Optical Absorption of Cl₂⁻ Molecule-Ions in Irradiated Potassium Chloride*

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Pure KCl, KCl-Ag, KCl-Tl, and KCl-Pb, after x-ray irradiation at liquid nitrogen temperature, show identical Cl_2^- paramagnetic resonance spectra. The optical absorption of the Cl_2^- molecule-ion has been identified and is found to consist of two bands, one at 365 m μ and one at 750 m μ . This identification was made possible by bleaching irradiated samples with polarized light and observing the effects on the optical and also the paramagnetic resonance spectra. It was possible to reorient the Cl_2^{-} molecule-ions such that a large fraction of them were in only one of the six possible $\langle 110\rangle$ directions.

INTRODUCTION

KÄNZIG¹ and Castner and Känzig² have reported V-center type paramagnetic resonance spectra in V-center type paramagnetic resonance spectra in pure potassium chloride crystals which have been heavily irradiated with x-rays at liquid nitrogen temperature. They have identified the species giving rise to the resonance as Cl_2^- molecule-ions.

Delbecg, Smaller, and Yuster^{3,4} have observed the identical Cl2⁻ resonance spectra in KCl-Tl, KCl-Ag, and KCl-Pb crystals which have been weakly irradiated with x-rays at liquid nitrogen temperature.

Since the V_1 band⁵ in the optical absorption spectrum disappears in the same temperature range as the Cl₂resonance absorption in pure KCl, there has been some speculation² that the V_1 band is an optical absorption band of the Cl2⁻ molecule-ions; however, no positive correlation has been reported. This investigation was undertaken, using both optical absorption and paramagnetic resonance absorption measurements, to determine the optical absorption of Cl₂⁻ molecule-ions in potassium chloride crystals.

Experimental Procedures

The electron spin resonance detection system used is the same as that described previously⁶ except for two modifications. The dual modulation system was modified to operate at 25 cps and 17 kc/sec. In addition, the method of holding the sample was changed to allow orientation of the crystal in the magnetic field. The crystal was mounted on the end of a Lucite rod, the axis of which could be tipped with respect to the magnetic field. The crystal could be further oriented by rotation about this axis as indicated by a pointer at the top of the rod.

Some of the pure KCl single crystals used in these experiments were obtained from the Harshaw Chemical Company. Other single crystals of pure KCl as well as the single crystals of KCl containing either Tl⁺, Ag⁺, or Pb⁺⁺ were grown in this laboratory using the Kyropoulos method. In general, the crystals used for optical absorption measurements were $12 \times 7 \times 1$ mm, and those used for electron spin resonance measurements were $2.5 \times 2.5 \times 10$ mm. The crystals were irradiated with x-rays produced in a Machlett tube operating at 50 kv and 50 ma, and filtered through 1 mm of quartz.

Polarized light was obtained by using an Ahrens or Glan-Thompson prism. The prism was mounted on the Dewar flask used to hold the crystal. The axis of the prism was oriented parallel to one of the cubic axes of the crystal; the prism could be rotated about its axis so that the electric vector of the transmitted polarized light made any desired angle with the other two cubic axes of the crystal. Absorption spectra were taken using, in general, polarized light with the electric vector oriented in either the $\lceil 011 \rceil$ or the $\lceil 011 \rceil$ direction, the light from the monochromator passing through the polarizing prism and then through the crystal to the detector. In bleaching a crystal with polarized light, the light from an AH-6 mercury lamp was passed through an appropriate filter combination and then through a polarizing prism before impinging upon the crystal; in most bleaching experiments, polarized light with the electric vector oriented in the $\lceil 011 \rceil$ direction was used. Bleaching of samples used in electron spin resonance measurements was done while the crystal, mounted on a Lucite holder, was immersed in a Dewar flask filled with liquid nitrogen; the polarized bleaching light entered the Dewar flask through quartz windows.

EXPERIMENTAL RESULTS

Figure 1 shows the absorption spectrum of a crystal of pure potassium chloride which had been exposed to x-rays for four hours at liquid nitrogen temperature. The well-known $V_2(235 \text{ m}\mu)$, $V_1(350 \text{ m}\mu)$, and $F(540 \text{ m}\mu)$ bands are observed. The intensity of the F band was so great that the absorption could not be measured at the center of the band. The F' band is also observed as a

^{*} Based on work performed under the auspices of the U.S. Atomic Energy Commission. ¹W. Känzig, Phys. Rev. 99, 1890 (1955).

