

**Thallium-203.**—Extensive measurements of the internal conversion coefficients of the 279-keV transition in  $Tl^{203}$  have recently been made by several groups of workers.<sup>3,4</sup> Starting with experimental values for  $\alpha^K$ ,  $\alpha^{L_I}$ , and  $\alpha^{L_{II}}$  and with the assumption  $G_K = G_{L_I} = G_{L_{II}}$  they found  $E2/M1 = 1.38 \pm 0.25$  and  $G_K = 0.53 \pm 0.08$ . Since this reduction factor depends rather decisively on the value of  $E2/M1$ , we felt that a measurement of  $E2/M1$  independent of internal conversion coefficients would be desirable. The measured angular distribution of the 279-keV  $\gamma$  rays following Coulomb excitation could be fitted equally well by  $(E2/M1)^{\frac{1}{2}} = 1$  to 2. However, a polarization-direction measurement is very sensitive to this range of  $(E2/M1)^{\frac{1}{2}}$  for a transition of the type  $3/2(E2+M1)^{\frac{1}{2}}$  and the value observed for  $E2/M1$  is listed in Table I.

**Rhenium-187, -185.**—These transitions are of limited value as evidence for a reduction in  $\beta_1^K$  because of the uncertainty in  $E2/M1$ . The determination of the  $E2/M1$  value from the angular distribution is unfavorable for these transitions because the transition of the type  $7/2(E2+M1)5/2$  is nearly isotropic for a wide range of  $E2/M1$ . The angular distributions have been measured and they are found to be isotropic. A  $K/L$  measurement has been made for the transition in  $Re^{187}$ ,<sup>12</sup> and the  $E2/M1$  value of 1/9 is based on this measurement.

<sup>1</sup> Rose, Goertzel, Spinrad, Harr, and Strong, *Phys. Rev.* **83**, 79 (1951) and "Tables of internal conversion coefficients" (privately circulated by M. E. Rose).

<sup>2</sup> L. A. Sliv, *Zhur. Eksptl. i Teort. Fiz.* **21**, 77 (1951); L. A. Sliv and M. A. Listengarten, *Zhur. Eksptl. i teort. Fiz.* **22**, 29 (1952).

<sup>3</sup> A. H. Wapstra and G. J. Nijgh, *Nuclear Phys.* **1**, 245 (1956).

<sup>4</sup> Nordling, Siegbahn, Sokolowski, and Wapstra, *Nuclear Phys.* **1**, 326 (1956).

<sup>5</sup> F. K. McGowan, *Phys. Rev.* **93**, 471 (1954).

<sup>6</sup> H. Paul and R. M. Steffen, *Phys. Rev.* **98**, 231 (1955).

<sup>7</sup> Heer, Ruetschi, Gimmi, and Kundig, *Helv. Phys. Acta* **28**, 336 (1955).

<sup>8</sup> P. H. Stelson and F. K. McGowan, *Bull. Am. Phys. Soc. Ser. II*, **1**, 264 (1956).

<sup>9</sup> F. K. McGowan and P. H. Stelson, *Phys. Rev.* **99**, 127 (1955).

<sup>10</sup> P. H. Stelson and F. K. McGowan, *Bull. Am. Phys. Soc. Ser. II*, **1**, 164 (1956).

<sup>11</sup> J. V. Kane and S. Frankel, *Bull. Am. Phys. Soc. Ser. II*, **1**, 171 (1956).

<sup>12</sup> Cork, Brice, Nester, LeBlanc, and Martin, *Phys. Rev.* **89**, 1291 (1953).

### Beta Decay of a $C^9$ Nucleus\*

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**I**N a systematic survey of photographic emulsions, exposed to 3-Bev protons, for excited nuclear fragments, a connected double star was found which is interpreted to be the disintegration of a  $C^9$  nucleus. It was thought worthwhile to describe the event in detail

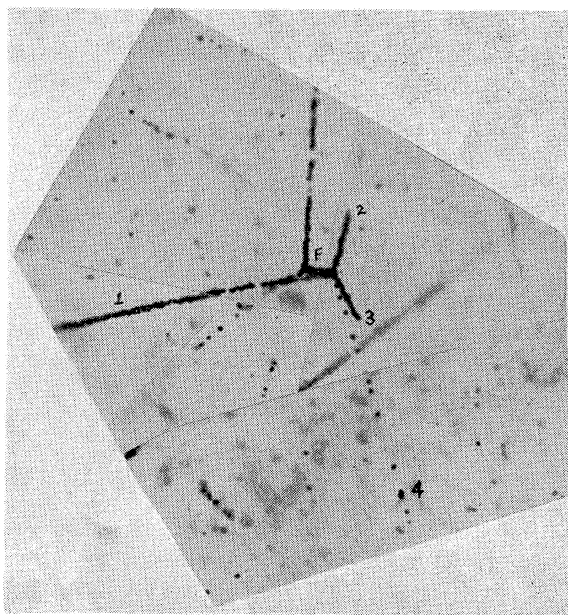


FIG. 1. A photograph of an event interpreted as the beta decay of  $C^9$ . The  $C^9$  nucleus (track  $F$ ) was produced in star ( $A$ ) and disintegrated into a proton, two alpha particles, and a positron (tracks 1, 2, 3, and 4, respectively).

