0.04	1.38	3.77±0.10
10.86 ±0.02 2 <sup>+</sup>	$\beta_6$ 2.02±0.05	0.097	4.40±0.13
11.28 ±0.04 2 <sup>+</sup>	$\beta_7$ 1.60±0.06	0.032	4.53±0.16

<sup>a</sup> Reference 1.

<sup>b</sup> The endpoint of  $\beta_0$  is the measured one. The others are deduced from the energy levels obtained in the  $\alpha$  measurement and correspond to the transitions  $\alpha_1$  to  $\alpha_7$  of Ref. 1.

from the Na<sup>20</sup> decay. Making use of the intensity ratio reported in the delayed- $\alpha$  measurements<sup>1</sup> and of the levels of Ne<sup>20</sup> from which they originate, it is possible to infer the relevant  $\beta$ -branching ratios and to propose a complete Na<sup>20</sup> decay scheme. Table I presents a summary of both experimental data, where the intensities of the  $\alpha$  transitions<sup>1</sup> have been renormalized to the total  $\beta$  decay. The Na<sup>20</sup> decay scheme is shown in Fig. 3.

The fact that only the 2<sup>+</sup> level of the ground-state rotational band of Ne<sup>20</sup> is populated by  $\beta$  decay determines the spin and parity of Na<sup>20</sup> to be 2<sup>+</sup>. This result is supported by similar observations<sup>6,8</sup> from the analog decay of F<sup>20</sup>. With the exception of the 8.74-MeV level, all the levels populated by  $\beta$  decay and reported in Table I have spin and parity 2<sup>+</sup>. This is deduced from the  $\log ft$  values of the involved  $\beta$  transitions and from the  $\alpha$ -decay selection rules. The  $\beta_3$  transition feeding the 8.74-MeV level has a  $\log ft$  of 6.35. The transition is then

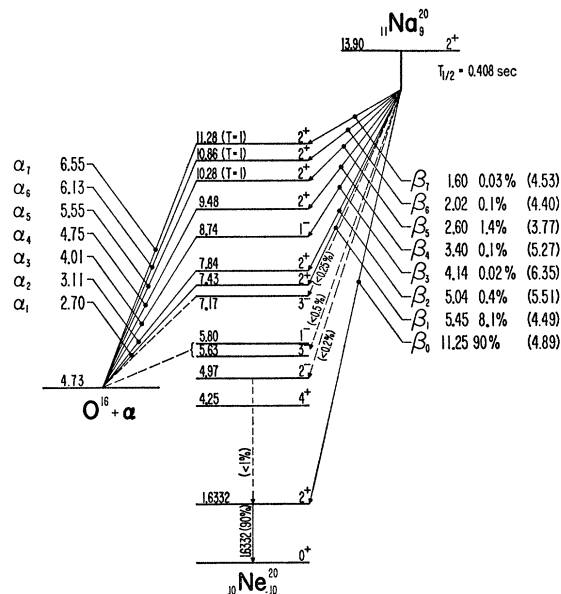


FIG. 3. The proposed Na<sup>20</sup> decay scheme. The broken lines indicate unobserved transitions whose branching might be higher than 0.1% (see text and Ref. 6). The  $\alpha$  transition energies and intensities are taken from Ref. 1.

likely to be first forbidden, in agreement with the assignment of 1<sup>-</sup> for the 8.74-MeV level, previously inferred by Lauritsen and Ajzenberg-Selove.<sup>6</sup>

The 10.28-MeV level of Ne<sup>20</sup> was first identified by Pearson and Spear<sup>9</sup> in an experiment on radiative capture of  $\alpha$ 's by O<sup>16</sup>. The  $\log ft$  value of the deduced 2.60-MeV  $\beta$  transition confirms that this level is the  $T=1$ ,  $T_z=0$  analog of the mass-20 triplet. The 2.60-MeV  $\beta$  branch is clearly superallowed, although slower than one might expect from a comparison with the  $ft$  value of similar transitions in neighboring nuclei. This retardation may be due to some  $T=0$  impurity in the 10.28-MeV level. This must be true to some extent because otherwise the isospin selection rule would forbid its  $\alpha$  decay.