² T. Castner and W. Känzig, J. Phys. Chem. Solids 3, 178 (1957)

³ Delbecq, Smaller, and Yuster, Color Center Symposium at Argonne National Laboratory, 1956 (unpublished). ⁴ Yuster, Delbecq, and Smaller, Bull. Am. Phys. Soc. Ser. II, 2,

^{302 (1957).}

⁶ Casler, Pringsheim, and Yuster, J. Chem. Phys. 18, 1564 (1950); F. Seitz, Revs. Modern Phys. 26, 7 (1954). ⁶ Delbecq, Smaller, and Yuster, Phys. Rev. 104, 599 (1956).



FIG. 1. Absorption spectra of Argonne grown "pure" KCl: after 4 hours x-ray irradiation at liquid nitrogen temperature, - after warming to -155° C for 2 minutes and remeasurement at liquid nitrogen temperature.

broad absorption to the long-wavelength side of the F band. Figure 2 shows the absorption spectra of irradiated KCl-Tl, KCl-Ag, and KCl-Pb crystals. It will be noted that in each case there is an absorption band at about 365 m μ and that the x-ray dosage needed to obtain equivalent absorption intensity in this region of the spectrum in these three crystals is a good deal less than that needed in the pure potassium chloride. All of the above-mentioned irradiated crystals (pure KCl, KCl-Tl, KCl-Ag, and KCl-Pb) show identical Cl₂electron spin resonance spectra. In order to compare the thermal stability of the V_1 center and the Cl_2 molecule-ion in pure potassium chloride, pulse annealing experiments were carried out in which the optical absorption of the V_1 band and the resonance intensity of Cl2⁻ were followed as functions of annealing temperature. In both cases the crystals were held at liquid nitrogen temperature for measurement, warmed rapidly to the desired annealing temperature, held at this temperature for two minutes, and then cooled to liquid nitrogen temperature for remeasurement. It is quite clear from Fig. 3 that the optical absorption of the



FIG. 2. Absorption spectra after x-ray irradiation at liquid nitrogen temperature. (a) KCl-Pb, 10-min x-ray, (b) KCl-Tl, 10min x-ray, (c) KCl-Ag, 1/2-min x-ray.

 V_1 band and the spin resonance of Cl_2 have their maximum decay rates at appreciably different temperatures. The thermal stabilities are sufficiently different that holding the crystal at -155° C for a period of time causes the V_1 band to disappear while leaving the intensity of the Cl₂⁻ resonance essentially unchanged. Warming the crystal again to -155° C causes little change in the intensity of the Cl₂⁻ resonance; this result indicates that the decay of V_1 centers does not result in the formation of Cl₂⁻ centers.

Another method used to check the correlation between the V_1 center and the Cl_2^- molecule-ion was the comparison of the intensity of the Cl_2^- resonance for equal optical absorption in the region of the V_1 band for pure KCl and for KCl containing either Tl⁺, Ag⁺, or Pb++. Since in the case of irradiated KCl-Pb crystals the absorption band at 365 m μ is relatively free from overlapping neighboring absorption bands, KCl-Pb was chosen to be used as comparison with pure KCl. These two crystals were irradiated such that the intensity of the absorption in the region of the V_1 band was the



FIG. 3. Fraction of (a) V_1 absorption, and (b) Cl_2^- electron spin resonance absorption remaining after pulse anneal, plotted as a function of annealing temperature.

same in both crystals, and then the intensity of the Cl₂⁻ resonance was measured in each. The intensity of the Cl₂⁻ resonance was ten times greater in the KCl-Pb crystal. This result, as well as the results of the thermal annealing experiment, shows that the V_1 band is not an absorption band of the Cl_2^- molecule ion.

Castner and Känzig² have shown that the Cl₂molecule-ions are oriented in the six (110) directions of the KCl crystal. It therefore might be possible with the use of polarized light to bleach preferentially those Cl₂⁻⁻ molecule-ions having certain orientations. Ueta⁷ has already shown that it is possible to obtain such anisotropic bleaching in the case of *M* centers in KCl.

