Equating constant terms on both sides of the energy-balance Eq. (A5) we find that

$$2\omega\tau_0\sum_{l=l'}\sum(S_lT_{l'}-S_{l'}T_l)l'$$
$$=F_0X_0^{\varphi}+2F_1U_z+(A\tau_0/kT_L)S_0$$
$$-\frac{A\tau_0}{kT_L}\{S_0^3+\tfrac{3}{4}\sum_{l''=l+l'}\sum\sum(S_lS_{l'}S_{l''}+S_lT_{l'}T_{l''})\}, \quad \text{(A14)}$$

which reduces in zeroth order to

$$F_0X_0^{\varphi}+2F_1U_z+(A\tau_0/kT_L)S_0(1-S_0^2)=0. \quad \text{(A15)}$$

Substituting X_0^{φ} from (A12) and U_z from (A13) we get

$$\frac{F_0^2B(B^2+\nu^2+\nu_c^2)}{(\nu_c^2-\nu^2)^2+2B^2(\nu^2+\nu_c^2)+B^4}+\frac{2F_1^2}{B}-\frac{A\tau_0}{kT_L}S_0(1-S_0^2)=0. \quad \text{(A16)}$$

This last equation can be used to find $S_0^2=\gamma$, which in turn can be inserted into an expression for the current as follows:

$$j=\sigma E=qn\mu_0{}^{\rm ac}\gamma^{-1/2}\int_0^{\infty}\frac{\gamma^2x^3e^{-x}dx}{(\gamma x)^2+l(\gamma x)^{3/2}}. \quad \text{(A17)}$$

Equation (A17) takes into account acoustic and neutral impurity scattering mechanisms only and can be evaluated with the aid of a digital computer. Results are shown in Figs. 4 and 5 and discussed in Sec. II.

Errata

Interband Contributions to Optical Harmonic Generation at a Metal Surface, SUDHANSHU S. JHA AND CHINDU S. WARKE [Phys. Rev. **153**, 751 (1967)] It is not correct to say that the contribution to the bilinear volume current density due to the second term in Eq. (3.33) involving $\mathfrak{B}(\omega)$ is a purely three-band contribution. It also contains a nonzero term involving only two bands. For low frequencies this term approximately changes m to m^* in Eq. (3.39).

In the right-hand sides of Eqs. (3.14)–(3.16), $\pm$ signs should be inserted. ω should be changed to 2ω in the left-hand sides of Eqs. (5.18) and (5.19).

Morse-Potential Evaluation of Second- and Third-Order Elastic Constants of Some Cubic Metals, R. C. LINCOLN, K. M. KOLIWAD, and P. B. GHATE [Phys. Rev. **157**, 463 (1967)]. It has been brought to our attention that the values in Table II of C_{11}, C_{12}, and C_{44} for sodium attributed to Daniels are in error. These values should be $C_{11}=0.739$, $C_{12}=0.622$, and $C_{44}=0.419$ instead of $C_{11}=0.738$, $C_{12}=0.521$, and $C_{44}=0.661$.

Localized-Mode Energy Losses in Large Excursions, J. H. WEINER [Phys. Rev. **152**, 597 (1966)]. There is a gap in the argument leading to the simplified form Eqs. (3.26) of the general equations of motion in curvilinear coordinates Eqs. (3.16). These simplifications for large N were based only on the vanishing of the individual terms of certain sums as $N\rightarrow\infty$. However, the limits of these sums do not vanish in general when $N\rightarrow\infty$. Equation (3.25) and the statement $\omega_{\alpha\beta}\equiv 0$ should be regarded therefore as additional simplifying assumptions on the potential-energy function $V(\theta,s)$.