Errata

Space-Charge Limited Current Relation in High-Pressure Gas Diodes, R. FORMAN [Phys. Rev. **123**, 1537 (1961)]. In our paper we derived an expression, consistent with our experimental data, showing that the space-charge current, J, in a high-pressure gas diode varies as $V^{\frac{3}{2}}$ and $p^{-\frac{1}{2}}$ (V and p are voltage and pressure, respectively). Subsequently, it was brought to our attention by Dr. H. F. Ivey that he had theoretically derived[1] similar space-charge expressions for conduction in dense gases showing that $J\propto V^{\frac{3}{2}}$ for the case of very high electric fields. We regret the omission of this reference.

Since publication of the manuscript, we have discovered that Richardson and Bazzoni in some very early work[2] also predicted $J\propto V^{\frac{3}{2}}$ in space-charge limited high-pressure gas diodes. Their attempts at experimental verification, however, were not conclusive.

[1] H. F. Ivey, *Advances in Electronics and Electron Physics*, edited by L. Marton (Academic Press, Inc., New York, 1954), Vol. 6, p. 137.

[2] O. W. Richardson and C. B. Bazzoni, Phil. Mag. **32**, 426 (1916).

Nuclear Quadrupole Interaction in Pure Metals, T. P. DAS AND M. POMERANTZ [Phys. Rev. **123**, 2070 (1961)]. In Eq. (2), the fifth term in the brackets should be $Y_{lm}(\theta_j\phi_j)/r_j^{l+1}$ instead of $Y_{lm}(\theta_j\phi_j)$; in the same equation, in the sixth term the factor i^l should be replaced by $(-i)^l$ and $\sum_j{}'$ should be replaced by $\sum_\lambda{}'$. The first and third of these errors were typographical errors in the manuscript. The second error has no effect on our results because l is even in our calculations. We are grateful to Dr. F. W. deWette for bringing these errors to our attention.

Charged Boson Gas, LESLIE L. FOLDY [Phys. Rev. **124**, 649 (1961)]. A factor of $\frac{1}{2}$ was inadvertently omitted in the expression for the ground-state energy U_0 as given in Eq. (12). This same factor is then absent in the constant S as defined by Eq. (18), whose numerical value is thus changed to -0.803. The final equation (22) for the ground-state energy per particle, u_0, should then have in place of $-1.606r_s^{-\frac{3}{4}}+0.425$, the expression $-0.803r_s^{-\frac{3}{4}}+0.213$.

Fission of a Hot Plasma, NORMAN ROSTOKER AND ALAN C. KOLB [Phys. Rev. **124**, 965 (1961)]. On p. 965, the second sentence of column two should read: "The ion cyclotron frequency is estimated to be $\omega_I\cong 2\times 10^8$ sec^{-1} for ions immersed in the plasma at the onset time of the instability, and the ion Larmor radius is $a_I\cong 1$ mm." On p. 969, in the first sentence of column one, "$\Omega\cong 2-5\times\Omega_0$" should read: "$\Omega\cong 0.2-0.5\times\Omega_0$."

Intermediate Meson Contributions to Hyperon Decays, DAVID R. HARRINGTON [Phys. Rev. **124**, 1290 (1961)]. Because of an algebraic error the expression given for $D(s)$ in Eq. (27) is one-half of the correct expression. The numerical results given in Tables III and IV should therefore be multiplied by two. This increases the importance of the two-meson contribution to the decays $Y\rightarrow N+\pi$, but otherwise the conclusions drawn in Sec. 5 are essentially unchanged.

Reactions of Alpha Particles with Tin-124, R. L. HAHN AND J. M. MILLER [Phys. Rev. **124**, 1879 (1961)]. Recent data[1] on the decay of the Sb^{126} isomers indicate that our cross-section measurements for the (α,pn) reaction are high by a factor of 2. That is, the 0.68-Mev γ ray, which was assumed to occur in 100% of the decays, appears to be a composite of two coincident gamma rays, of energy 0.665 and 0.695 Mev. Each of these gamma rays is reported to occur with unit probability per disintegration.

[1] *Nuclear Data Sheets* (National Academy of Sciences-National Research Council, Washington, D. C., 1961).

Relaxation Equations for Two-Magnon and Magnon-Phonon Processes in Ferrimagnetic Resonance, P. E. SEIDEN [Phys. Rev. **124**, 1110 (1961)]. The equation for n_k in the case of small spin-wave amplitudes [equation directly above Eq. (11)] should read:

$$n_k=n_0{}^{(0)}\lambda_{0k}(\lambda_{k\sigma}-\gamma\Delta H)^{-1}\times[\exp(-\gamma\Delta Ht)-\gamma\Delta H\lambda_{k\sigma}{}^{-1}\exp(-\lambda_{k\sigma}t)];$$

and Eq. (11) should read:

$$M_z-M_0=-\gamma\hbar n_0{}^{(0)}\{[1+\sum_k\lambda_{0k}(\lambda_{k\sigma}-\gamma\Delta H)^{-1}]\times\exp(-\gamma\Delta Ht)-\gamma\Delta H\sum_k\lambda_{0k}\lambda_{k\sigma}{}^{-1}\times(\lambda_{k\sigma}-\gamma\Delta H)^{-1}\exp(-\lambda_{k\sigma}t)\}.$$

Spin and Parity of the ω Meson, M. L. STEVENSON, L. W. ALVAREZ, B. C. MAGLIĆ, AND A. H. ROSENFELD [Phys. Rev. **125**, 687 (1962)]. The following "Notes Added in Proof" were returned with galleys of this paper but did not appear in the published article:

NOTES ADDED IN PROOF

I. Is the G Parity of the ω Meson -1?

A. Dalitz Plot Evidence for $G=-1$

If the true width of the 3-pion resonance were consistent with zero, one might question whether the observed $\pi^+\pi^-\pi^0$ decay mode is an "allowed" transition.