

Composite Particle Model for the Nucleon and the (3,3) Resonance, J. S. BALL AND D. Y. WONG [Phys. Rev. **133**, B179 (1964)]. An error has been discovered in the computer program used to calculate the results given in Table I. After correction Table I should read as follows:

	γ_{33}	γ_1	m	g^2	Γ_{33}	a_1	a_3	c_{11}	c_{31}	c_{13}	c_{33}	d_{13}	d_{33}	W_c
1	0.05	-1.0	7.26	24.0	1.01	-0.501	-0.192	-0.620	-0.056	-0.048	0.298	-0.0006	0.0008	18.6
2	0.05	-0.5	7.22	21.0	1.00	-0.451	-0.185	-0.484	-0.051	-0.051	0.297	-0.0011	0.0011	18.2
3	0.05	0	7.52	14.0	1.02	-0.369	-0.194	-1.06	-0.041	-0.059	0.303	-0.0014	0.0013	18.6
4	0.06	-1.0	7.47	23.4	1.00	-0.477	-0.168	-1.20	-0.056	-0.050	0.288	-0.0007	0.0009	18.9
5	0.06	-0.5	7.50	16.9	1.01	-0.425	-0.167	-1.17	-0.050	-0.055	0.302	-0.0012	0.0010	18.4
6	0.06	0	7.82	12.0	1.05	-0.338	-0.195	1.44	-0.036	-0.064	0.304	-0.0015	0.0012	18.6
7	0.07	-1.0	7.85	15.9	0.96	-0.400	-0.155	+1.51	-0.053	-0.057	0.267	-0.0007	0.0009	20.5
8	0.07	-0.5	7.7	20.6	1.04	-0.367	-0.154	...	-0.050	-0.059	0.297	-0.0012	0.0010	18.5
9	0.07	0	8.14	14.2	1.05	-0.296	-0.197	0.482	-0.030	-0.069	0.307	-0.0015	0.0012	18.7

The only qualitative change in the results is that the nucleon is not sufficiently bound over the entire range of coupling constants investigated.

We would like to thank G. L. Shaw and P. W. Coulter of Stanford University for bringing this error to our attention.

Pion Exchange Currents in Deuteron Photodisintegration Dispersion Theory, MALCOLM H. SKOLNICK [Phys. Rev. **136**, B1493 (1964)]. Equation (4.33) should read

$$A(p) = e^{\delta(v)} b \left\{ A_0(p) + r(p) + \frac{2\bar{R}M}{r_s(\gamma^2 - \alpha_+^2)} \left(\frac{p^2 + \alpha_+^2}{p^2 + \alpha_+^2} \right)^{1/2} \left[\frac{p^2 + \gamma^2}{l_c^2 - \alpha_+^2} \left(\frac{\alpha_+}{p^2 + \alpha_+^2} + \frac{l_c(\gamma^2 - \alpha_+^2)}{(p_c^2 + \gamma^2)(p^2 + l_c^2)} \right) \right] - \frac{\gamma}{p_c^2 + \gamma^2} \right\}. \quad (4.33)$$

Similarly Eqs. (4.39) and (4.41) should read

$$A_{\text{pole}} = \frac{\bar{R}M}{p^2 + l_c^2} + \frac{2\bar{R}M}{r_s(\gamma^2 - \alpha_+^2)} \left(\frac{p^2 + \alpha_+^2}{p^2 + \alpha_+^2} \right)^{1/2} \left\{ \left(\frac{p^2 + \gamma^2}{l_c^2 - \alpha_+^2} \right) \left[\frac{\alpha_+}{p^2 + \alpha_+^2} + \frac{l_c(\gamma^2 - \alpha_+^2)}{(p_c^2 + \gamma^2)(p^2 + l_c^2)} \right] - \frac{\gamma}{p_c^2 + \gamma^2} \right\}, \quad (4.39)$$

$$a(v_c) = \frac{1}{v_c + B} + K \left\{ \left(\frac{\gamma^2}{v_c + L/M} \right) \left[(1/\alpha_+) + \frac{L/M}{v_c(Mv_c + \gamma^2)^{1/2}} \right] - \gamma/v_c \right\}. \quad (4.41)$$

Electron Spin-Echo Envelope Modulation, L. G. ROWAN, E. L. HAHN, AND W. B. MIMS [Phys. Rev. **137**, A61 (1965)]. Directly following Eq. (10), the next line gives a relationship $2\hbar\omega_1 = g\beta H_{\text{rf}}$. The symbol g should be replaced by $g' = (g_1 g_{11}/g)$ for the case where H_{rf} is perpendicular to the plane containing H_0 and the crystalline c axis. In the next line below, $(2g\beta)$ should be replaced by $(g'\beta/2)$. The symbol g appearing in Eq. (15) and everywhere in the Appendix should be replaced by g' .

This change does not effect the physical results of echo envelope determinations because the over-all amplitude, determined by the g' factor associated with off-diagonal transition matrix elements, is normalized to arbitrary units of amplitude.

Internal Field in General Dipole Lattices, F. W. DE WETTE AND G. E. SCHACHER [Phys. Rev. **137**, A78 (1965)]. In Eq. (23), for $(j_2 + j_2\xi_2)$, read $(j_2 + j_3\xi_2)$. In Eqs. (44) and (45), replace $\Omega_{j_3}(\mu_1, \mu_2)$ by $\tilde{\Omega}_{j_3}(\mu_1, \mu_2)$. The definition of $\tilde{\Omega}_{j_3}$ is obtained from Eq. (24) by placing a minus sign in front of the first term inside the curly brackets. In Eqs. (51) and (52), replace $\Omega_0(\mu_1, \mu_2)$ by $\tilde{\Omega}_0(\mu_1, \mu_2)$. The definition of $\tilde{\Omega}_0$ is obtained from Eq. (25) by placing a minus sign in front of the first term inside the curly brackets. The

exponent in the right-hand side of Eq. (72) should read

$$-2\pi i \left\{ [j_1 + (\lambda_3 + j_3)\xi_1]\mu_1 + [j_2 + (\lambda_3 + j_3)\xi_2]\mu_2 + k_3\lambda_3 \right\}.$$

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Pressure Dependence of the Emission from Ga-(As_{1-x}P_x) Electroluminescent Diodes, G. E. FENNER [Phys. Rev. **137**, A1000 (1965)]. The expression $\eta(x) = \eta(0) \exp -x/\Lambda$ following Eq. (1) should read $\eta(x) = \eta(0) - x/\Lambda$. The factor $\eta(0)$ in Eq. (2) should be dropped and the calculated curves in Fig. 5 shifted by the appropriate factor $\eta(0)$.

Weak Interactions and Self-Consistent Theories, MAHIKO SUZUKI [Phys. Rev. **136**, B769 (1964)]. The statement made at the end of the paper, "present arguments lead to the contradictions independently of the behavior of the weak vertices near the light cone" is erroneous. In fact, if one admits a singular behavior like $1/q^2$, implying a massless boson, the theorem in the paper fails and therefore one can avoid the claim that self-consistent theories