If we let the faces of a cube be parallel to the (001), (010), and (100) planes of a KCl lattice, then the face diagonals indicate the six possible orientations of the Cl_2^- center in the lattice. It is possible that bleaching a crystal containing Cl_2^- centers with $\lceil 0\overline{11} \rceil$ polarized light⁸ of the appropriate wavelength will excite all Cl₂-

⁷ M. Ueta, J. Phys. Soc. Japan 7, 107 (1952). ⁸ $[0\overline{11}]$ polarized light indicates polarized light with the electric vector oriented parallel to $[0\overline{11}]$ in an alkali halide lattice.

centers except those oriented in [011]. The result of such an excitation might be that the number of Cl₂centers oriented in [011] remains constant and that the number of Cl₂⁻ centers in the other five orientations will decrease; or, if the excited Cl_2^- center is not very mobile and loses its excitation energy before annihilation, there will be a reorientation of the Cl_2^- centers into [011] from the other five directions. Since the intensity of the Cl₂⁻ resonance obtained upon irradiation is greatest in the case of the crystals with added impurities, it was convenient to do most of the polarized light experiments with these crystals rather than with pure KCl. After x-ray irradiation, the optical absorption of these crystals was measured with [011] and also $[0\overline{1}1]$ polarized light. No absorption measurements were made at wavelengths lower than 290 m μ because of the high absorption of the polarizing prism at lower wavelengths. Bleaching with light was usually done with [011] polarized light. After bleaching, the absorp-



FIG. 4. Absorption spectra of KCl-Pb (a) after 10-min x-ray irradiation at liquid nitrogen temperature as measured with either [011] or [011] polarized light; and then after bleaching with [011] polarized light (436 m μ) as measured with (b) [011] polarized light, and (c) [011] polarized light.

tion of the crystals was again measured with $\lceil 011 \rceil$ as well as [011] polarized light. Figure 4 shows the results of such a bleaching experiment carried out at liquid nitrogen temperature on a KCl-Pb crystal. The curve labeled [011], [011] shows the absorption spectrum taken immediately after x-ray irradiation; the absorption spectrum is the same whether taken with [011] or [011] polarized light. The effect of bleaching with [011] polarized light (436 m μ) is shown by the upper and lower curves. The absorption spectrum taken with [011] polarized light shows the $365\text{-m}\mu$ band to be increased by over a factor of 2 while the spectrum taken with $[0\overline{1}1]$ polarized light shows the 365-mµ band to be decreased by more than a factor of 5. These results show that, although there may have been some "real bleaching" due to recombination of electrons and holes, the principal effect of this bleach with $\lceil 0\overline{1}1 \rceil$ polarized light is to cause a reorientation of a very large fraction



FIG. 5. Subtraction curves, after normalization to give the same peak optical density, for \bigcirc KCl-Pb, \times KCl-Ag, and \square KCl-Tl, showing dichroism introduced by bleaching with [011] polarized light absorbed by the 365-m μ band.

of the centers such that the dipole moment associated with the transition giving rise to the $365\text{-m}\mu$ band is parallel to either [011] or [111]. That the dipole moment is oriented in the [011] direction is shown by the fact that bleaching with [001] polarized light in the $365\text{-m}\mu$ band produces a dichroism in this band, as observed in the absorption spectrum taken with [001] and [010] polarized light.

If the [011] absorption is subtracted from the [011] absorption, the absorption spectrum obtained is due to those centers which bleach in an anisotropic manner. Any center which shows no anisotropy contributes the same absorption to both the [011] and the [011] absorptions and therefore is eliminated in a subtraction. Figure 5 shows the subtraction curve not only for irradiated KCl-Pb but also for KCl-Ag and KCl-Tl obtained in the same manner; the data were normalized to give the same optical density at the peak in each case. These subtraction curves, after normalization, are identical within experimental error, indicating that a common species gives rise to the absorption band in each of these crystals.

The final proof that the Cl_2^- centers give rise to the $365\text{-m}\mu$ band rests on a correlation of the optical bleaching experiments with changes in the resonance spectra. A crystal of KCl-Pb was irradiated with x-rays at liquid nitrogen temperature so that the Cl₂⁻ resonance was readily observed. With Cl_2^- centers distributed at random among the six (110) directions and with the magnetic field parallel to [011], a resonance spectrum similar to that shown in Fig. 6(A) is observed. As pointed out by Castner and Känzig,2 this spectrum consists of three basic patterns-0°, 60°, and 90° (the angle indicated is that between the magnetic field and the molecular axis). Figure 6(B) shows the resonance spectrum after the crystal had been exposed to $[0\overline{1}1]$ polarized light absorbed in the 365-m μ band. The 0° pattern has been enhanced considerably after such a bleach (the ratio of the intensity of the 0° pattern to that of the 60° pattern has changed by a factor of 16), and therefore a large fraction of the Cl₂- centers have been reoriented such that the molecular axis is parallel to $\lceil 011 \rceil$. The correlation between the behaviors of the



FIG. 6. Electron spin resonance spectrum $(H \parallel [011])$ of a KCl-Pb (A) immediately after x-raying at liquid nitrogen temperature, and (B) after bleaching with $[0\overline{11}]$ polarized light absorbed by the 365-m μ band.