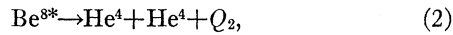
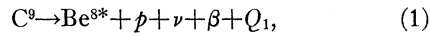
since there has been no evidence for a long-lived  $C^9$  nucleus.

A photograph of the connected stars is shown in Fig. 1. The primary star ( $A$ ), which was probably produced by a neutron, has three outgoing tracks. Track  $F$  is saturated, 6.2 microns long and was caused by a slow multiply-charged particle. The absence of  $\delta$  rays and the presence of some visible scattering along  $F$  suggest that the particle came to rest before it gave rise to the secondary star ( $B$ ). The secondary star which appears at the end of track  $F$  consists of four charged particles. The main characteristics of the secondary star and the fragment are given in Table I. Measurements of multiple Coulomb scattering and grain density along track 4 indicate that it was caused by a 3.1-Mev electron. Tracks 1, 2, and 3 are coplanar to within one degree. The coplanarity strongly suggests that the fragment that caused track  $F$  came to rest before it decayed and that no neutrons were involved in the decay. Since track 1 was produced by a singly-charged particle, the coplanarity and the momentum balance uniquely determine that the particles which produced

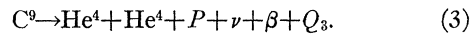
TABLE I. Characteristics of  $C^9$  decay.

Track	Range in microns	Identification	Energy in Mev	Angles
$F$	6.2	$C^9$	9.4	$65^\circ$
1	341.0	P	7.6	$115^\circ$
2	9.9	$He^4$	2.7	
3	7.8	$He^4$	2.1	$136.5^\circ$
4	...		3.1	$11^\circ$

tracks 1, 2, and 3 were a proton and two alpha particles, respectively. With this assignment, the momentum unbalance for tracks 1, 2, and 3 is  $15 \pm 10$  Mev/ $c$ . Depending on the sign of the charge of the beta particle, the fragment (track  $F$ ) was either  $\text{Be}^9$  or  $\text{C}^9$ . Since  $\text{Be}^9$  is stable, the fragment was probably a  $\text{C}^9$  nucleus and the beta particle was positive. The possible decay schemes are



or



The similarity in the energy of the two alpha particles

suggests the formation of  $\text{Be}^8$  in an excited state, in which case  $Q_2$  is found to be 3.8 Mev.

For the decay scheme (3), the value of  $Q_3$  is found to be greater than 15.4 Mev. Of course the energy of the neutrino cannot be measured. However, if  $Q_3$  were greater than 16.4 Mev,  $\text{C}^9$  would be unstable against decay into  $\text{B}^8$  and a proton. Hence the maximum energy of the neutrino is 1.0 Mev. The limits on the mass of  $\text{C}^9$  are 8408.6 and 8409.6 Mev.

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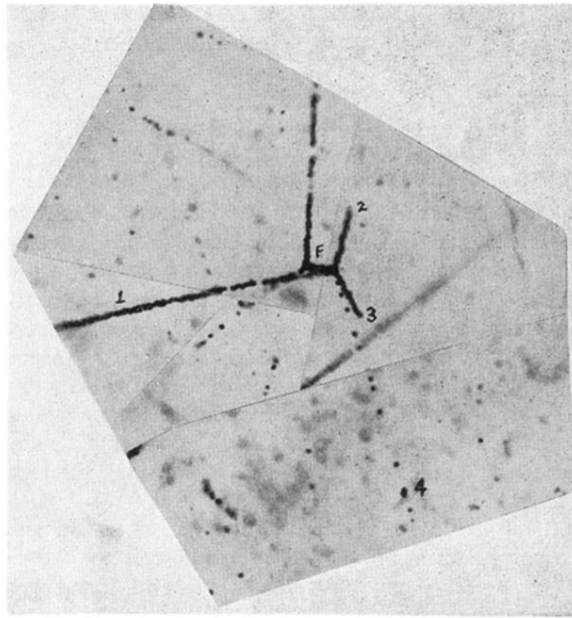


FIG. 1. A photograph of an event interpreted as the beta decay of  $C^{13}$ . The  $C^{13}$  nucleus (track  $F$ ) was produced in star ( $A$ ) and disintegrated into a proton, two alpha particles, and a positron (tracks 1, 2, 3, and 4, respectively).