One may draw several conclusions from the data concerning the  $\beta_6$  and  $\beta_7$  transitions. Since both  $\beta$  branches are followed by  $\alpha$  decay and have  $ft$  values which are allowed, it follows that the 10.86-MeV and 11.28-MeV states of Ne<sup>20</sup> have  $J^\pi=2^+$ . If one compares the excitation of these levels relative to the lowest  $T=1$  state at 10.28 MeV, one finds a close correspondence with the  $T=1$  levels at 0.65 MeV and 0.99 MeV in F<sup>20</sup>. It is tempting to suggest that these two pairs of levels are in fact  $T=1$  isobaric analog states, which would be supported by the  $ft$  values of both  $\beta$  transitions involved. This would in turn suggest  $J^\pi=2^+$  for the corresponding levels in F<sup>20</sup>.

From our  $\gamma$  measurements, we can deduce some limits on the intensities of possible first-forbidden  $\beta$  transitions to negative-parity states of Ne<sup>20</sup>. The first candidate for such a transition is the 2<sup>-</sup> level at 4.97 MeV. This state is known to decay with a branching >95% by a cascade  $\gamma$  transition to the 1.63-MeV level. Our measurement gives a limit of  $\log ft > 6.1$  to the corresponding  $\beta$  transition. This result can be compared to the value of  $\log ft=6.9$ , reported by Alburger and Jones<sup>7</sup> for the analog  $\beta$  transition from F<sup>20</sup>. From charge independence of nuclear forces, analog  $\beta$  transitions should have nearly equal comparative half-lives. Assuming equality, one can deduce that the branching ratio of a  $\beta$  transition to the 4.97 (2<sup>-</sup>)-MeV level is  $<1.6 \times 10^{-3}$ .

The next two candidates for first-forbidden  $\beta$ -decay feeding are the 3<sup>-</sup> state at 5.63 MeV and the 1<sup>-</sup> state at 5.80 MeV. Both states are known to decay predominantly by  $\alpha$ -particle emission<sup>6,10</sup> to the ground state of O<sup>16</sup>. These  $\alpha$  particles have not been observed in the work of Polichar *et al.*<sup>1</sup> and recent measurements<sup>11</sup> have placed an upper limit of 0.5% to their total intensity relative to the  $\beta$  decay of Na<sup>20</sup>. This result gives to the corresponding  $\beta$  branches a  $\log ft > 6.5$ , which is in reasonable agreement with the systematics of  $\log ft$  values for  $\Delta J=1$  first-forbidden  $\beta$  transitions.

<sup>9</sup> J. D. Pearson and R. H. Spear, Nucl. Phys. 54, 434 (1964).

<sup>10</sup> J. D. Pearson, E. Almqvist, and J. A. Kuehner, Can. J. Phys. 42, 489 (1964).

<sup>11</sup> J. E. Steigerwalt (private communication).

From our  $\gamma$  measurements, a  $\beta$  branching to the 7.03(4<sup>-</sup>)-MeV level could be as high as 2%. This would give a  $\log ft$  of 5.0 which is about three orders of magnitude too small. It is reasonable to infer, from consideration of  $\log ft$  systematics, that such a transition has a probable branching ratio  $<3 \times 10^{-5}$ . Similar arguments could be used to estimate a possible  $\beta$  feeding of the  $\alpha$  unstable states at 7.17(3<sup>-</sup>), 8.90(1<sup>-</sup>), and 9.16(3<sup>-</sup>) MeV. Assuming a  $\log ft > 6.0$ , the corresponding branching ratios would be  $<0.25$ ,  $<0.06$ , and  $<0.04\%$ , respectively. These estimates are realistic and agree with the fact that those transitions were not observed

in delayed- $\alpha$  measurements. Along the same lines, branching to 0<sup>+</sup> and 4<sup>+</sup> states by second-forbidden transitions ought to be  $<10^{-7}$ . This is well below the limit of 1–2% placed upon them by our  $\gamma$  measurements and the estimate of  $<0.05\%$  obtained for a possible  $\beta$  transition to the 0<sup>+</sup> ground state of Ne<sup>20</sup>.