 $365\text{-}m\mu$ band and the Cl_2^- resonance spectrum shows that the Cl_2^- centers give rise to the $365\text{-}m\mu$ absorption band. It can be further concluded that the dipole moment associated with the transition giving rise to the $365\text{-}m\mu$ band is oriented parallel to the molecular axis of the Cl_2^- center.

Theory^{9,2} predicts that the Cl_2^- molecule-ion has, in addition to the absorption in the ultraviolet, two absorptions at longer wavelengths. A weak absorption band associated with Cl2⁻ centers and having about 1/60 of the intensity of the band in the ultraviolet was found with peak at 750 m μ (Fig. 7). Both KCl-Ag and KCl-Pb after x-irradiation at liquid nitrogen temperature show this peak in the red; however, in the case of KCl-Tl the absorption due to the trapped electron is sufficiently intense in this region that it masks the long wavelength absorption of the Cl₂⁻ center. Upon bleaching a KCl-Pb or a KCl-Ag crystal which contains Cl₂⁻ centers with $\lceil 0\overline{11} \rceil$ polarized light in the 750-mµ band, no dichroism is produced in the crystal. However, that this long-wavelength absorption band is associated with Cl₂⁻ centers is shown by the following experiment. After a KCl-Pb or a KCl-Ag crystal which contains Cl_2^- centers is bleached with polarized light absorbed in the $365-m\mu$ band so as to produce a dichroism in that band, it is found that subsequent bleaching with light absorbed in the 750-mµ band causes the dichroism

to disappear. It should be pointed out at this time that the stimulation peaks³ in the infrared, observed in KCl-Tl crystals after x-ray irradiation at liquid nitrogen temperature, cannot be associated with Cl_2^- centers as suggested by Castner and Känzig,² since these stimulation peaks are still present after all the Cl_2^- centers have been destroyed.

The V_1 band in pure KCl was also bleached with polarized light in experiments similar to those performed on the 365-m μ band. Only a very slight anisotropy occurred in bleaching with either [001] or [011] polarized light. It is quite clear that the behavior of the V_1 band is not the same as that of the 365-m μ band. However, the slight anisotropy which developed on bleaching showed optical properties in the ultraviolet and infrared characteristic of Cl_2^- . The intensity of this anisotropic absorption was about that expected from the intensity of Cl_2^- as measured in paramagnetic resonance on a similar sample. It appears that the absorption in this region of the spectrum consists of two bands, the major absorption being due to V_1 centers and another smaller absorption being due to Cl_2^-



FIG. 7. 750-m μ band absorption introduced into KCl-Ag by 2 hours x-ray irradiation at liquid nitrogen temperature.

centers. In irradiated pure KCl, the Cl_2^- centers represent only a small fraction of the hole centers formed, while in the case of the irradiated doped crystals it is believed that the Cl_2^- centers account for essentially all of the hole centers.

Upon x-raying a KCl-Tl, KCl-Pb, or KCl-Ag crystal at liquid helium temperature, it is found that $Cl_2^$ centers are formed just as rapidly as at liquid nitrogen temperature. In addition, the 365- and 750-m μ $Cl_2^$ bands, when exposed to polarized light, behave the same at liquid helium temperature as at liquid nitrogen temperature.

Additional information was obtained through a study of the thermal stability of the Cl_2^- center. Upon pulse annealing a crystal of KCl-Ag, KCl-Pb, or KCl-Tl which contains preferentially oriented Cl_2^- centers, it is observed that the dichroism begins to disappear at about -100° C; at this temperature the Cl_2^- centers are just beginning to move through the lattice at an appreciable rate and as a result of the jumping lose their preferred orientation. However, it is believed that at -100° C the Cl_2^- center does not move very rapidly

⁹ M. H. Cohen, Phys. Rev. 101, 1432 (1956); Inui, Harasawa, and Obata, J. Phys. Soc. Japan 11, 612 (1956).

through the crystal, since no appreciable decrease in the concentration of Cl₂⁻ centers is observed until about -70° C. It is observed that the maxima in the rates of decay of both the Cl₂⁻ resonance and the 365-m μ band occur at about -65°C in KCl-Ag or KCl-Tl, and at -55° C in KCl-Pb. At the same temperature at which the Cl₂⁻ centers disappear, a glow peak¹⁰ is observed; the spectral distribution of the luminescence depends on the impurity present in the crystal (Tl+, Ag+, Pb++). It is believed that these impurity ions emit as a result of the recombination of trapped electrons with wandering Cl_2^- centers.

DISCUSSION

Three different types of experiments reported above show that the V_1 band is not the optical absorption of Cl_2^- molecule-ions. (1) The V_1 band and Cl_2^- resonance in pure KCl do not decay at the same temperature. (2) The intensity of the Cl_2^- resonance absorption is not proportional to the intensity of the V_1 band. (3) The V_1 band does not bleach in an anisotropic manner with $\lceil 0\overline{1}1 \rceil$ or $\lceil 001 \rceil$ polarized light. This last result is in agreement with those obtained by Lambe and West.11

Castner and Känzig² have pointed out that two models which would be consistent with the Cl₂⁻ resonance absorption observed in pure KCl are (1) Cl₂with no vacancies, (2) Cl_2^- with two positive-ion vacancies. Their conclusion, as well as ours, is that in the case of KCl crystals with added impurities it is highly unlikely that the rate of production of positiveion vacancies could be as great as the observed rate of production of Cl2⁻ centers, and therefore the twovacancy model is improbable. Also we have found that the rate of production of Cl_2^- is essentially the same at liquid helium temperature as it is at liquid nitrogen temperature, which fact also refutes the two-positiveion vacancy model. The bleaching experiments with polarized light show perhaps as well as any of the other experiments that vacancies are not involved. In order for a reorientation to take place, if two positive-ion vacancies are associated with the Cl_2^- , it is necessary (a) that there be many other pairs of positive-ion vacancies that have no Cl_2^- attached or (b) that the pair of positive-ion vacancies reorient as well as the Cl_2 . The first alternative would require an exceedingly large concentration of pairs of positive-ion vacancies, while the second would require the movement of positive-ion vacancies at liquid helium temperature due to excitation with 750-m μ light; both of these possibilities seem unlikely. In the case of Cl_2^- molecule-ions formed in the potassium chloride crystals containing added impurities, there is another possible modela Cl_2^- with an impurity associated with it. Since the

resonance spectra are all identical, it means that the electric fields closely surrounding the Cl₂ are very similar. It is possible that Ag⁺ and Tl⁺ have electric fields which are so similar to K⁺ that their presence cannot be detected in the resonance spectra. However, in the case of KCl-Pb this possibility does not seem likely, since lead is believed to be present as a divalent ion. Whether or not the vacancy which enters the crystal along with the lead is directly associated with the Pb⁺⁺, the neighborhood of the Cl₂⁻ should probably be sufficiently distorted to affect the electron spin resonance spectrum. Since no difference is observed, it is concluded that the Cl₂- has no impurities as nearest neighbors.12

Two other sets of experimental data, (1) that the optical absorption of Cl_2^- is the same in KCl-Ag, KCl-Tl, and KCl-Pb, and (2) that Cl_2^- molecule ions begin to move through the lattice at the same temperature in the doped crystals, indicate further that the Cl₂⁻ molecule ions are located in regions in the lattice characteristic of pure KCl and not influenced by the added impurities.

The experimental facts that the maximum rate of disappearance of the Cl₂⁻ molecule-ion occurs at -140°C in pure KCl, at -65°C in KCl-Tl and KCl-Ag, and at -55° C in KCl-Pb appear to be inconsistent with the conclusion that the Cl_2^- molecule-ion has no impurities as nearest neighbors. The following is offered tentatively as an explanation for these differences. Since the reorientation experiments show that Cl₂begins to move through the lattice at the same temperature in KCl-Pb, KCl-Tl, and KCl-Ag, it is suggested that the difference in decay temperature in KCl-Pb is not because of a difference in the Cl2⁻ but rather a difference in the retrapping and recombination processes involving the mobile Cl2-. No successful reorientation experiments have been performed on pure potassium chloride because of the low concentration of Cl_2 centers; consequently no reorientation temperature has been measured. Nevertheless, it is believed that in pure KCl, the Cl₂⁻ molecule-ion does not have mobility at -140 °C but is destroyed by some electron which is released at this temperature.