#### ACKNOWLEDGMENTS

It is a pleasure to thank Ralph de Vries and John Steigerwalt for their comments and help during the measurements.

### Beta Decay of <sup>37</sup>S†

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(Received 5 July 1967)

Radioactive <sup>37</sup>S has been prepared by the <sup>37</sup>Cl(*n,p*)<sup>37</sup>S reaction and its decay products have been observed with a magnetic  $\beta$ -ray spectrometer and with NaI(Tl) crystal and Ge(Li) crystal  $\gamma$ -ray spectrometers. The  $\beta$ -ray groups have endpoint energies, relative intensities, and comparative half-lives of  $4.75 \pm 0.04$  MeV ( $5.6 \pm 0.6\%$ ),  $\log f_{\beta}t = 8.01$ ;  $1.64 \pm 0.04$  MeV, 94.0%,  $\log f_{\beta}t = 4.38$ ; and  $1.04 \pm 0.04$  MeV,  $\sim 0.4\%$ ,  $\log f_{\beta}t = 5.65$ . The  $\gamma$  rays have energies and relative intensities of  $3.107 \pm 0.002$  MeV (99.6%) and  $3.708 \pm 0.004$  MeV ( $\sim 0.4\%$ ). No other  $\gamma$  rays were observed; an upper limit of 0.5% is placed on their relative intensity. Coincidence counting places an upper limit of 1% on the possibility of a cascade transition from the 3.107-MeV state of <sup>37</sup>Cl. Spin-parity assignments for <sup>37</sup>Cl are  $\frac{3}{2}^-$  for the 3.107-MeV state and  $\frac{5}{2}^-$  for the 3.708-MeV state. The ground state of <sup>37</sup>S has odd parity and probable spin  $\frac{7}{2}$ .

#### INTRODUCTION

IN this investigation, NaI(Tl) scintillation crystals, a Ge(Li) crystal detector, and a magnetic  $\beta$ -ray spectrometer have been used in determining the energies and relative intensities of the  $\beta$  rays and the  $\gamma$  rays which are emitted in the  $\beta$  decay of <sup>37</sup>S. A primary purpose was to locate negative-parity excited states of the daughter  $d_{3/2}$  subshell nucleus <sup>37</sup>Cl.

#### Previous Studies

Before this study was initiated, there had been reported<sup>1,2</sup>  $\beta$ -ray transitions to the ground state of <sup>37</sup>Cl (4.7 MeV,  $\sim 10\%$ ) and to a 3.1-MeV excited state (1.6 MeV,  $\sim 90\%$ ). Only a 3.1-MeV  $\gamma$  ray had been observed; cascading transitions or another  $\beta$ -ray transi-

tion which would produce  $\gamma$  rays of energy of less than 2 MeV had been reported<sup>2</sup> to be of less than 1% of the intensity of the 3.1-MeV  $\gamma$  ray. The  $\log ft$  value of 7.3 for the ground-state transition was recognized as being consistent with a unique first-forbidden transition; similarly, the  $\log ft$  value of 4.2 for the transition to the 3.1-MeV state was consistent with an allowed transition. In <sup>37</sup>Cl, the ground-state spin-parity was well-established<sup>3</sup> as  $\frac{3}{2}^+$ . Magnetic analysis of the charged particles from the <sup>37</sup>Cl(*p,p'*)<sup>37</sup>Cl\* reaction had made possible the identification of states<sup>4,5</sup> at 0.835 (conflicting reports), 1.72, 3.087, and 3.105 MeV. The doubtful 0.835-MeV state was not observed in the <sup>37</sup>Cl(*p,p'*) reaction<sup>5</sup> with the use of separated isotopes of Cl, it was not observed in the <sup>37</sup>Cl(*n,n'*) $\gamma$  reaction,<sup>6</sup> and it was not observed in the <sup>40</sup>Ar(*p, $\alpha'$ ) $\gamma$  reaction.<sup>7</sup>*

† Supported in part by the National Science Foundation.

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