The changes in the Cl₂⁻ resonance absorption spectrum along with the changes in the optical absorption spectrum, which are produced by irradiation with polarized light, show that the Cl2- molecule-ion has two optical absorption bands with peaks at 365 m μ and 750 m μ . Absorption of light in either of these two bands leads to an excited state from which a Cl₂molecule-ion reforms, oriented at random in one of the six (110) directions in the crystal. The result that very little real bleaching occurs after excitation suggests

¹⁰ These glow peaks and the processes which take place at higher temperatures will be the subject of a later paper. ¹¹ J. Lambe and E. J. West, Phys. Rev. **108**, 634 (1957).

¹² In irradiated KCl-Ag, electron spin resonance spectra have been observed which arise from the presence of a hole trapped at and also from an electron trapped at Ag⁺(Ag⁰). These results will be included in a later paper.

that the Cl_2^- molecule-ion reforms very close to its original neighborhood. Therefore, when a potassium chloride crystal is irradiated with x-rays at liquid nitrogen temperature, it seems likely that a positive hole which has been created moves only a very short distance before forming a Cl_2^- molecule-ion. Since the rate of formation of positive holes by x-ray irradiation should be the same in pure potassium chloride as in doped potassium chloride, Cl_2^- molecule-ions should be formed at the same rate in each of these crystals; yet it is observed that the net rate of formation of $Cl_2^$ centers is from 100 to 1000 times greater in the case of the doped crystals. It is believed that the added impurities act as very efficient electron traps,¹² decreasing the rate of recombination of electrons and positive holes, while in the pure crystals electron traps must be created slowly by the irradiation and consequently the recombination rate is much greater.

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Control of Luminescence by Charge Extraction*

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It has been discovered that the application of a potential of a few volts to phosphors of the ZnS group can quench the fluorescence. The effect has been observed in single-crystal and recrystallized wafers of the phosphors by illuminating one side with nonpenetrating light and applying a field, through suitable contacts, normal to the illuminated surface. With the illuminated surface negative, the fluorescence is reduced by an amount dependent on the voltage for as long as the voltage is maintained, and total extinction can often be obtained. The luminescent output follows variations in the applied voltage up to a maximum frequency which depends on the fluorescence time constant.

The effect is shown to be a fundamental property of phosphor materials possessing large differences in hole and electron mobilities or capture cross section. It is explained in terms of a decrease in the excess electron density due to the extraction of holes from the illuminated surface by the field. A simple mathematical theory is proposed which accounts in detail for the observed effects and which predicts current-voltage relationships which have been confirmed experimentally.

I. INTRODUCTION

I N the literature on phosphors there are many reports on effects of applying an electric field to photoluminescent materials. These include the Gudden-Pohl effect¹ (the transient stimulation of luminescence), Dechene's² experiment on the quenching of luminescence with a high field, and the Destriau³ effect of stimulation of luminescence by an alternating field. Halsted⁴ has observed modulation of fluorescence by an alternating field and explained the effect by assuming that the optically created hole-electron pairs are first separated and then brought back together under the influence of the field. These, and a number of less clearly defined effects have been reviewed recently in an article by Ivey.⁵

The purpose of this paper is to describe a new electrophotoluminescent effect and to propose an explanation for it. The effect is unusual in that dc as well as ac fields can be used, that these fields are very low, and that the effect produced can be very large, almost complete extinction or significant stimulation of fluorescence being obtained depending on the direction of the field.

In Sec. II the experimental arrangements are described. In Sec. III experiments are described which essentially repeat those of Halsted except that singlecrystal phosphor material was used instead of a phosphor powder in a dielectric matrix. This allowed the field modulation effect to be studied in a more controlled manner. The results observed when a transparent conducting coating is placed in direct contact with the single crystals (and also with sintered layers of phosphor material) are described in Sec. IV and the interpretation of these results is discussed. A mathematical description of the observed effects is presented in an Appendix.

II. EXPERIMENTAL METHOD

In the majority of experiments, wafers of singlecrystal cadmium sulfide were used. These were activated by heating in the presence of cupric chloride and

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 ¹ B. Gudden and R. W. Pohl, Z. Physik 2, 192 (1920).
² G. Dechene, Compt. rend. 201, 139 (1935).

³ M. Destriau, Compt. rend. 238, 2298 (1954).

⁴ R. E. Halsted, Phys. Rev. 99, 1897 (1955).

⁵ H. F. Ivey, J. Electrochem. Soc. 104, 740 (1957).